

The Fractional Poisson Process and the Inverse Stable Subordinator

Mark M. Meerschaert* Erkan Nane† P. Vellaisamy‡

Abstract

The fractional Poisson process is a renewal process with Mittag-Leffler waiting times. Its distributions solve a time-fractional analogue of the Kolmogorov forward equation for a Poisson process. This paper shows that a traditional Poisson process, with the time variable replaced by an independent inverse stable subordinator, is also a fractional Poisson process. This result unifies the two main approaches in the stochastic theory of time-fractional diffusion equations. The equivalence extends to a broad class of renewal processes that include models for tempered fractional diffusion, and distributed-order (e.g., ultraslow) fractional diffusion. The paper also discusses the relation between the fractional Poisson process and Brownian time.

Key words: Fractional Poisson process; Inverse stable subordinator; Renewal process; Mittag-Leffler waiting time; Fractional difference-differential equations; Caputo fractional derivative; Generalized Mittag-leffler function; Continuous time random walk limit; Distributed order derivative; Tempered fractional derivative.

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*Department of Statistics and Probability, Michigan State University, East Lansing, MI 48823. E-mail: mcubed@stt.msu.edu. URL: <http://www.stt.msu.edu/~mcubed>. The research of MMM was partially supported by NSF grants DMS-1025486, DMS-0803360, EAR-0823965, and NIH grant R01-EB012079-01.

†Corresponding author. Address: Department of Mathematics and Statistics, 221 Parker Hall, Auburn University, Auburn, Alabama 36849, USA. E-mail: nane@auburn.edu. URL: <http://www.auburn.edu/~ezn0001>

‡Department of Mathematics, Indian Institute of Technology Bombay, Powai, Mumbai 400076, INDIA. E-mail: pv@math.iitb.ac.in. URL: <http://www.math.iitb.ac.in/~pv>

1 Introduction

The fractional Poisson process (FPP) was introduced and studied by Repin and Saichev [42], Jumarie [25], Laskin [30], Mainardi et al. [32, 33], Uchaikin et al. [46] and Beghin and Orsingher [6, 7]. The FPP is a natural generalization of the usual Poisson process, with an interesting connection to fractional calculus. This renewal process has IID waiting times J_n that satisfy

$$\mathbb{P}(J_n > t) = E_\beta(-\lambda t^\beta) \quad (1.1)$$

for $0 < \beta \leq 1$, where

$$E_\beta(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(1 + \beta k)}. \quad (1.2)$$

denotes the Mittag-Leffler function. When $\beta = 1$, the waiting times are exponential with rate λ , since $e^z = E_1(z)$. Let $T_n = J_1 + \dots + J_n$ be the time of the n th jump. Then the FPP

$$N_\beta(t) = \max\{n \geq 0 : T_n \leq t\} \quad (1.3)$$

is a renewal process with Mittag-Leffler waiting times.

A compound FPP is obtained by subordinating a random walk to the FPP. The resulting process is non-Markovian (unless $\beta = 1$) and the distribution of that process solves a “master equation” analogous to the Kolmogorov equation for Markov processes, with the usual integer order time derivative replaced by a fractional derivative.

The continuous time random walk (CTRW) is another useful model in fractional calculus. Consider a CTRW whose IID particle jumps Y_n have PDF $w(x)$, and whose IID waiting times (J_n) are Mittag-Leffler variables independent of (Y_n) . The particle location after n jumps is $S(n) = Y_1 + \dots + Y_n$, and the CTRW $S(N_\beta(t))$ gives the particle location at time $t \geq 0$. Hilfer and Anton [22] show that the PDF $p(x, t)$ of the CTRW $S(N_\beta(t))$ solves the fractional master equation

$$\partial_t^\beta p(x, t) = -\lambda p(x, t) + \lambda \int_{-\infty}^{\infty} p(x - y, t) w(y) dy \quad (1.4)$$

where ∂_t^β denotes the Caputo fractional derivative. The Caputo fractional derivative, defined for $0 \leq n - 1 < \beta < n$ by

$$\partial_t^\beta g(t) = \frac{1}{\Gamma(n - \beta)} \int_0^t (t - r)^{n-1-\beta} g^{(n)}(r) dr, \quad (1.5)$$

where $g^{(k)}$ denotes the k -th derivative of g , was invented to properly handle initial values [13].

If $\beta = 1$, then ∂_t^β is the usual first derivative. The corresponding CTRW $S(N_1(t))$ is a compound Poisson process, and (1.4) reduces to

$$\partial_t p(x, t) = -\lambda p(x, t) + \lambda \int_{-\infty}^{\infty} p(x - y, t) w(y) dy, \quad (1.6)$$

the Cauchy problem associated with this infinitely divisible Lévy process. Then a general result on Cauchy problems [3, Theorem 3.1] implies that the PDF of the time-changed process $S(N_1(E(t)))$ solves the fractional Cauchy problem (1.4), where

$$E(t) = \inf\{r > 0 : D(r) > t\} \quad (1.7)$$

is the right-continuous inverse (hitting time, first passage time) of $D(t)$, a standard β -stable subordinator with $\mathbb{E}[e^{-sD(t)}] = e^{-ts^\beta}$ for some $0 < \beta < 1$.

Since the PDF of both $S(N_\beta(t))$ and $S(N_1(E(t)))$ solve the same governing equation (1.4), with the same point-source initial condition (i.e., both processes start at the origin), these two processes have the same one dimensional distributions. Heuristically, the degenerate case $Y_n \equiv 1$ gives $S(n) = n$, which strongly suggest that the FPP $N_\beta(t)$ and the process $N_1(E(t))$ have the same one dimensional distributions. We will call $N_1(E(t))$ the fractal time Poisson process (FTPP), since it comes from a self-similar time change (see, e.g., [35, Proposition 3.1]). In this paper, we will prove that the FPP and the FTTP are in fact the same process, by showing that the waiting times between jumps in the FTTP are IID Mittag-Leffler. This strong connection between the FPP and the FTTP unifies the two main approaches in the stochastic theory of fractional diffusion. For example, the FPP approach was used recently in the work of Behgin and Orsingher [6], while the inverse stable subordinator is a key ingredient in [39].

2 Two equivalent formulations

Recall that the fractional Poisson process (FPP) $N_\beta(t)$ is a renewal process with Mittag-Leffler waiting times (1.1), and the fractal time Poisson process (FTTP) $N_1(E(t))$ is Poisson process, with rate $\lambda > 0$, time-changed via the inverse stable subordinator (1.7). The proof that the FPP and the FTTP are the same process requires the following simple lemma.

Lemma 2.1. *Let $D(t)$ be a strictly increasing right-continuous process with left-hand limits, and let $E(t)$ be its right-continuous inverse defined by (1.7). Then*

$$D(r-) = \sup\{t > 0 : E(t) < r\} \tag{2.1}$$

for any $r > 0$.

