

Non-Associative Geometry of Quantum Tori

Francesco D'ANDREA ^{†‡} and Davide FRANCO [‡]

[†] *Dipartimento di Matematica e Applicazioni, Università di Napoli "Federico II",
Complesso MSA, Via Cintia, 80126 Napoli, Italy*

[‡] *I.N.F.N. Sezione di Napoli, Complesso MSA, Via Cintia, 80126 Napoli, Italy*
E-mail: francesco.dandrea@unina.it, dfranco@unina.it

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Abstract. We describe how to obtain the imprimitivity bimodules of the noncommutative torus from a “principal bundle” construction, where the total space is a quasi-associative deformation of a 3-dimensional Heisenberg manifold.

Key words: noncommutative torus; quasi-Hopf algebras; cochain quantization

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1 Introduction

In differential geometry a standard way to construct vector bundles is from a principal bundle and a representation of the structure group. This construction works in noncommutative geometry too [6, 13, 15], with spaces replaced by algebras, vector bundles replaced by finitely generated projective modules, structure groups replaced by Hopf algebras or compact quantum groups, and principal bundles replaced by algebra extensions with suitable additional properties (see, e.g., [3, 14]). When the structure group is $U(1)$, it is possible to reconstruct the total space of the “bundle” (more precisely, a strongly graded C^* -algebra) from the base space and a noncommutative “line bundle” (a self-Morita equivalence bimodule), cf. [2] (see also [1, 7]).

A case study is provided by the C^* -algebra of the noncommutative torus A_θ , $\theta \in \mathbb{R} \setminus \mathbb{Q}$ [5, 20]. The group $SL_2(\mathbb{Z})$ acts on $\mathbb{R} \setminus \mathbb{Q}$ by fractional linear transformations,

$$g\theta := \frac{a\theta + b}{c\theta + d}, \quad \forall \theta \in \mathbb{R} \setminus \mathbb{Q} \quad \text{and} \quad g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{Z}), \quad (1.1)$$

and for every g as above there is a full right Hilbert A_θ -module $E_g(\theta)$ which is a Morita equivalence bimodule between A_θ and the algebra $\text{End}_{A_\theta}(E_g(\theta)) \simeq A_{g\theta}$ [5, 21]. The equivalence class of the right module $E_g(\theta)$ only depends on the ‘degree’ and ‘rank’ of the module, defined by $\text{deg}(E_g(\theta)) = c$ and $\text{rank}(E_g(\theta)) = c\theta + d$ (see, e.g., [19]). Every other finitely generated projective right A_θ -module is a direct sum of such modules.

In this paper we show how one can formally derive some of these bimodules – those with rank 1 for $\theta \rightarrow 0$ (hence, line bundles) – from some kind of $U(1)$ -bundle, but at the price of working with quasi-associative algebras, i.e., monoids in the category of representations of a quasi-Hopf algebra. Several well-known properties of these bimodules can be shown to be a consequence of the explicit form of the coassociator (cf. Lemma 3.1), including the associativity property of the pairing of bimodules of Polishchuk and Schwarz [19].

The plan of the paper is the following. In Section 2 we recall some basic definitions about twist deformations of quasi-Hopf algebras and algebra modules. In Section 3 we will introduce the twist we are interested in, based on the universal enveloping algebra of the 3-dimensional Heisenberg Lie algebra $\mathfrak{h}_3(\mathbb{R})$; we will also prove some generalized associativity property of

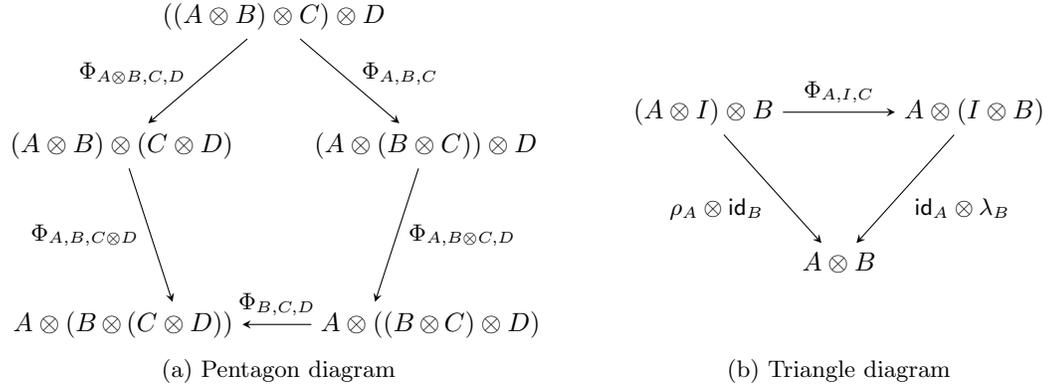


Figure 1

twisted module algebras, namely Proposition 3.2. In Section 4 we will give a concrete realization of $\mathfrak{h}_3(\mathbb{R})$ in terms of differential operators on the total space of a principal $U(1)$ -bundle $M_3 \rightarrow \mathbb{T}^2$ on the 2-torus, and apply to it the deformation recipe of Section 2 with the twist introduced in Section 3: the total space becomes a quasi-associative \mathbb{Z} -graded algebra, which in degree 0 is the algebra of smooth functions on the noncommutative torus; in degree $n \neq 0$ we get bimodules that we compare with Connes–Rieffel imprimitivity bimodules. In Section 5 we give a slightly different version of the construction, more natural if one is interested in complex structures: we show that, in this case, the twisted algebra of functions on the Heisenberg manifold has an associative commutative subalgebra given by ordinary theta functions. In Section 6 we study vector bundles of any rank.

Notations. By a *algebra* we shall always mean a unital algebra (not necessarily associative nor commutative) over a commutative unital ring R ; the algebraic tensor product over R will be denoted by \otimes_R , or simply \otimes if there is no risk of confusion; by a Hopf algebra over $\mathbb{C}[[\nu]]$ we shall always mean a *topological* Hopf algebra, completed in the h -adic topology, and by $\otimes_{\mathbb{C}[[\nu]]}$ the completed tensor product [4, Section 4.1A].

2 Mathematical preliminaries

In this section, we recall some definitions and properties of quasi-Hopf algebras and Drinfeld twists, from [4, 10, 11, 16].

2.1 Monoidal categories

A *monoidal category* is a category \mathbf{C} equipped with a functor $\otimes: \mathbf{C} \times \mathbf{C} \rightarrow \mathbf{C}$ which, modulo natural isomorphisms, is associative and unital [4, Section 5.1]. More precisely, there is an object I and three natural isomorphisms Φ between the functors $(- \otimes -) \otimes -$ and $- \otimes (- \otimes -)$, λ between $I \otimes -$ and the identity, ρ between $- \otimes I$ and the identity $-$ such that the diagrams 1a and 1b commute for all objects A, B, C, D . Examples are the category of modules over a field (vector spaces), over a group or a Hopf algebra (representations), over a topological Hopf algebra (with completed tensor product).

A *monoid* [16, Section VII.3] in a monoidal category is an object A together with two arrows $m: A \otimes A \rightarrow A$ and $\eta: I \rightarrow A$, the “multiplication” and “unit”, satisfying the usual axioms of a unital associative algebra modulo natural transformations, namely the diagrams 2a and 2b must commute. Monoids in the category of vector spaces are associative algebras, in the category of representations of a Hopf algebra H are H -module algebras.