Proof. Let $t_0 = \sup\{t > 0 : E(t) < r\}$. Then there exists a sequence of points $t_n \uparrow t_0$ such that $E(t_n) < r$ for all n . Let $\varepsilon_n = r - E(t_n) > 0$. If $r > E(t)$ then, since $D(t)$ is strictly increasing, $D(r) > t$. Since $D(r)$ is right-continuous, it follows that $D(E(t)) \geq t$ for all $t > 0$. Then we have $t_n \leq D(E(t_n)) = D(r - \varepsilon_n) < D(r-)$. Letting $n \rightarrow \infty$ shows that $D(r-) \geq t_0$.

Since D has left-hand limits, for any $r_n \uparrow r$ we have $D(r_n) \rightarrow D(r-)$ as $n \rightarrow \infty$. If $D(r-) > t_0$, then for some $r_n < r$ we have $D(r_n) > t_0$. Since $E(t)$ is nondecreasing and continuous, this implies that $E(D(r_n)) \geq r$, by definition of t_0 . But, $E(D(r)) = r$ for all $r > 0$ implying that $r_n \geq r$, which is a contradiction. Thus, (2.1) follows. \square

Theorem 2.2. *For any $0 < \beta < 1$, the FTTP $N_1(E(t))$ is also a FPP. That is, the waiting times between jumps of the FTTP are IID Mittag-Leffler.*

Proof. Let W_n be an IID sequence with $\mathbb{P}(W_n > t) = e^{-\lambda t}$ and $V_n = W_1 + \dots + W_n$ so that the Poisson process $N_1(t) = \max\{n \geq 0 : V_n \leq t\}$. Let

$$\tau_n = \sup\{t > 0 : N_1(E(t)) < n\} \tag{2.2}$$

denote the jump times of the FTTP. This definition of the jump times takes into account the fact that $E(t)$ has constant intervals corresponding to the jumps of the process $D(t)$. Using the fact that $\{N_1(t) < n\} = \{V_n > t\}$ for the Poisson process, along with (2.2), we have

$$\tau_n = \sup\{t > 0 : E(t) < V_n\}.$$

Then Lemma 2.1 implies that $\tau_n = D(V_n-)$. Define $X_1 = \tau_1$ and $X_n = \tau_n - \tau_{n-1}$ for $n \geq 2$, the waiting times between jumps of the FTTP. In order to show that the FTTP is an FPP, it suffices to show that X_n are IID Mittag-Leffler, i.e., they are IID with the same distribution (1.1) as J_n .

Recall that the Laplace transform of the exponential distribution $\mathbb{E}(e^{-sW_n}) = \lambda/(\lambda + s)$. Also recall that $\mathbb{E}(e^{-sD(t)}) = e^{-ts^\beta}$. Since $D(t)$ is a Lévy process, it has no fixed points of discontinuity and hence $D(t-), D(t)$ are identically distributed for all $t \geq 0$. (Indeed, $D(t) = D(t-) a.s.$ [2, Lemma 2.3.2]).

Then a conditioning argument yields

$$\begin{aligned} \mathbb{E}(e^{-s\tau_1}) &= \mathbb{E}(e^{-sD(W_1^-)}) = \mathbb{E}\left[\mathbb{E}\left(e^{-sD(W_1^-)} \middle| W_1\right)\right] \\ &= \mathbb{E}\left[\mathbb{E}\left(e^{-sD(W_1)} \middle| W_1\right)\right] = \mathbb{E}\left[e^{-W_1 s^\beta}\right] = \frac{\lambda}{\lambda + s^\beta}. \end{aligned} \quad (2.3)$$

Let $f_\beta(x) = \partial_x[1 - E_\beta(-\lambda x^\beta)]$ be the Mittag-Leffler PDF of J_n . It is well known that

$$\int_0^\infty e^{-sx} E_\beta(-\lambda x^\beta) dx = \frac{s^{\beta-1}}{\lambda + s^\beta},$$

see for example [39, Eq. (3.4)]. Now integrate by parts to see that

$$\begin{aligned} \mathbb{E}(e^{-sT_1}) &= \int_0^\infty e^{-sx} f_\beta(x) dx \\ &= \int_0^\infty s e^{-sx} (1 - E_\beta(-\lambda x^\beta)) dx \\ &= s \left[\frac{1}{s} - \frac{s^{\beta-1}}{\lambda + s^\beta} \right] = \frac{\lambda}{\lambda + s^\beta} = \mathbb{E}(e^{-s\tau_1}) \end{aligned} \quad (2.4)$$

and then the uniqueness theorem for LT implies that T_1, τ_1 are identically distributed. In particular, X_1 has the same Mittag-Leffler distribution as J_1 .

A straightforward extension of this argument shows that (T_1, \dots, T_n) is identically distributed with (τ_1, \dots, τ_n) for any positive integer n . To ease notation, we only write the case $n = 2$. First observe that

$$\mathbb{E}(e^{-s_1 T_1} e^{-s_2 T_2}) = \mathbb{E}(e^{-s_1 J_1} e^{-s_2 (J_1 + J_2)}) = \frac{\lambda}{\lambda + (s_1 + s_2)^\beta} \cdot \frac{\lambda}{\lambda + s_2^\beta},$$

using the independence of J_1 and J_2 . Next write

$$\begin{aligned} \mathbb{E}(e^{-s_1 D(t_1)} e^{-s_2 D(t_1 + t_2)}) &= \mathbb{E}(e^{-s_1 D(t_1)} e^{-s_2 [D(t_1) + D(t_1 + t_2) - D(t_1)]}) \\ &= \mathbb{E}(e^{-(s_1 + s_2) D(t_1)} e^{-s_2 [D(t_1 + t_2) - D(t_1)]}) \\ &= e^{-t_1 (s_1 + s_2)^\beta} e^{-t_2 s_2^\beta}, \end{aligned}$$

using the fact that $D(t)$ has independent increments. Then

$$\begin{aligned} \mathbb{E}(e^{-s_1\tau_1-s_2\tau_2}) &= \mathbb{E}(e^{-s_1D(W_1)-s_2D([W_1+W_2]-)}) \\ &= \mathbb{E}\left[\mathbb{E}\left(e^{-s_1D(W_1)-s_2D(W_1+W_2)}\middle|W_1, W_2\right)\right] \\ &= \mathbb{E}\left[e^{-W_1(s_1+s_2)^\beta} e^{-W_2s_2^\beta}\right] = \frac{\lambda}{\lambda + (s_1 + s_2)^\beta} \cdot \frac{\lambda}{\lambda + s_2^\beta} = \mathbb{E}(e^{-s_1T_1}e^{-s_2T_2}). \end{aligned}$$

Now an application of the continuous mapping theorem shows that (J_1, \dots, J_n) is identically distributed with (X_1, \dots, X_n) for any positive integer n . Then (X_n) is an IID sequence, so $N_1(E(t))$ is a renewal process. \square

Remark 2.3. Theorem 2.2 extends a result in Behgin and Orsingher [6]. They define (in our notation) a random variable $E(t)$ and show that the two random variables $N_\beta(t)$ and $N_1(E(t))$ have the same density function, by comparing their Laplace transforms. They identify $E(t)$ only through its density function, which they express in terms of an integral involving the density of $D(t)$, see Remark 3.2 for more detail. Cahoy, Uchaikin, Woyczynski [12] also connect the Mittag-Leffler distribution with a stable law. They note that (in our notation) $P(J_n > t) = E[\exp(-\lambda t^\beta/D(1)^\beta)]$, which is useful in simulations. To connect this with our work, note that $E(t) = (t/D(1))^\beta$ in distribution (see Corollary 3.1 in [35]), so that $P(J_n > t) = E[\exp(-\lambda E(t))]$. A result of Bingham [9] shows that the Laplace transform of the stable hitting time $E(t)$ is Mittag-Leffler, so that (1.1) holds.