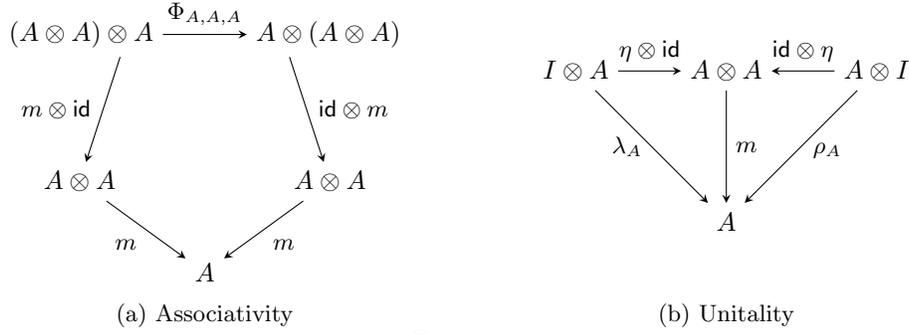


Figure 2

2.2 Quasi-Hopf algebras

Let H be an algebra and $\Delta: H \rightarrow H \otimes H$ a linear map. For all $n \geq m$ and all $1 \leq i_1 < i_2 < \dots < i_m \leq n$, we will denote by $h \mapsto h_{i_1 \dots i_m}$ the linear map $H^{\otimes m} \rightarrow H^{\otimes n}$ defined on homogeneous tensors $h = a_1 \otimes a_2 \otimes \dots \otimes a_m$ as follows: we put a_k in the leg i_k of the tensor product for all $k = 1, \dots, m$, and fill the additional $n - m$ legs with 1. The subscript $(i_k i_{k+1})$ in parenthesis means that we apply Δ to a_k and put the first leg of $\Delta(a_k)$ in position i_k and the second in position i_{k+1} . So, for example, if $h = a \otimes b$, $h_{13} = a \otimes 1 \otimes b$, $h_{(12)3} = \Delta(a) \otimes b$, etc.

A *quasi-bialgebra* (over a commutative ring R) is an associative algebra H together with two homomorphisms $\Delta: H \rightarrow H \otimes H$ and $\epsilon: H \rightarrow R$ (the ‘‘coproduct’’ and ‘‘counit’’) and an invertible element $\Phi \in H \otimes H \otimes H$ (the ‘‘coassociator’’) satisfying the following conditions (here we follow the notations of [10, 11]):

$$\Phi_{12(34)} \Phi_{(12)34} = \Phi_{234} \Phi_{1(23)4} \Phi_{123}, \quad (2.1)$$

$$(\text{id} \otimes \epsilon \otimes \text{id})(\Phi) = 1, \quad (2.2)$$

and

$$(\text{id} \otimes \Delta)\Delta(h) = \Phi(\Delta \otimes \text{id})\Delta(h)\Phi^{-1}, \quad (2.3)$$

$$(\text{id} \otimes \epsilon)\Delta(h) = h = (\epsilon \otimes \text{id})\Delta(h),$$

for all $h \in H$. A *quasi-Hopf algebra* is a quasi-bialgebra satisfying an additional condition which for $\Phi = 1$ reduces to the existence and bijectivity of an antipode, see, e.g., [4, Definition 16.1.1].

Any Hopf algebra is a quasi-Hopf algebra with trivial coassociator, $\Phi = 1$. In fact, in the Hopf case any Φ which is *invariant* (i.e., commutes with the image of the iterated coproduct) does the job. If Φ is invariant, (2.3) reduces to the coassociativity condition $(\text{id} \otimes \Delta)\Delta = (\Delta \otimes \text{id})\Delta$ and we get the usual definition of bialgebra/Hopf algebra.

If H is a commutative Hopf algebra, (2.1) can be interpreted as a 3-cocycle condition in the Hopf algebra cohomology of H [17, Section 2.3]. In general, if A, B, C are three H -modules, we may think of Φ as a module map $(A \otimes B) \otimes C \rightarrow A \otimes (B \otimes C)$, with the natural isomorphism of vector spaces understood. Condition (2.1) becomes the associativity condition 1a for the category of H -modules, which is then monoidal with unit object $I = R$ given by the ground ring (with module structure being given by the counit, so that (2.2) implies the commutativity of the diagram 1b, and module structure on a tensor product defined by the coproduct).

A monoid A in this category is a H -module algebra, that is an algebra A with multiplication satisfying the quasi-associativity condition (cf. diagram 2a):

$$m(m \otimes \text{id}) = m(\text{id} \otimes m)\Phi, \quad (2.4)$$

Commutativity of the diagram 2b means that A is unital, with $1_A := \eta(1_R)$. The condition that m and η are morphisms in the category, i.e., H -module maps, gives $h.1_A = \epsilon(h)1_A$ and

$$h \circ m = m \circ \Delta(h),$$

for all $h \in H$. Equation (2.4) becomes the usual associativity condition if the action of Φ on $A \otimes A \otimes A$ is trivial (a sufficient condition of course is that $\Phi = 1$).

The above definitions remain valid if the ground ring is $R = \mathbb{C}[[\nu]]$, with \otimes_R the completed tensor product.

The advantage of quasi-bialgebras (resp. quasi-Hopf algebras) over ordinary bialgebras (resp. Hopf algebras) is that there is a “gauge action” of the group of invertible elements in $H \otimes H$ that doesn't change the category of modules.

Let H be a quasi-bialgebra (resp. quasi-Hopf algebra) and $F \in H \otimes H$ an invertible element satisfying $(\text{id} \otimes \epsilon)(F) = (\epsilon \otimes \text{id})(F) = 1$. One can define a new quasi-bialgebra (resp. quasi-Hopf algebra) H_F given by H as an algebra and with the same counit, but with a new coproduct Δ_F and coassociator Φ_F defined by

$$\begin{aligned}\Delta_F(h) &:= F\Delta(h)F^{-1}, \quad \forall h \in H, \\ \Phi_F &:= F_{23}F_{1(23)}\Phi(F^{-1})_{(12)3}(F^{-1})_{12}.\end{aligned}$$

It turns out that the categories of H -modules and H_F -modules are equivalent as monoidal categories. In particular, if (A, m) is a H -module algebra, there is a H_F -module algebra (A, m_F) given by A as a vector space, with the same unit element, and with multiplication:

$$m_F := m \circ F^{-1}. \tag{2.5}$$

We will refer to H_F as a twist deformation of H , and to F as a *twisting element* based on H .

2.3 Quantum universal enveloping algebras

By a *deformation* of a quasi-bialgebra over a field \mathbb{k} we mean a topological quasi-bialgebra H_ν over $\mathbb{k}[[\nu]]$ such that $H_\nu/\nu H_\nu \simeq H$ as quasi-bialgebras and $H_\nu \simeq H[[\nu]]$ as $\mathbb{k}[[\nu]]$ -modules. If H is a quasi-Hopf algebra, any deformation as a quasi-bialgebra is in fact a quasi-Hopf algebra [4, Section 16.1C].

If $H = \mathcal{U}(\mathfrak{g})$ is the universal enveloping algebra of a Lie algebra \mathfrak{g} , with standard Hopf algebra structure and trivial coassociator, a deformation H_ν with coassociator $\Phi_\nu \equiv 1 \pmod{\nu^2}$ will be called a *quantum universal enveloping algebra* (or QUEA). Any twist of $\mathcal{U}(\mathfrak{g})$ by a twisting element F_ν based on $\mathcal{U}(\mathfrak{g})[[\nu]]$ (and satisfying $F_\nu \equiv 1 \pmod{\nu}$) is a QUEA, and roughly speaking every QUEA arises in this way (see Theorem 16.1.11 of [4] for the precise statement).

3 A twist based on the Heisenberg Lie algebra

The example which is of interest to us is based on the Lie algebra $\mathfrak{h}_3(\mathbb{R})$ of the 3-dimensional Heisenberg group, or more precisely on the Hopf algebra $\mathcal{U}(\mathfrak{h}_3(\mathbb{R}))$.

The Lie algebra $\mathfrak{h}_3(\mathbb{R})$ has a basis of three elements \mathfrak{p} , \mathfrak{q} , \mathfrak{t} , with \mathfrak{t} central and

$$[\mathfrak{p}, \mathfrak{q}] = \mathfrak{t}. \tag{3.1}$$

Let us call *parameter space* $\Theta \subset \mathbb{C}[[\nu]]$ the ideal:

$$\Theta := \nu\mathbb{C}[[\nu]].$$

Note that for all $\{a_n\}_{n \geq 0}$ belonging to a complex vector space A and for all $\theta \in \Theta$, the expression $\sum_{n \geq 0} a_n \theta^n$ is a well defined element of $A[[\nu]]$ (for each $N \geq 0$, the coefficient of ν^N is a finite linear combination of elements of A).

There is an action α of \mathbb{Z} by automorphisms on the vector space Θ given by

$$\alpha_k(\theta) = \frac{\theta}{1+k\theta} = \theta - k\theta^2 + \dots, \quad k \in \mathbb{Z}, \quad \theta \in \Theta. \quad (3.2)$$

Clearly α maps Θ into itself, and $\alpha_j\alpha_k = \alpha_{j+k} \forall j, k \in \mathbb{Z}$.

For all $\theta \in \Theta$, a twisting element based on $\mathcal{U}(\mathfrak{h}_3(\mathbb{R}))[[\nu]]$ is

$$F_\theta = \exp \{ \theta \mathfrak{p} \otimes \mathfrak{q} \}.$$

Let A be any $\mathcal{U}(\mathfrak{h}_3(\mathbb{R}))$ -module algebra (associative, since the coassociator is trivial), and for all $n \in \mathbb{Z}$ define

$$A_n := \{ a \in A : \mathfrak{t}.a = na \}. \quad (3.3)$$

Since \mathfrak{t} is central, each A_n is a $\mathcal{U}(\mathfrak{h}_3(\mathbb{R}))$ -module itself, and $A_\bullet := \bigoplus_{n \in \mathbb{Z}} A_n$ is a graded (associative) module subalgebra of A .

We let $a *_\theta b := m \circ F_\theta^{-1}(a \otimes b)$ be the product on $A[[\nu]]$ defined in (2.5), and denote by $A^\theta = (A[[\nu]], *_\theta)$ the deformed algebra. Since \mathfrak{t} is central, each set $A_n[[\nu]]$ is stable under the action of $\mathfrak{p}, \mathfrak{q}$. So $A_\bullet^\theta = (A_\bullet[[\nu]], *_\theta)$ is a \mathbb{Z} -graded subalgebra of A^θ , and $A_0^\theta = (A_0[[\nu]], *_\theta)$ is a yet smaller subalgebra.

Lemma 3.1. *For all $\theta, \theta' \in \Theta$:*

$$(\Delta \otimes \text{id})(F_\theta^{-1})(F_{\theta'}^{-1} \otimes 1) = (\text{id} \otimes \Delta)(F_{\theta'}^{-1})(1 \otimes F_\theta^{-1})\Phi_{\theta, \theta'}, \quad (3.4)$$

where

$$\Phi_{\theta, \theta'} := \exp \{ -\mathfrak{p} \otimes (\theta - \theta' - \theta\theta'\mathfrak{t}) \otimes \mathfrak{q} \}. \quad (3.5)$$

Proof. From (3.4), recalling that $\mathfrak{p}, \mathfrak{q}$ are primitive elements

$$\begin{aligned} \Phi_{\theta, \theta'} &= (1 \otimes F_\theta)(\text{id} \otimes \Delta)(F_{\theta'})(\Delta \otimes \text{id})(F_\theta^{-1})(F_{\theta'}^{-1} \otimes 1) \\ &= e^{\theta 1 \otimes \mathfrak{p} \otimes \mathfrak{q}} e^{\theta'(\mathfrak{p} \otimes 1 \otimes \mathfrak{q} + \mathfrak{p} \otimes \mathfrak{q} \otimes 1)} e^{-\theta(\mathfrak{p} \otimes 1 \otimes \mathfrak{q} + 1 \otimes \mathfrak{p} \otimes \mathfrak{q})} e^{-\theta' \mathfrak{p} \otimes \mathfrak{q} \otimes 1}. \end{aligned} \quad (3.6)$$

Next, we use Baker–Campbell–Hausdorff formula, which for two elements X, Y of an associative algebra with central commutator $[X, Y]$ reduces to $e^{\nu X} e^{\nu Y} = e^{\nu^2[X, Y]} e^{\nu Y} e^{\nu X}$. We get

$$e^{-\theta(\mathfrak{p} \otimes 1 \otimes \mathfrak{q} + 1 \otimes \mathfrak{p} \otimes \mathfrak{q})} e^{-\theta' \mathfrak{p} \otimes \mathfrak{q} \otimes 1} = e^{\theta\theta' \mathfrak{p} \otimes \mathfrak{t} \otimes \mathfrak{q}} e^{-\theta' \mathfrak{p} \otimes \mathfrak{q} \otimes 1} e^{-\theta(\mathfrak{p} \otimes 1 \otimes \mathfrak{q} + 1 \otimes \mathfrak{p} \otimes \mathfrak{q})}.$$

Using this in (3.6) we get

$$\begin{aligned} \Phi_{\theta, \theta'} &= e^{\theta\theta' \mathfrak{p} \otimes \mathfrak{t} \otimes \mathfrak{q}} e^{\theta 1 \otimes \mathfrak{p} \otimes \mathfrak{q}} (e^{\theta'(\mathfrak{p} \otimes 1 \otimes \mathfrak{q} + \mathfrak{p} \otimes \mathfrak{q} \otimes 1)} e^{-\theta' \mathfrak{p} \otimes \mathfrak{q} \otimes 1}) e^{-\theta(\mathfrak{p} \otimes 1 \otimes \mathfrak{q} + 1 \otimes \mathfrak{p} \otimes \mathfrak{q})} \\ &= e^{\theta\theta' \mathfrak{p} \otimes \mathfrak{t} \otimes \mathfrak{q}} e^{\theta 1 \otimes \mathfrak{p} \otimes \mathfrak{q}} e^{\theta' \mathfrak{p} \otimes 1 \otimes \mathfrak{q}} e^{-\theta(\mathfrak{p} \otimes 1 \otimes \mathfrak{q} + 1 \otimes \mathfrak{p} \otimes \mathfrak{q})}. \end{aligned}$$

Now all the exponents mutually commute, and after some simplification we arrive at (3.5). \blacksquare

In particular, we can observe that the coassociator:

$$\Phi_{\theta, \theta} = e^{\theta^2 \mathfrak{p} \otimes \mathfrak{t} \otimes \mathfrak{q}} \quad (3.7)$$

of the QUE algebra is not trivial nor invariant, so that we are dealing with a genuine quasi-Hopf deformation of $\mathcal{U}(\mathfrak{h}_3(\mathbb{R}))$. Many properties of A^θ can be deduced from Lemma 3.1.

Proposition 3.2. *If $\theta' = \alpha_n(\theta)$, with α as in (3.2), one has the generalized associativity law:*

$$(a *_{\theta'} b) *_{\theta} c = a *_{\theta'} (b *_{\theta} c) \quad \forall a, c \in A[[\nu]], \quad b \in A_n[[\nu]]. \quad (3.8)$$

Proof. If $\theta' = \alpha_n(\theta)$, from the definition of the star product and the observation that $\mathfrak{t}.b = nb$, we deduce that $\Phi_{\theta, \theta'}$ is the identity on $a \otimes b \otimes c$ and then (3.8) holds. \blacksquare

Corollary 3.3.