Remark 2.4. The proof of Theorem 2.2 uses the fact that, if $D(t)$ is a β -stable subordinator and W_1 is exponential, then $D(W_1)$ has a Mittag-Leffler distribution. This fact was first noticed by Pillai [41], who showed that $W_1^{1/\beta}D(1)$ is Mittag-Leffler. These are equivalent because $D(t)$ is identically distributed with $t^{1/\beta}D(1)$. This Mittag-Leffler distribution is also known as the positive Linnik law, e.g., see Huillet [24]. It has the property of geometric stability: A geometric random sum of Mittag-Leffler variables is again Mittag-Leffler, e.g., see Kozubowski [28].

Next we want to show that the FTTP $N_1(E(t))$, and hence also the FPP $N_\beta(t)$, occurs naturally as a CTRW scaling limit. This provides a further justification for the FPP as a robust physical model, see for example Laskin [30]. Suppose now that $\mathbb{P}(J_n > t) = t^{-\beta}L(t)$, where $0 < \beta < 1$ and L is slowly varying. For example, this is true of the Mittag-Leffler waiting times. Then J_1 belongs to the strict domain of attraction of some stable law D with index $0 < \beta < 1$, i.e., there exist $b_n > 0$ such that

$$b_n(J_1 + \dots + J_n) \Rightarrow D, \tag{2.5}$$

where $D(1) = D > 0$ almost surely, and \Rightarrow denotes convergence in distribution. Let $b(t) = b_{[t]}$. Then $b(t) = t^{-1/\beta}L_0(t)$ for some slowly varying function $L_0(t)$ (e.g., see [19, XVII.5]). Since b varies regularly with index $-1/\beta$, b^{-1} is regularly varying with index $1/\beta > 0$ and so by [45, Property 1.5.5] there exists a regularly varying function \tilde{b} with index β such that $1/b(\tilde{b}(c)) \sim c$, as $c \rightarrow \infty$. Here we use the notation $f \sim g$ for positive functions f, g if and only if $f(c)/g(c) \rightarrow 1$ as $c \rightarrow \infty$. Let $T_n = J_1 + \dots + J_n$ and define a renewal process

$$R(t) = \max\{n \geq 0 : T_n \leq t\} \tag{2.6}$$

with these waiting times. Next, construct a CTRW with iid Bernoulli jumps $Y_n^{(p)}$ with $\mathbb{P}(Y_n^{(p)} = 1) = p$ and $\mathbb{P}(Y_n^{(p)} = 0) = 1 - p$, independent of (J_n) . Let $S^{(p)}(n) = Y_1^{(p)} + \dots + Y_n^{(p)}$, a binomial random variable. Then $S^{(p)}(R(t))$ is a CTRW with heavy tailed waiting times and Bernoulli jumps.

Theorem 2.5. *The FTTP is the process limit of a CTRW sequence:*

$$\{S^{(1/\tilde{b}(c))}([\lambda R(ct)])\}_{t \geq 0} \Rightarrow \{N_1(E(t))\}_{t \geq 0} \quad (2.7)$$

as $c \rightarrow \infty$ in the M_1 topology on $D([0, \infty), \mathbb{R})$.

Proof. Since the sequence (J_n) is in the strict domain of attraction of a β -stable random variable D , [35, Corollary 3.4] shows that

$$\{\tilde{b}(c)^{-1}R(ct)\}_{t \geq 0} \Rightarrow \{E(t)\}_{t \geq 0} \quad \text{as } c \rightarrow \infty.$$

in the Skorokhod J_1 topology, where $D(t)$ is the stable subordinator with $D(1) = D$, and $E(t)$ is given by (1.7).

Since the binomial random variable $S^{(p)}(n)$ has LT $\mathbb{E}(e^{-sS^{(p)}(n)}) = (1 + (e^{-s} - 1)p)^n$ for any $n \geq 0$, it follows that

$$\mathbb{E}(e^{-sS^{(p)}([\lambda t/p])}) = (1 + (e^{-s} - 1)p)^{[\lambda t/p]} \rightarrow \exp(-\lambda t(1 - e^{-s})),$$

as $p \rightarrow 0$, using the fact that $(1 + ap)^{1/p} \rightarrow e^a$ as $p \rightarrow 0$. It follows by the continuity theorem for LT that $S^{(p)}([\lambda t/p]) \Rightarrow N_1(t)$ for any $t > 0$, since $\exp(-\lambda t(1 - e^{-s}))$ is the LT of the Poisson random variable $N_1(t)$. Then a standard argument (e.g., see [34, Example 11.2.18] shows that we also get

$$\{S^{(p)}([\lambda t/p])\}_{t \geq 0} \xrightarrow{f.d.} \{N_1(t)\}_{t \geq 0},$$

as $p \rightarrow 0$, where $\xrightarrow{f.d.}$ denotes convergence of all finite dimensional distributions. Since the sample paths of $S^{(p)}([\lambda t/p])$ are increasing and $N_1(t)$ is continuous in probability, being a Lévy process, J_1 convergence follows using [9, Theorem 3].

Since the CTRW waiting times (J_n) are independent of the jumps $(Y_n^{(p)})$, and since $1/\tilde{b}(c) \rightarrow 0$ as $c \rightarrow \infty$, it follows that

$$(S^{(1/\tilde{b}(c))}([\lambda t \tilde{b}(c)]), \tilde{b}(c)^{-1}R(ct)) \Rightarrow (N_1(t), E(t))$$

in the J_1 topology of the product space $D([0, \infty), \mathbb{R} \times \mathbb{R})$, by [8, Theorem 3.2]. Since the process $E(t)$ is nondecreasing and continuous, [47, Theorem 13.2.4] along with the continuous mapping theorem yields

$$S^{(1/\tilde{b}(c))}([\lambda R(ct)]) = S^{(1/\tilde{b}(c))}([\lambda \cdot \tilde{b}(c)^{-1}R(ct) \cdot \tilde{b}(c)]) \Rightarrow N_1(E(t))$$

in the M_1 topology on $D([0, \infty), \mathbb{R})$. □

Remark 2.6. For the specific case of Mittag-Leffler waiting times, where $\mathbb{P}(J_n > t) = E_\beta(-t^\beta)$, we can take $b_n = n^{-1/\beta}$ in (2.5). To check this, note that

$$\mathbb{E}(e^{-sb_n T_n}) = \left(\frac{1}{1 + (sb_n)^\beta} \right)^n = \left(1 - \frac{s^\beta}{n + s^\beta} \right)^n \rightarrow e^{-s^\beta} = \mathbb{E}(e^{-sD(1)})$$

as $n \rightarrow \infty$. Then $\tilde{b}(c) = c^\beta$ and the CTRW convergence (2.7) reduces to $S^{(c^{-\beta})}([\lambda R(ct)]) \Rightarrow N_1(E(t))$ as $c \rightarrow \infty$. Substitute $p = c^{-\beta}$ to get

$$S^{(p)}([\lambda R(p^{-1/\beta} t)]) \Rightarrow N_1(E(t)), \quad \text{as } p \rightarrow 0. \quad (2.8)$$