- 1) A_0^{θ} is an associative unital subalgebra of A^{θ} .
- 2) $E_n := A_n[[\nu]]$ is an $A_0^{\alpha_n(\theta)}$ - A_0^{θ} -bimodule (with left module structure given by $*_{\alpha_n(\theta)}$ and right module structure given by $*_{\theta}$).
- 3) $*_{\theta}: E_j \otimes E_k \rightarrow E_{j+k}$ descends to a map $E_j \otimes_{A_0^{\theta}} E_k \rightarrow E_{j+k}$ (where the left and right module structure are given by the $*_{\theta}$ multiplication).
- 4) For all $m, n, p \in \mathbb{Z}$ and $\theta' = \alpha_n(\theta)$, the following diagram commutes:

$$\begin{array}{ccc}
 & E_m \otimes_{A_0^{\theta'}} E_n \otimes_{A_0^{\theta}} E_p & \\
 \text{id} \otimes *_{\theta} \swarrow & & \searrow *_{\theta'} \otimes \text{id} \\
 E_m \otimes_{A_0^{\theta'}} E_{n+p} & & E_{m+n} \otimes_{A_0^{\theta}} E_p \\
 *_{\theta'} \swarrow & & \searrow *_{\theta} \\
 & E_{m+n+p} &
 \end{array}$$

4 Line bundles on the noncommutative torus

Let $H_3(\mathbb{R})$ be the group of upper triangular matrices

$$(x, y, t) := \begin{bmatrix} 1 & x & t \\ 0 & 1 & y \\ 0 & 0 & 1 \end{bmatrix}, \quad x, y, t \in \mathbb{R}.$$

With a slight abuse of notations, we identify elements of $\mathcal{U}(\mathfrak{h}_3(\mathbb{R}))$ with left invariant vector fields on $H_3(\mathbb{R})$, that are generated by the differential operators

$$\mathfrak{p} := \frac{1}{\sqrt{2\pi}} \frac{\partial}{\partial x}, \quad \mathfrak{q} := \frac{-i}{\sqrt{2\pi}} \left(\frac{\partial}{\partial y} + x \frac{\partial}{\partial t} \right), \quad \mathfrak{t} := \frac{1}{2\pi i} \frac{\partial}{\partial t}.$$

The choice of normalization will be clear later on (one can check that (3.1) is satisfied).

Let $H_3(\mathbb{Z}) := \{(x, y, t) \in H_3(\mathbb{R}) : x, y, t \in \mathbb{Z}\}$. By left invariance, the above vector fields descend to the 3-dimensional Heisenberg manifold $M_3 := H_3(\mathbb{Z}) \backslash H_3(\mathbb{R})$. Thinking of functions on $K \backslash G$ as left K -invariant functions on G , we get

$$C^{\infty}(M_3) = \{f \in C^{\infty}(\mathbb{R} \times \mathbb{T}^2) : f(x+1, y, t+y) = f(x, y, t)\},$$

where $\mathbb{T}^2 = \mathbb{R}^2/\mathbb{Z}^2$ (so $f \in C^{\infty}(M_3)$ is periodic with period 1 in y and t). The action of central elements $(0, 0, t) \in H_3(\mathbb{R})$ descends to a principal action of $U(1)$ on M_3 , and $U(1) \backslash M_3 \simeq \mathbb{T}^2$. We identify $C^{\infty}(\mathbb{T}^2)$ with the subset of $f \in C^{\infty}(M_3)$ that do not depend on t (and so are periodic in both x and y):

$$C^{\infty}(\mathbb{T}^2) = \{f \in C^{\infty}(M_3) : f(x, y, t+t') = f(x, y, t) \forall x, y, t, t'\}.$$

Let $A := C^\infty(M_3)$. In the notation of previous section, $A_0 = C^\infty(\mathbb{T}^2)$ and, for $n \neq 0$, every element $f \in A_n$ can be written in the form

$$f(x, y, t) = \sum_{k \in \mathbb{Z}} \tilde{f}(x + \frac{k}{n}; k) e^{2\pi i(ky + nt)} \quad (4.1)$$

for a unique Schwartz function $\tilde{f}: \mathbb{R} \times \mathbb{Z}/n\mathbb{Z} \rightarrow \mathbb{C}$. The bijection $A_n \rightarrow \mathcal{S}(\mathbb{R} \times \mathbb{Z}/n\mathbb{Z})$, $f \mapsto \tilde{f}$, is known as Weil–Brezin–Zak transform [12, Section 1.10]. Functions (4.1) can be interpreted as smooth sections of a non-trivial smooth line bundle on \mathbb{T}^2 [8]. The algebra $A_\bullet = \bigoplus_{n \in \mathbb{Z}} A_n$ is dense in A (in the uniform topology); indeed, by periodicity in t , the only weight spaces of \mathfrak{t} appearing in the decomposition of A are those with integer weight. It is *strongly* \mathbb{Z} -graded, which is the algebraic counterpart of the principality of the bundle $M_3 \rightarrow \mathbb{T}^2$ (see, e.g., [1] or [7]).

By point (1) of Corollary 3.3, A_0^θ is an associative subalgebra of A^θ . By standard Fourier analysis, it is not difficult to verify that A_0^θ is generated by two unitary elements, the functions

$$V(x, y, t) := e^{2\pi i x}, \quad U(x, y, t) := e^{2\pi i y}, \quad (4.2)$$

with relation

$$U *_\theta V = e^{2\pi i \theta} V *_\theta U. \quad (4.3)$$

It is then the formal analogue of the smooth algebra of the noncommutative torus.

We can extend $\mathbb{C}[[\nu]]$ -linearly the map $f \mapsto \tilde{f}$ in (4.1) to a bijective map

$$E_n := A_n[[\nu]] \rightarrow \tilde{E}_n := \mathcal{S}(\mathbb{R} \times \mathbb{Z}/n\mathbb{Z})[[\nu]].$$

For $\theta, \theta' \in \Theta$, we can define a left action of $A_0^{\theta'}$ on \tilde{E}_n and a right action of A_0^θ on \tilde{E}_n by

$$a \cdot \tilde{f} := \widetilde{a *_\theta' f}, \quad \tilde{f} \cdot a := \widetilde{f *_\theta a},$$

for all $\tilde{f} \in \tilde{E}_n$ and $a \in A_0$. A computation using (4.1) gives the explicit formulas (for $n \neq 0$):

$$(U \cdot \tilde{f})(x; k) = \tilde{f}(x - \frac{1}{n}; k - 1), \quad (\tilde{f} \cdot U)(x; k) = \tilde{f}(x - \frac{1}{n} - \theta; k - 1), \quad (4.4a)$$

$$(V \cdot \tilde{f})(x; k) = e^{2\pi i(x - \frac{k}{n})} e^{-2\pi i n \theta' x} \tilde{f}(x; k), \quad (\tilde{f} \cdot V)(x; k) = e^{2\pi i(x - \frac{k}{n})} \tilde{f}(x; k), \quad (4.4b)$$

where for a smooth function ψ , by $\psi(y + \theta')$ we mean the formal power series

$$\psi(y + \theta') := e^{\theta' \partial_y} \psi(y) = \sum_{k \geq 0} \frac{\theta'^k}{k!} \partial_y^k \psi(y).$$

(If we replace θ' by a real number, although the above series is convergent only for ψ analytic, the formulas (4.4) are well defined for any Schwartz function \tilde{f} .)

By comparing these formulas with equations (2.1)–(2.5) of [18] we recognize the formal analogue of Connes–Rieffel imprimitivity bimodule $E_g(\theta)$, in the special case

$$g = \begin{pmatrix} 1 & 0 \\ n & 1 \end{pmatrix}. \quad (4.5)$$

Finally, for all $\tilde{f}_1 \in \tilde{E}_{n_1}$ and $\tilde{f}_2 \in \tilde{E}_{n_2}$, one can compute $\widetilde{f_1 *_\theta f_2} \in \tilde{E}_{n_1+n_2}$ using (4.1) and find the explicit formula

$$\widetilde{f_1 *_\theta f_2}(x; k) = \sum_{k_1+k_2=k} \tilde{f}_1\left(\left(1 - n_2\theta\right)x + \left(1 + n_1\theta\right)\frac{k_1 n_2 - k_2 n_1}{n_1(n_1+n_2)}; k_1\right) \tilde{f}_2\left(x - \frac{k_1 n_2 - k_2 n_1}{n_2(n_1+n_2)}; k_2\right), \quad (4.6)$$

which is valid for $n_1, n_2, n_1 + n_2 \neq 0$. If, on the other hand, $f_1 \in E_n$ and $f_2 \in E_{-n}$ ($n \neq 0$), using the identity

$$\psi(x) = \sum_{m \in \mathbb{Z}} e^{2\pi i m x} \int_0^1 e^{-2\pi i m q} \psi(q) dq,$$

valid for $\psi \in \mathcal{S}(\mathbb{R})$, we get

$$f_1 *_{\theta} f_2 = \sum_{r, s \in \mathbb{Z}} U^r V^s (\tilde{f}_1 | \tilde{f}_2 \cdot V^{-s} U^{-r}), \quad (4.7)$$

where

$$(\tilde{f}_1 | \tilde{f}_2) := \sum_{k=1}^n \int_{-\infty}^{\infty} \tilde{f}_1((1+n\theta)x; k) \tilde{f}_2(x; -k) dx.$$

Equations (4.6) and (4.7) are a formal version of [19, Proposition 1.2(a)].

We can now give an interpretation to Corollary 3.3: point (2) is the analogue of, e.g., equation (2.6) of [18], stating that the algebra of endomorphisms of a finitely generated projective module over a noncommutative torus is another noncommutative torus with a different deformation parameter; point (4) is the associativity of the pairing of bimodules in [19, Proposition 1.2(b)].

5 Complex structures and theta functions

In order to include a complex structure in the construction, it is convenient to start from a different realization of the Heisenberg group and of the principal bundle of previous section.

Let $\tau \in \mathbb{C}$ be a complex number with imaginary part $\Im(\tau) > 0$ and $\Lambda := \mathbb{Z} + \tau\mathbb{Z} \subset \mathbb{C}$ a lattice. We now construct a principal $U(1)$ bundle over the elliptic curve $E_{\tau} := \mathbb{C}/\Lambda$, isomorphic to \mathbb{T}^2 if one forgets about the complex structure.

We parametrize the Heisenberg group as follows: we set $H_3(\mathbb{R}) := \mathbb{C} \times \mathbb{R}$ with multiplication

$$(z_1, t_1) \cdot (z_2, t_2) = (z_1 + z_2, t_1 + t_2 + \Im(\bar{z}_1 z_2)), \quad (5.1)$$

where the bar denotes complex conjugation and \Im the imaginary part. Right invariant vector fields are spanned by

$$\mathfrak{p} = \sqrt{\frac{\Im(\tau)}{\pi}} \left(\frac{\partial}{\partial z} + \frac{i\bar{z}}{2} \frac{\partial}{\partial t} \right), \quad \mathfrak{q} = \sqrt{\frac{\Im(\tau)}{\pi}} \left(\frac{\partial}{\partial \bar{z}} - \frac{iz}{2} \frac{\partial}{\partial t} \right), \quad \mathfrak{t} = \frac{\Im(\tau)}{\pi i} \frac{\partial}{\partial t}, \quad (5.2)$$

where we identify elements of $\mathcal{U}(\mathfrak{h}_3(\mathbb{R}))$ with their representation as differential operators.

Let $H_3(\mathbb{Z}) \subset H_3(\mathbb{R})$ be the subgroup generated by the elements $(1, 0)$ and $(\tau, 0)$.

Remark 5.1. For a more explicit description, one can verify that $H_3(\mathbb{Z})$ coincides with the group $G := \{g_{m,n,k} := (m + n\tau, (mn + 2k)\Im(\tau)) : m, n, k \in \mathbb{Z}\}$. Indeed, from (5.1) we get $g_{m,n,k} g_{m',n',k'} = g_{m+m',n+n',k+k'-nm'}$, proving that G is a group. Clearly $H_3(\mathbb{Z}) \subset G$, since G contains the generators $(1, 0)$ and $(\tau, 0)$ of $H_3(\mathbb{Z})$. On the other hand, $g_{m,n,0} = (1, 0)^m (\tau, 0)^n$ and $g_{0,0,1} = (1, 0) \cdot (\tau, 0) \cdot (-1, 0) \cdot (-\tau, 0)$ belong to $H_3(\mathbb{Z})$, and since $g_{m,n,k} = g_{m,n,0} (g_{0,0,1})^k$, this proves that $H_3(\mathbb{Z}) \supset G$. Note that $H_3(\mathbb{Z})$ is a proper subgroup of $\Lambda \times \Im(\tau)\mathbb{Z}$, since for example it doesn't contain the element $(0, \Im(\tau))$, and contains properly the group $2(\Lambda \times \Im(\tau)\mathbb{Z})$.

We define $M_3 := H_3(\mathbb{R})/H_3(\mathbb{Z})$. As in previous section, the action of central elements $(0, t) \in H_3(\mathbb{R})$ descends to a principal action of $U(1)$ on M_3 , and $M_3/U(1) \simeq E_\tau$. In the notations of Section 2.3, we set $A := C^\infty(M_3)$ and A_n as in (3.3), so that $A_0 \simeq C^\infty(E_\tau)$. The right invariance of $f \in A$ under the action of $(0, 2\Im(\tau)) \in H_3(\mathbb{Z})$ proves that f is periodic in t with period $2\Im(\tau)$, hence the differential operator \mathfrak{t} in (5.2) has integer spectrum (this explains the choice of normalization) and the subalgebra $A_\bullet = \bigoplus_{n \in \mathbb{Z}} A_n$ is dense in A .

The map (4.1) is replaced by the bijection (5.3) below.

Proposition 5.2. *For all $n \neq 0$ there is a bijection $A_n \ni f \mapsto \tilde{f} \in \mathcal{S}(\mathbb{R} \times \mathbb{Z}/n\mathbb{Z})$ given by*

$$f(z, t) = e^{\pi i n \left\{ \frac{1}{\Im(\tau)} t + xy \right\}} \sum_{k \in \mathbb{Z}} e^{2\pi i k x} \tilde{f}\left(y + \frac{k}{n}; k\right), \quad (5.3)$$

where the real coordinates $(x, y) \in \mathbb{R}^2$ are defined by $z := x + \tau y$.