3 Fractional calculus

This section develops some interesting connections between the fractional Poisson process and fractional calculus. In the process, some connections between results in the existing literature will be established. Behgin and Orsingher [6, Eq. (2.17)] show that the FPP of order $0 < \beta < 1$ has distribution

$$\mathbb{P}(N_\beta(t) = k) = \int_0^\infty e^{-\lambda x} \frac{(\lambda x)^k}{k!} V(x, t) dx, \quad (3.1)$$

where $V(x, t)$ is a “folded PDF” defined on $x > 0$, for each $t > 0$, by $V(x, t) = 2v(x, t)$, and $v(x, t)$ is another PDF with $x \in \mathbb{R}$ for each $t > 0$ that solves

$$\begin{aligned} \partial_t^{2\beta} v(x, t) &= \partial_x^2 v(x, t); \\ v(x, 0) &= \delta(x); \\ \partial_t v(x, 0) &\equiv 0, \quad \text{if } 1/2 < \beta < 1. \end{aligned} \quad (3.2)$$

It is also stated in [6, Eq. (1.9)] that the FPP $N_\beta(t) = N_1(T_t)$, where T_t is a random variable with PDF $V(x, t)$ for $t > 0$.

On the other hand, a simple conditioning argument shows that the equivalent FTTP process has distribution

$$\mathbb{P}(N_1(E(t)) = k) = \int_0^\infty P(N_1(x) = k) h(x, t) dx = \int_0^\infty e^{-\lambda x} \frac{(\lambda x)^k}{k!} h(x, t) dx \quad (3.3)$$

where $h(x, t)$ is the density of $E(t)$, a PDF on $x > 0$ for each $t > 0$. It follows from [37, Theorem 4.1] that this PDF solves

$$\partial_t^\beta h(x, t) = -\partial_x h(x, t); \quad h(x, 0) = \delta(x). \quad (3.4)$$

In view of Theorem 2.2, the two distributions (3.1) and (3.3) must be equal. Thus, the main purpose of this section is to reconcile the two fractional differential equations (3.2) and (3.4).

Theorem 3.1. *Let $N_\beta(t)$ be a fractional Poisson process (1.3) with $0 < \beta < 1$, so that (3.1) holds. Let $N_1(E(t))$ be the equivalent fractal time Poisson process, where $E(t)$ is the standard inverse β -stable subordinator with PDF $h(x, t)$, so that (3.3) holds. Then*

$$h(x, t) = 2v(x, t) \quad \text{for all } x > 0 \text{ and } t > 0. \quad (3.5)$$

In particular, the two fractional partial differential equations (3.2) and (3.4) are consistent, in the sense that the folded solution $V(x, t) = 2v(x, t)$ to (3.2) coincides with the solution $h(x, t)$ to (3.4).

Proof. Mainardi [31, Eq. (3.2)] shows that the solution to the fractional diffusion-wave equation (3.2) has LT

$$\tilde{v}(x, s) = \int_0^\infty e^{-st} v(x, t) dt = \frac{1}{2} s^{\beta-1} e^{-|x|s^\beta} \quad (3.6)$$

while [37, Eq. (3.13)] shows that

$$\tilde{h}(x, s) = s^{\beta-1} e^{-xs^\beta}. \quad (3.7)$$

Since both are differentiable in t , they are also continuous, so LT uniqueness for continuous functions implies (3.5).

Take Fourier transforms in (3.6) to see that the solution to (3.2) has Fourier-Laplace transform (FLT)

$$\bar{v}(k, s) = \int_0^\infty e^{-st} \int_{-\infty}^\infty e^{-ikx} v(x, t) dx dt = \frac{s^{2\beta-1}}{s^{2\beta} + k^2}, \quad (3.8)$$

where we have used the fact that $e^{-a|x|}$ has FT $2a/(a^2 + k^2)$. Rearrange to get

$$s^{2\beta} \bar{v}(k, s) - s^{2\beta-1} = -k^2 \bar{v}(k, s)$$

and invert the FT to get

$$s^{2\beta} \tilde{v}(x, s) - s^{2\beta-1} v(x, 0) = \partial_x^2 \tilde{v}(x, s),$$

using the fact that $\partial_x f(x)$ has FT $(ik)\hat{f}(k)$ and $v(x, 0) = \delta(x)$ has FT $\hat{v}(k, 0) \equiv 1$. To invert the LT, note that the Caputo fractional derivative $\partial_t^\beta f(t)$ has LT $s^\beta \hat{f}(s) - s^{\beta-1} f(0)$ if $0 < \beta \leq 1$, and LT $s^\beta \hat{f}(s) - s^{\beta-1} f(0) - s^{\beta-2} f'(0)$ if $1 < \beta \leq 2$. This is easy to verify from the definition (1.5), using the corresponding formula for the integer derivative, along with the fact that $s^{\beta-1}$ is the LT of $t^{-\beta}/\Gamma(1-\beta)$. Now use the remaining initial condition $\partial_t v(x, 0) \equiv 0$ for $1/2 < \beta < 1$ to invert the LT, and arrive at (3.2).

Likewise, the solution to (3.4) has FLT

$$\bar{h}(k, s) = \int_0^\infty e^{-ikx} s^{\beta-1} e^{-xs^\beta} dx = \frac{s^{\beta-1}}{s^\beta + ik}, \quad (3.9)$$

using the fact that $e^{ax} I(x \geq 0)$ has FT $1/(a + ik)$. To see that these are consistent, compute the FLT of $h(|x|, t)$:

$$\begin{aligned} \int_0^\infty e^{-st} \int_{-\infty}^\infty e^{-ikx} h(|x|, t) dx dt &= \int_0^\infty e^{-st} \left(\int_0^\infty e^{-ikx} h(x, t) dx + \int_0^\infty e^{ikx} h(x, t) dx \right) dt \\ &= \frac{s^{\beta-1}}{s^\beta + ik} + \frac{s^{\beta-1}}{s^\beta - ik} = 2 \left(\frac{s^{2\beta-1}}{s^{2\beta} + k^2} \right) = 2\bar{v}(k, s). \end{aligned}$$

Invert the FLT to see that $h(|x|, t) = 2v(x, t)$ for all $x \in \mathbb{R}$ and $t > 0$. To verify the LFT solution, take FT in (3.4) to get

$$\partial_t^\beta \hat{h}(k, t) = -ik \hat{h}(k, t)$$

and apply the LT to get $s^\beta \bar{h}(k, s) - s^{\beta-1} = -ik \bar{h}(k, s)$, using the point source initial condition $\hat{h}(k, 0) \equiv 1$. \square

Remark 3.2. Behgin and Orsingher [6, Eq. (2.21)] show that

$$v(x, t) = \frac{1}{2\Gamma(1-\beta)} \int_0^t (t-w)^{-\beta} p(|x|, t) dx,$$

where $p(x, t)$ is the density of the stable subordinator $D(t)$, while [37, Theorem 3.1] implies that

$$h(x, t) = \frac{1}{\Gamma(1-\beta)} \int_0^t (t-w)^{-\beta} p(x, t) dx.$$

This gives another proof that $h(|x|, t) = 2v(x, t)$.