Proof. This is essentially (4.1) modulo a reparametrization. $f \in A$ iff it is right invariant under the action of the two generators $(1, 0)$ and $(\tau, 0)$ of $H_3(\mathbb{Z})$. In real coordinates, we get the conditions $f(x+1, y, t - y\Im(\tau)) = f(x, y, t)$ and $f(x, y+1, t + x\Im(\tau)) = f(x, y, t)$. As mentioned above, f is also $2\Im(\tau)$ -periodic in t , and the condition $\mathfrak{t}f = nf$ says that every $f \in A_n$ is given by $e^{\pi i n \frac{1}{\Im(\tau)} t}$ times a function of x, y . Define $F(x, y)$ implicitly by $f(x, y, z) = e^{\pi i n \left\{ \frac{1}{\Im(\tau)} t + xy \right\}} F(x, y)$. The two invariance conditions above become $F(x+1, y) = F(x, y)$ and $F(x, y+1) = F(x, y)e^{-2\pi i n x}$. From the former, $F(x, y) = \sum_{k \in \mathbb{Z}} e^{2\pi i k x} F_k(y)$ for some functions F_k .

The latter condition gives $F_k(y+1) = F_{k+n}(y)$; if we define $\tilde{f}(y; k) := F_k(y - \frac{k}{n})$, the condition becomes $\tilde{f}(y; k+n) = \tilde{f}(y; k)$. Thus, $\tilde{f}(y; k)$ is periodic in k with period n . Finally, from [8, Lemma 3.2] it follows that f is C^∞ iff \tilde{f} is Schwartz. ■

We can apply the same recipe of previous section, and deform the algebra A with the twist $F_\theta = \exp\{\theta \mathfrak{p} \otimes \mathfrak{q}\}$, $\theta \in \Theta$. The advantage is that now A^θ has, besides A_0^θ , two additional associative subalgebras. The coassociator $\Phi_{\theta, \theta} = e^{\theta^2 \mathfrak{p} \otimes \mathfrak{t} \otimes \mathfrak{q}}$, cf. (3.7), is 1 on the kernels of \mathfrak{p} , \mathfrak{t} and \mathfrak{q} . The second kernel is A_0^θ , the first and third are related by a conjugation $z \mapsto \bar{z}$. We will focus on the latter.

As in previous section, A_0^θ is generated by two unitary functions

$$U(x, y, t) := e^{2\pi i x}, \quad V(x, y, t) := e^{2\pi i y}. \quad (5.4)$$

Since $U *_\theta V = e^{-\frac{\pi \theta \bar{\tau}}{\Im(\tau)}} UV$ and $V *_\theta U = e^{-\frac{\pi \theta \tau}{\Im(\tau)}} VU$, we get the usual noncommutative torus commutation relation¹

$$U *_\theta V = e^{2\pi i \theta} V *_\theta U. \quad (5.5)$$

Let $A_{\text{hol}} := \ker \mathfrak{q}$ and $A_{\text{hol}}^\theta := (A_{\text{hol}}[[\nu]], *_\theta)$. Since F_θ is 1 on $A_{\text{hol}} \otimes A_{\text{hol}}$, the product is undeformed: $a *_\theta b = ab \forall a, b \in A_{\text{hol}}^\theta$ and A_{hol}^θ is a commutative associative subalgebra of A^θ .

With an explicit computation we now check that elements of A_{hol} are (essentially) classical theta functions on the torus. Clearly $A_0 \cap A_{\text{hol}} = \mathbb{C}$ is the set of constant functions. For $n \neq 0$, the set $A_n \cap A_{\text{hol}}$ is described by the following lemma.

Lemma 5.3. *If $n < 0$, $A_{\text{hol}} \cap A_n = \{0\}$. If $n > 0$, elements $f \in A_{\text{hol}} \cap A_n$ are in bijection with elements $c_f \in \mathbb{C}^{\mathbb{Z}/n\mathbb{Z}}$ via the formula*

$$f(z, t) = e^{\pi i n \left\{ \frac{1}{\Im(\tau)} t + xy + \tau y^2 \right\}} \sum_{k \in \mathbb{Z}} c_f(k) q^{k^2/n} e^{2\pi i k z}, \quad (5.6)$$

where $q := e^{\pi i \tau}$.

¹The exchange $U \leftrightarrow V$ in (4.2) and (5.4) is needed to get the same commutation relation (4.3) and (5.5).

Proof. With some algebraic manipulation we can rewrite (5.3) as follows

$$f(z, t) = e^{\pi i n \left\{ \frac{1}{\Im(\tau)} t + xy + \tau y^2 \right\}} \sum_{k \in \mathbb{Z}} q^{k^2/n} e^{2\pi i k z} \left(e^{-\pi i n \tau \left(y + \frac{k}{n} \right)^2} \tilde{f} \left(y + \frac{k}{n}; k \right) \right). \quad (5.7)$$

In real coordinates, \mathfrak{q} is proportional to the differential operator $\nabla := \tau \frac{\partial}{\partial x} - \frac{\partial}{\partial y} + \Im(\tau) z \frac{\partial}{\partial t}$. Both the factor $e^{\pi i n \left\{ \frac{1}{\Im(\tau)} t + xy + \tau y^2 \right\}}$ and the holomorphic exponential $e^{2\pi i k z}$ are in the kernel of such an operator, and since the product in parenthesis in (5.7) only depends on y :

$$\nabla f(z, t) = -e^{\pi i n \left\{ \frac{1}{\Im(\tau)} t + xy + \tau y^2 \right\}} \sum_{k \in \mathbb{Z}} q^{k^2/n} e^{2\pi i k z} \frac{\partial}{\partial y} \left(e^{-\pi i n \tau \left(y + \frac{k}{n} \right)^2} \tilde{f} \left(y + \frac{k}{n}; k \right) \right).$$

One has $\nabla f = 0$ iff $e^{-\pi i n \tau \left(y + \frac{k}{n} \right)^2} \tilde{f} \left(y + \frac{k}{n}; k \right) =: c_f(k)$ does not depend on $y + \frac{k}{n}$, and in this case (5.7) reduces to (5.6). Since $\Im(\tau) > 0$, for $n > 0$ the function $\tilde{f}(y; k) = c_f(k) e^{\pi i n \tau y^2}$ is of Schwartz class for any $c_f \in \mathbb{C}^{\mathbb{Z}/n\mathbb{Z}}$, while for $n < 0$ it is of Schwartz class only if it is zero. ■

The series in (5.6) are the usual theta functions on E_τ . If $n = 1$, for example, the series in (5.6) is proportional to the Jacobi's theta function $\vartheta(z; q) = \sum_{k \in \mathbb{Z}} q^{k^2/n} e^{2\pi i k z}$.

The algebraic structure of theta functions is encoded in the formula (5.8) below.

Proposition 5.4. *Let $n_1, n_2 > 0$. For all $f_i \in A_{\text{hol}} \cap A_{n_i}$, $i = 1, 2$,*

$$c_{f_1 f_2}(k) = \sum_{\substack{k_1, k_2 \in \mathbb{Z} \\ k_1 + k_2 = k}} c_{f_1}(k_1) c_{f_2}(k_2) q^{\frac{1}{n_1 n_2 (n_1 + n_2)} (k_1 n_2 - k_2 n_1)^2}, \quad (5.8)$$

where $f \mapsto c_f$ is the map in (5.6).

Proof. A simple computation using (5.6) and the algebraic identity:

$$\frac{k_1^2}{n_1} + \frac{k_2^2}{n_2} = \frac{(k_1 + k_2)^2}{n_1 + n_2} + \frac{(k_1 n_2 - k_2 n_1)^2}{n_1 n_2 (n_1 + n_2)}. \quad \blacksquare$$

Remark 5.5. In the C^* -algebraic setting, for any fixed modular parameter τ , any g as in (1.1) and $\theta \in \mathbb{R}$ solution of $g\theta = \theta$, a ring of “quantum theta functions” $B_g(\theta, \tau)$ can be defined as a suitable “holomorphic” subalgebra of the tensor algebra $\bigoplus_{k \geq 0} E_g(\theta)^{\otimes k}$ (with tensor product over the algebra A_θ of the noncommutative torus and $E_g(\theta)^{\otimes 0} := A_\theta$). A product formula for quantum theta functions appeared first in [9] in terms of generators and relations (see also [19]), while an alternative formula which is closer to our notations is [22, equation (7.4)].