Remark 3.3. A closed form expression for the probability mass function $p(n, t) = \mathbb{P}(N_1(E(t)) = n) = P(N_\beta(t) = n)$ is

$$p(n, t) = \frac{(\lambda t^\beta)^n}{n!} \sum_{r=0}^{\infty} \frac{(n+r)!}{r!} \frac{(-\lambda t^\beta)^r}{\Gamma(\beta(r+n)+1)}. \quad (3.10)$$

See for example Jumarie [25], Laskin [30], Beghin and Orsingher [6, 7] and Cahoy [11] for different derivations of this result

The equivalence in Theorem 3.1 results from folding the solution to the fractional diffusion-wave equation (3.2). Another fractional partial differential equation for the density $h(x, t)$ of the standard inverse β -stable subordinator $E(t)$, which is closer to the form (3.2), can be obtained by arguments similar to those used in [4] to connect the inverse stable subordinator to iterated Brownian motion. In that theory, it is customary to avoid distributions by imposing a functional initial condition.

Theorem 3.4. *Let $E(t)$ be the standard inverse β -stable subordinator with density $h(x, t)$. Then for any $f \in L_2(\mathbb{R}) \cap C^1(\mathbb{R})$, the function*

$$u(x, t) = \mathbb{E}_x[f(E(t))] = \int_0^\infty f(x+y)h(y, t)dy \quad (3.11)$$

solves the fractional differential equation

$$\partial_t^{2\beta} u(x, t) = -\partial_x f(x) \frac{t^{\beta-1}}{\Gamma(1-\beta)} + \partial_x^2 u(x, t); \quad u(x, 0) = f(x). \quad (3.12)$$

In particular, when $\beta = 1/2$, (3.11) solves

$$\partial_t u(x, t) = \frac{-\partial_x f(x)}{\sqrt{\pi t}} + \partial_x^2 u(x, t); \quad u(0, x) = f(x), \quad (3.13)$$

and in this case we also have $u(x, t) = \mathbb{E}_x[f(|B(t)|)]$, where $B(t)$ is a Brownian motion with variance $2t$.

Remark 3.5. Equation (3.13) in Theorem 3.4 is a special case of Theorem 0.1 in [1] using the Markov process $X(t) = X(0) - t$. The semigroup associated with $X(t)$ is $T(t)f(x) = \mathbb{E}_x[f(X(t))] = f(x-t)$, with infinitesimal generator $-\partial_x$.

Proof of Theorem 3.4. From (3.9), we have

$$\bar{u}(k, s) = \frac{s^{\beta-1} \hat{f}(k)}{s^\beta + ik} = \frac{s^{\beta-1} \hat{f}(k)}{s^\beta + ik} \cdot \frac{s^\beta - ik}{s^\beta - ik} = \frac{s^{2\beta-1} - iks^{\beta-1}}{s^{2\beta} + k^2} \hat{f}(k)$$

so that $s^{2\beta} \bar{u}(k, s) - s^{2\beta-1} \hat{f}(k) = -ik \hat{f}(k) s^{\beta-1} - k^2 \bar{u}(k, s)$, which inverts to (3.12). It is well known that the Brownian motion first passage time $D(y) = \inf\{t > 0 : B(t) > y\}$ is a stable subordinator with index $\beta = 1/2$ [2, Example 1.3.19]. Then it is easy to see that

$$E(t) = \inf\{y > 0 : D(y) > t\} = \sup\{B(r) : 0 \leq r \leq t\}$$

and this recovers the fact, typically proven using the reflection principle, that

$$P(E(t) > y) = 2P(B(t) > y).$$

Then $E(t)$ and $|B(t)|$ have the same one dimensional distributions, so we also have $u(x, t) = \mathbb{E}_x[f(|B(t)|)]$. Note that $X(t) = X(0) - t$ is *a fortiori* a continuous Markov process associated with the shift semigroup $T(t)f(x) = \mathbb{E}_x[f(X(t))] = f(x - t)$ with generator

$$L_x f(x) = \lim_{t \rightarrow 0^+} \frac{T(t)f(x) - f(x)}{t} = -\partial_x f(x).$$

Then [4, Corollary 3.4] and [1, Theorem 0.1] imply that $u(x, t)$ solves the equation

$$\partial_t u(x, t) = \frac{L_x f(x)}{\sqrt{\pi t}} + L_x^2 u(x, t); \quad u(0, x) = f(x).$$

When $L_x = -\partial_x$, this reduces to (3.13), a special case of (3.12) with $\Gamma(1/2) = \sqrt{\pi}$. □

Remark 3.6. In the case $f(x) = \delta(x)$, Theorem 3.4 gives an alternative governing equation for $h(x, t)$. Note that (3.12) is very similar to the governing equation (3.2) for the unfolded PDF

Remark 3.7. The process $|B(t)|$ in Theorem 3.4 is not the same process as the inverse 1/2-stable subordinator $E(t)$ in Theorem 2.2, although they have the same one dimensional distributions. One reason is the fact that as a process $E(t)$ has nondecreasing sample paths, whereas $|B(t)|$ does not have this property. Hence, the FTTP $N_1(E(t))$ is not the same as the Brownian time Poisson process $N_1(|B(t)|)$. However, we do have $E(t) = \sup\{B(r) : 0 \leq r \leq t\}$, see for example [10], so that a Poisson process subordinated to the supremum of a Brownian motion is an FPP with $\beta = 1/2$.

Remark 3.8. Let $E(t)$ be the standard inverse stable subordinator of index $\beta = 1/m$ for integer $m > 1$. Then [4, Remark 3.11], [40, Theorem 1.1] and Keyantuo and Lizama [26, Theorem 3.3] imply that $u(x, t) = \mathbb{E}_x[f(E(t))]$ solves

$$\partial_t u(x, t) = \sum_{j=1}^{m-1} \frac{t^{j/m-1}}{\Gamma(j/m)} (-\partial_x)^j f(x) + (-\partial_x)^m u(x, t); \quad u(0, x) = f(x),$$

for $t > 0$ and $x \in \mathbb{R}$, which is then equivalent to (3.4). The proof is similar to Theorem 3.4. For example, when $\beta = 1/3$ use

$$\bar{u}(s, k) = \frac{s^{-2/3} \hat{f}(k)}{s^{1/3} + ik} \cdot \frac{s^{2/3} - s^{1/3} ik + k^2}{s^{2/3} - s^{1/3} ik + k^2} = \frac{1 - s^{-1/3} ik + s^{-2/3} k^2}{s + ik^3} \hat{f}(k).$$

4 Renewal processes and inverse subordinators

Theorem 2.2 shows that a Poisson process, time-changed by an inverse stable subordinator, yields a renewal process with Mittag-Leffler waiting times. This section extends that result to arbitrary subordinators that are strictly increasing. Let $D_\psi(t)$ be a strictly increasing Lévy process (subordinator) with $\mathbb{E}[e^{-sD_\psi(t)}] = e^{-t\psi(s)}$, where the Laplace exponent

$$\psi(s) = bs + \int_0^\infty (e^{-sx} - 1)\phi(dx), \tag{4.1}$$

$b \geq 0$, and ϕ is the Lévy measure of D_ψ . Then we must have either

$$\phi(0, \infty) = \infty, \tag{4.2}$$

or $b > 0$, or both. Let

$$E_\psi(t) = \inf\{\tau \geq 0 : D_\psi(\tau) > t\} \quad (4.3)$$

be the inverse subordinator (1.7), and recall that $N_1(t)$ is a Poisson process with rate λ .

Theorem 4.1. *The time-changed Poisson process $N_1(E_\psi(t))$ is a renewal process whose IID waiting times (J_n) satisfy*

$$\mathbb{P}(J_n > t) = \mathbb{E}[e^{-\lambda E_\psi(t)}]. \quad (4.4)$$

Proof. The proof is similar to Theorem 2.2. Take $N_1(t) = \max\{n \geq 0 : V_n \leq t\}$, where $V_n = W_1 + \dots + W_n$, with W_n IID as $\mathbb{P}(W_n > t) = e^{-\lambda t}$. Let

$$\tau_n = \sup\{t > 0 : N_1(E_\psi(t)) < n\} = \sup\{t > 0 : E_\psi(t) < V_n\}$$

and apply Lemma 2.1 to get $\tau_n = D_\psi(V_n^-)$. Then, as in the proof of Theorem 2.2, we have

$$\begin{aligned} \mathbb{E}(e^{-s\tau_1}) &= \mathbb{E}(e^{-sD_\psi(W_1^-)}) \\ &= \mathbb{E}\left[\mathbb{E}\left(e^{-sD_\psi(W_1)} \mid W_1\right)\right] = \mathbb{E}\left[e^{-W_1\psi(s)}\right] = \frac{\lambda}{\lambda + \psi(s)}. \end{aligned} \quad (4.5)$$

By [37, Corollary 3.5], the IID random variables J_n in (4.4) satisfy

$$\int_0^\infty e^{-st} \mathbb{P}(J_n > t) dt = \int_0^\infty e^{-st} \mathbb{E}[e^{-\lambda E_\psi(t)}] dt = \frac{\psi(s)}{s(\lambda + \psi(s))}. \quad (4.6)$$

Integrate by parts to get

$$\int_0^\infty e^{-st} \mathbb{P}_{J_n}(dt) = \int_0^\infty se^{-st} [1 - \mathbb{P}(J_n > t)] dt = 1 - \frac{\psi(s)}{\lambda + \psi(s)} = \frac{\lambda}{\lambda + \psi(s)}, \quad (4.7)$$

which shows that $T_1 = J_1$ is identically distributed with τ_1 . Extend this argument, as in the proof of Theorem 2.2, to show that (T_1, \dots, T_n) is identically distributed with (τ_1, \dots, τ_n) for any positive integer n . For example, when $n = 2$, write

$$\mathbb{E}(e^{-s_1 D_\psi(t_1)} e^{-s_2 D_\psi(t_1+t_2)}) = \mathbb{E}(e^{-(s_1+s_2)D_\psi(t_1)} e^{-s_2[D_\psi(t_1+t_2)-D_\psi(t_1)]}) = e^{-t_1\psi(s_1+s_2)} e^{-t_2\psi(s_2)}$$

and condition to get

$$\begin{aligned} \mathbb{E}(e^{-s_1\tau_1-s_2\tau_2}) &= \mathbb{E}(e^{-s_1 D_\psi(W_1^-)-s_2 D_\psi([W_1+W_2]^-)}) \\ &= \mathbb{E}\left[\mathbb{E}\left(e^{-s_1 D_\psi(W_1)-s_2 D_\psi(W_1+W_2)} \mid W_1, W_2\right)\right] \\ &= \mathbb{E}\left[e^{-W_1\psi(s_1+s_2)} e^{-W_2\psi(s_2)}\right] = \frac{\lambda}{\lambda + \psi(s_1 + s_2)} \cdot \frac{\lambda}{\lambda + \psi(s_2)}. \end{aligned}$$

On the other hand,

$$\mathbb{E}(e^{-s_1 T_1} e^{-s_2 T_2}) = \mathbb{E}(e^{-s_1 J_1} e^{-s_2 (J_1+J_2)}) = \frac{\lambda}{\lambda + \psi(s_1 + s_2)} \cdot \frac{\lambda}{\lambda + \psi(s_2)}$$

using the fact that (J_n) are IID. To finish the proof, use continuous mapping to show that (J_1, \dots, J_n) is identically distributed with (X_1, \dots, X_n) , where $X_n = \tau_n - \tau_{n-1}$ are the waiting times between jumps for the process $N_1(E_\psi(t))$. \square

Remark 4.2. Let $N_\psi(t)$ denote the renewal process from Theorem 4.1, so that

$$N_\psi(t) = \max\{n \geq 0 : T_n \leq t\}, \quad (4.8)$$

where $T_n = \sum_{i=1}^n J_i$ and (J_n) are IID according to (4.4). Theorem 4.1 shows that $N_\psi(t) = N_1(E_\psi(t))$. This extends the relation $N_\beta(t) = N_1(E(t))$ from Theorem 2.2, the special case of an inverse stable subordinator $E(t)$ and Mittag-Leffler waiting times J_n , to a general inverse subordinator.

Remark 4.3. Let $M(t) = \mathbb{E}(N_\psi(t))$ denote the renewal function of the renewal process $N_\psi(t)$. Then using Lageras [29, Equation 4], it follows that the LT of $M(t)$ is $\lambda/\psi(s)$.