For g as in (4.5) and $n = 1$, as one can easily check, the product $(c_{f_1}, c_{f_2}) \mapsto c_{f_1 f_2}$ defined by (5.8) coincides with [22, equation (7.4)] and $B_g(\theta, \tau) \simeq A_{\text{hol}}$ (while $B_g(\theta, \tau)$ is a subalgebra of A_{hol} if $n > 1$ in (4.5)). However, this is not surprising since, for g as in (4.5), the only solution to $g\theta = \theta$ is $\theta = 0$. The construction in this section, on the other hand, works in a formal setting and there is no constrain on θ . For any $\theta \in \Theta$, A^θ is well-defined, although not associative, and has an associative (and commutative) subalgebra given by classical theta functions.

6 On higher rank vector bundles

In this section, we describe how to derive (formal) imprimitivity bimodules associated to an arbitrary element of $\text{SL}(2, \mathbb{Z})$.

Let us fix a g as in (1.1), with $d \neq 0$. We denote by $\mathcal{M}_{c,d} \subset C^\infty(\mathbb{R} \times \mathbb{T})$ the set of functions satisfying

$$f(x+d, y) = e^{-2\pi icy} f(x, y), \quad (6.1)$$

This is a $C^\infty(\mathbb{T}^2)$ -right module with product given by pointwise multiplication. One can verify that $\mathcal{M}_{c,d}$ is isomorphic to the module of (smooth) sections of a rank $|d|$ vector bundle on \mathbb{T}^2 as follows. Note that $\mathcal{M}_{c,d} \simeq \mathcal{M}_{-c,-d}$, so from now on we can assume that $d \geq 1$.

That $\mathcal{M}_{0,d}$ is a free module of rank d comes from the following observation.

Remark 6.1. By standard Fourier analysis, every smooth function φ with period d can be written (in a unique way) as

$$\varphi(x) = \sum_{k=1}^d e^{2\pi i \frac{k}{d} x} \varphi_k(x)$$

with φ_k of period 1. This gives a $C^\infty(\mathbb{R}/\mathbb{Z})$ -module isomorphism:

$$C^\infty(\mathbb{R}/d\mathbb{Z}) \rightarrow C^\infty(\mathbb{R}/\mathbb{Z}) \otimes \mathbb{C}^d, \quad \varphi \mapsto (\varphi_1, \dots, \varphi_d).$$

(Functions of period d form a free module of rank d over functions of period 1.)

For the reason above, $\mathcal{M}_{0,d} \simeq \mathcal{M}_{0,1} \otimes \mathbb{C}^d$ as modules over $\mathcal{M}_{0,1} = C^\infty(\mathbb{T}^2)$.

From now on we forget about free modules and assume that $c \neq 0$ (and $d \geq 1$ as above). Note that the condition $\det(g) = 1$ in (1.1) guarantees that c and d are coprime. Vice versa, by Bézout's lemma such a g exists for every coprime c, d (although it is not unique).

Lemma 6.2. *Every $n \in \mathbb{Z}$ can be written, in a unique way, as $n = kc + md$ for some $1 \leq k \leq d$ and $m \in \mathbb{Z}$.*

Proof. The map $\mathbb{Z}^2 \ni (k, m) \rightarrow kc + md \in \mathbb{Z}$ is surjective, due to the identity $n = (-nb)c + (na)d$ following from the determinant condition in (1.1); since (k, m) and $(k-d, m+c)$ have the same image, one can always choose $1 \leq k \leq d$; restricted to $[1, \dots, d] \times \mathbb{Z}$ the map is also injective, since $kc + md = k' + md' - \text{i.e., } (k-k')c + (m-m')d = 0$ - implies that d must divide $k-k'$. But $|k-k'| \leq d-1$, so it must be $k-k' = 0$, which also implies $m-m' = 0$. ■

Remark 6.1 can then be rephrase as follows.

Lemma 6.3. *Every $\varphi \in C^\infty(\mathbb{R}/d\mathbb{Z})$ can be written in a unique way as*

$$\varphi(x) = \sum_{k=1}^d e^{2\pi i \frac{c}{d} kx} \varphi_k(x)$$

with $\varphi_1, \dots, \varphi_d \in C^\infty(\mathbb{R}/\mathbb{Z})$.

Proof. From Lemma 6.2, we can write the Fourier series of φ as

$$\varphi(x) = \sum_{n \in \mathbb{Z}} e^{2\pi i \frac{1}{d} nx} \hat{\varphi}_k = \sum_{k=1}^d e^{2\pi i \frac{c}{d} kx} \left(\sum_{m \in \mathbb{Z}} e^{2\pi i mx} \hat{\varphi}_k \right)$$

and called $\varphi_k(x) = \sum_{m \in \mathbb{Z}} e^{2\pi i mx} \hat{\varphi}_k$ we get the desired result. ■

Proposition 6.4. $\mathcal{M}_{c,d}$ is isomorphic to the $C^\infty(\mathbb{T}^2)$ -module of smooth sections of the vector bundle

$$\begin{array}{c} (\mathbb{T} \times \mathbb{R} \times \mathbb{C}^d) / \mathbb{Z} \\ \downarrow \\ \mathbb{T} \times \mathbb{R} / \mathbb{Z} = \mathbb{T}^2, \end{array} \quad (6.2)$$

where the action of \mathbb{Z} on $\mathbb{T} \times \mathbb{R} \times \mathbb{C}^d$ is generated by the map

$$(x, y; v_1, v_2, \dots, v_d) \mapsto (x, y + 1; e^{2\pi i c x} v_d, v_1, \dots, v_{d-1}).$$

Proof. Any $f \in \mathcal{M}_{c,d}$ – since $e^{2\pi i \frac{c}{d} x y} f(x, y)$ is periodic in x with period d – can be written (in a unique way) in the form $f(x, y) = e^{-2\pi i \frac{c}{d} x y} \sum_{k=1}^d e^{2\pi i \frac{c}{d} k x} f_k(x, y)$, where f_1, \dots, f_d have period 1 in x . From (6.1) we get the conditions

$$f_k(x, y + 1) = f_{k-1}(x, y), \quad \forall k \neq 1, \quad f_1(x, y + 1) = e^{2\pi i c x} f_d(x, y).$$

We may then think of (f_1, \dots, f_d) as a section of the vector bundle (6.2). More precisely, the corresponding section is the map

$$\mathbb{T}^2 \ni [x, y] \mapsto [[x], y; f_1(x, y), \dots, f_d(x, y)] \in (\mathbb{T} \times \mathbb{R} \times \mathbb{C}^d) / \mathbb{Z}.$$

The transformation $f \mapsto (f_1, \dots, f_d)$ is the desired $C^\infty(\mathbb{T}^2)$ -module isomorphism. \blacksquare

Let U, V be the generators of $C^\infty(\mathbb{T}^2)$ in (4.2), and U', V' the $C^\infty(\mathbb{T}^2)$ -linear endomorphisms of $\mathcal{M}_{c,d}$ given by

$$(U'f)(x, y) := e^{2\pi i a y} f(x + b, y), \quad (V'f)(x, y) := e^{2\pi i \frac{1}{d} x} f(x, y).$$

(Here a, b are the elements in the first row of (1.1).) The operators U', V' are unitary if we equip $\mathcal{M}_{c,d}$ with the inner product given by $\langle f_1, f_2 \rangle := \int_{[0,d] \times [0,1]} f_1(x, y)^* f_2(x, y) dx dy$, and satisfy the defining relation of the rational noncommutative torus with deformation parameter b/d , namely $U'V' = e^{2\pi i \frac{b}{d}} V'U'$. Let us denote by $B_{b/d}$ the algebra of power series in U', V' with rapid decay coefficients.

In order to apply the deformation machinery of Section 3 we need an action of $\mathfrak{h}_3(\mathbb{R})$. The pointwise product gives a map $\mathcal{M}_{c,d} \times \mathcal{M}_{c',d'} \rightarrow \mathcal{M}_{cd'+c'd, dd'}$. An action of $\mathfrak{h}_3(\mathbb{R})$ by derivations is

$$\mathfrak{p}f = \frac{1}{\sqrt{2\pi}} \frac{\partial}{\partial x} f, \quad \mathfrak{q}f = \frac{-i}{\sqrt{2\pi}} \left(\frac{\partial}{\partial y} + 2\pi i \frac{c}{d} x \right) f, \quad \mathfrak{t}f = \frac{c}{d} f,$$

for all $f \in \mathcal{M}_{c,d}$. The condition that \mathfrak{q} is a derivation (i.e., satisfies the Leibniz rule) fixes the constant in front of the factor x , and consequently the normalization of \mathfrak{t} . In particular $\mathcal{M}_{0,1} = C^\infty(\mathbb{T}^2)$ is in the kernel of \mathfrak{t} .

A compatible action on $B_{b/d} \subset \text{End}_{C^\infty(\mathbb{T}^2)}(\mathcal{M}_{c,d})$ is given by commutators: for all primitive $X \in \mathfrak{h}_3(\mathbb{R})$ and all $\xi \in B_{b/d}$, $X(\xi)$ is the endomorphism defined by $X(\xi)f := X(\xi f) - \xi X(f)$ for all $f \in \mathcal{M}_{c,d}$. With a simple computation one checks that \mathfrak{t} is mapped to 0, while $\mathfrak{p} \mapsto \frac{1}{\sqrt{2\pi}} \delta_{V'}$ and $\mathfrak{q} \mapsto \frac{-i}{\sqrt{2\pi}} \delta_{U'}$, where $\delta_{U'}, \delta_{V'}$ are the derivations defined on generators by

$$\delta_{U'}(U') = \frac{2\pi i}{d} U', \quad \delta_{U'}(V') = 0, \quad \delta_{V'}(U') = 0, \quad \delta_{V'}(V') = \frac{2\pi i}{d} V'.$$

We can now use the twisting element in Section 3 to construct, for any $\theta, \theta' \in \Theta$, two new algebras $A_0^\theta := (C^\infty(\mathbb{T}^2)[[\nu]], *_\theta)$ and $B_{b/d}^{\theta'} := (B_{b/d}[[\nu]], *_{\theta'})$ and to deform $E_{c,d} := \mathcal{M}_{c,d}[[\nu]]$ into a left $B_{b/d}^{\theta'}$ -module and right A_0^θ -module.

Note that A_0^θ and $B_{b/d}^{\theta'}$ are both associative (they are in the kernel of \mathfrak{t} , hence the coassociator is trivial), the former is generated by the two unitaries U and V with relation (4.3), the latter by the two unitaries U' and V' with relation

$$U' *_{\theta'} V' = e^{2\pi i(d^{-2}\theta' + \frac{b}{d})} V' *_{\theta'} U'. \quad (6.3)$$

The action α in (3.2) extends in an obvious way to an action of \mathbb{Q} on Θ , and from Lemma 3.1 we derive the following analogue of Proposition 3.2.

Proposition 6.5. *The left $B_{b/d}^{\theta'}$ action and right A_0^θ action on $E_{c,d}$ commute, i.e.,*

$$(\xi *_{\theta'} f_1) *_{\theta} f_2 = \xi *_{\theta'} (f_1 *_{\theta} f_2) \quad \forall \xi \in B_{b/d}^{\theta'}, \quad f_1 \in E_{c,d}, \quad f_2 \in A_0^\theta, \quad (6.4)$$

if and only if $\theta' = \alpha_{\frac{c}{d}}(\theta)$.

Proof. If $\theta' = \alpha_{\frac{c}{d}}(\theta)$, since $\mathfrak{t}f = \frac{c}{d}f$ for all $f \in E_{c,d}$, we deduce that $\Phi_{\theta,\theta'}$ is the identity on $\xi \otimes f_1 \otimes f_2$ and then (6.4) holds. On the other hand,

$$m(m \otimes \text{id})(\Phi_{\theta,\theta'} - 1)(V' \otimes f \otimes U) = (e^{2\pi i d^{-1}(1 + \frac{c}{d}\theta)\{\alpha_{c/d}(\theta) - \theta'\}} - 1)V' f U$$

is zero for all $f \in E_{c,d}$ only if $\theta' = \alpha_{\frac{c}{d}}(\theta)$. ■

Note that if $\theta' = \alpha_{\frac{c}{d}}(\theta)$, the parameter in (6.3) is

$$d^{-2}\theta' + \frac{b}{d} = \frac{a\theta + b}{c\theta + d} = g\theta$$

as expected.

The action of the two algebras on $E_{c,d}$ can be explicitly computed on generators. The left action of U' and the right action of V are undeformed, while for the remaining two generators one gets

$$\begin{aligned} (f *_{\theta} U)(x, y) &= e^{2\pi i y} f(x - \theta, y), \\ (V' *_{\theta'} f)(x, y) &= e^{2\pi i \frac{1}{d}x} e^{-2\pi i c d^{-2}\theta' x} f(x, y - d^{-1}\theta'). \end{aligned}$$

It remains to compare what we obtained with the well-known formulas that one has in the C^* -algebraic setting, cf., e.g., equations (2.1)–(2.5) of [18]. This can be done by means of the transform $E_{c,d} \rightarrow \tilde{E}_{c,d} := \mathcal{S}(\mathbb{R} \times \mathbb{Z}/c\mathbb{Z})[[\nu]]$, $f \mapsto \tilde{f}$, defined by

$$f(x, y) = \sum_{k \in \mathbb{Z}} \tilde{f}(x + k\frac{d}{c}; k) e^{2\pi i k y}$$

and similar to (4.1) (except for $\frac{1}{n}$ replaced by $\frac{d}{c}$). Under this transform, the left/right module structure become:

$$\begin{aligned} \widetilde{U' *_{\theta'} f}(x; k) &= \tilde{f}(x - \frac{1}{c}; k - a), & \widetilde{f *_{\theta} U}(x; k) &= \tilde{f}(x - \frac{d}{c} - \theta; k - 1), \\ \widetilde{V' *_{\theta'} f}(x; k) &= e^{2\pi i \frac{1}{d}(x - k\frac{d}{c})} e^{-2\pi i c d^{-2}\theta' x} \tilde{f}(x; k), & \widetilde{f *_{\theta} V}(x; k) &= e^{2\pi i(x - k\frac{d}{c})} \tilde{f}(x; k), \end{aligned}$$

for all $\tilde{f} \in \tilde{E}_{c,d}$. These formulas reduce to equations (2.1)–(2.5) of [18] if $\theta' = \alpha_{\frac{c}{d}}(\theta)$ and after replacing formal deformation parameters by irrational numbers.

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