Remark 4.4. A simple conditioning argument shows that

$$p_\psi(n, t) = \mathbb{P}(N_\psi(t) = n) = \int_0^\infty P(N_1(x) = n) h_\psi(x, t) dx = \int_0^\infty e^{-\lambda x} \frac{(\lambda x)^n}{n!} h(x, t) dx$$

where $h_\psi(x, t)$ is the density of $E_\psi(t)$. A straightforward calculation shows that

$$\tilde{p}_\psi(n, s) = \frac{s^{-1}\psi(s)}{\lambda + \psi(s)} \cdot \frac{\lambda^n}{(\lambda + \psi(s))^n}. \quad (4.9)$$

Use (4.6) and (4.7) to see that the first factor in (4.9) is the LT ($t \mapsto s$) of $\tilde{h}_\psi(\lambda, t) = \mathbb{P}(J_n > t) = \mathbb{E}(e^{-\lambda E_\psi(t)})$, and the second is the LT of $T_n = J_1 + J_2 + \dots + J_n$ with J_n IID as in (4.4). Denote the distribution of T_n by $F^{(n,*)}$, the n -fold convolution of the distribution function F of J_1 . Invert the LT to get

$$p_\psi(n, t) = \int_0^t \tilde{h}_\psi(\lambda, t-s) F^{(n,*)}(ds) \quad (4.10)$$

which extends (3.10).

5 CTRW scaling limits and governing equations

In this section, we extend the fractional calculus results of Section 3 to the inverse subordinators of Section 4. A general theory of CTRW scaling limits and governing equations is developed in [37]. Consider a sequence of CTRW indexed by a scale parameter $c > 0$. Take J_n^c nonnegative IID random variables representing the waiting times between particle jumps and $T^c(n) = \sum_{i=1}^n J_i^c$, the time of the n th jump. Let Y_i^c be IID random vectors on \mathbb{R}^d representing the particle jumps, independent of the waiting times, and set $S^c(n) = \sum_{i=1}^n Y_i^c$, the location of the particle after n jumps. Define $N_t^c = \max\{n \geq 0 : T^c(n) \leq t\}$, the number of jumps by time $t \geq 0$ and

$$X^c(t) = S^c(N_t^c) = \sum_{i=1}^{N_t^c} Y_i^c \quad (5.1)$$

the position of the particle at time $t \geq 0$ and scale $c > 0$. Assume a triangular array limit

$$\{(S^c(ct), T^c(ct))\}_{t \geq 0} \Rightarrow \{(A(t), D_\psi(t))\}_{t \geq 0}, \quad \text{as } c \rightarrow \infty, \quad (5.2)$$

in the J_1 topology on $D([0, \infty), \mathbb{R}^d \times \mathbb{R}_+)$, so that $A(t)$ and $D_\psi(t)$ are independent Lévy processes on \mathbb{R}^d and \mathbb{R} , respectively. Since the waiting times are nonnegative, $D_\psi(t)$ is a subordinator. In this section, we assume the drift $b = 0$ in (4.1), as well as condition (4.2) and

$$\int_0^1 y |\ln y| \phi(dy) < \infty. \quad (5.3)$$

Assumption (4.2) implies that the process $\{D_\psi(t)\}$ is strictly increasing, i.e., $D_\psi(t)$ is not compound Poisson. Then [37, Theorem 3.1] shows that the inverse subordinator $E_\psi(t)$ in (4.3) has a Lebesgue density

$$h_\psi(x, t) = \int_0^t \phi(t - y, \infty) \mathbb{P}_{D_\psi(x)}(dy). \quad (5.4)$$

Write $\mathbb{E}[e^{-sD_\psi(t)}] = e^{-t\psi(s)}$, as before. Let $P(x, t) = \mathbb{P}(A(t) \leq x)$ be the distribution function of $A(t)$, and write

$$\hat{P}(k, t) = \int e^{-ik \cdot x} P(dx, t) = e^{-t\psi_A(k)},$$

where $\psi_A(k)$ is the Fourier symbol of A . The symbols define pseudo-differential operators: $\psi(\partial_t)f(t)$ has LT $\psi(s)\tilde{f}(s)$, and $\psi_A(-i\partial_x)f(x)$ has FT $\psi_A(k)\hat{f}(k)$, for suitable functions f . Then [37, Theorem 2.1] establishes the CTRW scaling limit

$$\{X^c(t)\}_{t \geq 0} \Rightarrow \{A(E_\psi(t))\}_{t \geq 0}, \quad \text{as } c \rightarrow \infty, \quad (5.5)$$

in the M_1 -topology on $D([0, \infty), \mathbb{R}^d)$. Recall that a function Q is a *mild solution* to a space-time pseudo-differential equation if its (Fourier-Laplace or Laplace-Laplace) transform solves the equivalent algebraic equation in transform space. The next result is a small extension and simplification of [37, Theorem 4.1].

Theorem 5.1. *Assume (5.2) holds, where $D_\psi(t)$ is a subordinator without drift such that conditions (4.2) and (5.3) hold. The distribution function of the CTRW limit process $A(E_\psi(t))$ in (5.5) is given by*

$$Q(x, t) = \int_0^\infty P(x, u) h_\psi(u, t) du \quad (5.6)$$

where $h_\psi(u, t)$ is the density (5.4) of the inverse subordinator $E_\psi(t)$. The distribution function $Q(x, t)$ solves the generalized Cauchy problem

$$\psi(\partial_t)Q(x, t) = -\psi_A(-i\partial_x)Q(x, t) + H(x)\phi(t, \infty) \quad (5.7)$$

in the mild sense, where $H(x) = I(x \geq 0)$ is the Heaviside function. Furthermore, $P(x, u)$ solves the Cauchy problem

$$\partial_t P(x, t) = -\psi_A(-i\partial_x)P(x, t); \quad P(x, 0) = H(x), \quad (5.8)$$

and $h_\psi(x, t)$ solves the Cauchy problem (in the x variable)

$$\partial_x h_\psi(x, t) = -\psi(\partial_t)h_\psi(x, t); \quad h_\psi(0, t) = \phi(t, \infty). \quad (5.9)$$

Proof. The proof is similar to [37, Theorem 4.1]. Equation (5.6) follows from a simple conditioning argument. Apply [37, Theorem 3.6] to see that $Q(x, t)$ has FLT

$$\bar{Q}(k, s) = \int_0^\infty e^{-st} \int_{\mathbb{R}^d} e^{-ik \cdot x} Q(dx, t) dt = \frac{1}{s} \frac{\psi(s)}{\psi_A(k) + \psi(s)} \quad (5.10)$$

and rearrange to get

$$\psi(s) \bar{Q}(k, s) = -\psi_A(k) \bar{Q}(k, s) + s^{-1} \psi(s). \quad (5.11)$$

From [37, Eq. (3.12)] we get

$$\int_0^\infty e^{-su} \phi(u, \infty) du = s^{-1} \psi(s). \quad (5.12)$$

Now invert the FLT (5.11), using (5.12) and $\int e^{-ik \cdot x} H(dx) \equiv 1$, to arrive at (5.7). It is well known that $P(x, t)$ solves the Cauchy problem (5.8), see for example [23].

Equation (4.6) shows that the bivariate Laplace transform (LLT)

$$\bar{h}_\psi(\lambda, s) = \int_0^\infty \int_0^\infty e^{-\lambda x - st} h_\psi(x, t) dt dx = \frac{1}{s} \frac{\psi(s)}{\lambda + \psi(s)}.$$

This rearranges to

$$\lambda \bar{h}_\psi(\lambda, s) - s^{-1} \psi(s) = -\psi(s) \bar{h}_\psi(\lambda, s). \quad (5.13)$$

Using (5.12) again, it follows from (5.4) and the convolution property of the LT that

$$\tilde{h}(x, s) = \int_0^\infty e^{-st} h_\psi(x, t) dx = s^{-1} \psi(s) e^{-x\psi(s)}$$

and hence $\tilde{h}(0, s) = s^{-1} \psi(s)$. Now invert the Laplace transform in the x variable in equation (5.13) to get

$$\partial_x \tilde{h}(x, s) = -\psi(s) \tilde{h}(x, s),$$

and then invert the LT in the t variable to see that $h_\psi(x, t)$ solves (5.9). \square

For any random walk $S(n) = \sum_{i=1}^n Y_i$, the compound Poisson process $A(t) = S(N_1(t))$ is a Lévy process. Introduce IID waiting times (4.4) between these random walk jumps to get a CTRW. In this case, the CTRW is exactly of the form $A(E_\psi(t))$, without passing to the limit. Then the governing equations in Theorem 5.1 pertain to the CTRW itself.

Theorem 5.2. *Assume $D_\psi(t)$ is a subordinator without drift such that conditions (4.2) and (5.3) hold, and let $E_\psi(t)$ be the inverse subordinator (4.3). Take J_n IID waiting times according to (4.4), and let $N_\psi(t)$ denote the renewal process (4.8). Take Y_n IID jumps on \mathbb{R}^d , independent from (J_n) , with common distribution μ , and let $S(n) = \sum_{i=1}^n Y_i$. Then the distribution function $P(x, t) = \mathbb{P}(X(t) \leq x)$ of the CTRW $X(t) = S(N_\psi(t))$ solves the generalized Cauchy problem*

$$\psi(\partial_t) P(x, t) = -\lambda P(x, t) + \lambda \int P(x - y, t) \mu(dy) + H(x) \phi(t, \infty) \quad (5.14)$$

in the mild sense. Furthermore, $X(t) = A(E_\psi(t))$, where $A(t) = S(N_1(t))$ is a compound Poisson process.

Proof. Theorem 4.1 yields $N_\psi(t) = N_1(E_\psi(t))$, and then the CTRW is

$$X(t) = S(N_\psi(t)) = S(N_1(E_\psi(t))) = A(E_\psi(t)).$$

A standard conditioning argument shows that the compound Poisson FT $\hat{P}(k, t) = e^{-t\psi_A(k)}$, where the Fourier symbol $\psi_A(k) = \lambda(1 - \hat{\mu}(k))$. The inverse FT of $\psi_A(k)\hat{f}(k)$ is

$$\psi_A(-i\partial_x)f(x) = -\lambda f(x) + \lambda \int f(x-y)\mu(dy) \quad (5.15)$$

using the FT convolution property. Now Theorem 5.1 implies that (5.14) holds. \square

Remark 5.3. In the situation of Theorem 5.2, where $A(t)$ is compound Poisson, the distribution function $P(x, t) = \mathbb{P}(A(t) \leq x)$ solves the Cauchy problem (5.8), which can be written in this case as

$$\partial_t P(x, t) = -\lambda P(x, t) + \lambda \int_{-\infty}^{\infty} P(x-y, t)\mu(dy); \quad P(x, 0) = H(x). \quad (5.16)$$

This is the Kolmogorov forward equation for the Markov process $A(t)$. If μ has density $w(x)$, apply ∂_x on both sides of (5.16) to see that the probability density $p(x, t) = \partial_x P(x, t)$ of $A(t)$ solves (1.6). If D is the stable subordinator with Laplace symbol $\psi(s) = s^\beta$, then (5.14) holds with $\phi(t, \infty) = t^{-\beta}/\Gamma(1-\beta)$ and $\psi(\partial_t) = \mathbb{D}_t^\beta$, the Riemann-Liouville fractional derivative. The Riemann-Liouville fractional derivative is defined for $0 \leq n-1 < \beta < n$ by

$$\mathbb{D}_t^\beta g(t) = \frac{1}{\Gamma(n-\beta)} \frac{d^n}{dt^n} \int_0^t (t-r)^{n-1-\beta} g^{(n)}(r) dr, \quad (5.17)$$

which differs from the Caputo derivative (1.5) in that the derivative is applied after the integration. The LT of $\mathbb{D}_t^\beta g(t)$ is $s^\beta \tilde{g}(s)$. Apply ∂_x to both sides of (5.14) in this case to get

$$\mathbb{D}_t^\beta p(x, t) = -\lambda p(x, t) + \lambda \int p(x-y, t)\mu(dy) + \delta(x) \frac{t^{-\beta}}{\Gamma(1-\beta)},$$

the fractional kinetic equation of Zaslavsky [48]. To recover (1.4), use $\partial_t^\beta g(t) = \mathbb{D}_t^\beta g(t) - g(0)t^{-\beta}/\Gamma(1-\beta)$ and $p(x, 0) = \delta(x)$.

Remark 5.4. In the special case where $\mu = \varepsilon_1$ is a point mass, so that $Y_n = 1$ almost surely, $A(t) = N_1(t)$ is a Poisson process with rate $\lambda > 0$. Then the distribution function $P(x, t)$ of the renewal process $N_\psi(t) = A(E_\psi(t))$ solves

$$\psi(\partial_t)P(x, t) = -\lambda[P(x, t) - P(x-1, t)] + H(x)\phi(t, \infty). \quad (5.18)$$

If D is the stable subordinator with Laplace symbol $\psi(s) = s^\beta$, Equation (5.18) reduces to

$$\partial_t^\beta P(x, t) = -\lambda[P(x, t) - P(x-1, t)]$$

as in Remark 5.3. The probability mass function $p(n, t) = P(n, 1) - P(n-1, t) = \Delta P(n, t)$ for $n > 0$. Apply the difference operator Δ on both sides to obtain

$$\partial_t^\beta p(n, t) = -\lambda[p(n, t) - p(n-1, t)]$$

as in Jumarie [25].

Remark 5.5. Scher and Lax [44] showed that a CTRW with waiting time distribution ω and jump distribution ν has FLT

$$\bar{Q}(k, s) = \frac{1}{s} \frac{1 - \tilde{\omega}(s)}{1 - \tilde{\omega}(s)\hat{\nu}(k)},$$

where $\hat{\nu}(k) = \int e^{-ik \cdot x} \nu(dx)$. To reconcile with Theorem 5.2, recall from (4.7) that the waiting times (4.4) in Theorem 5.2 have LT

$$\tilde{\omega}(s) = \int e^{-st} \omega(dt) = \frac{\lambda}{\lambda + \psi(s)}$$

and then it follows that $\psi(s) = \lambda(1 - \tilde{\omega}(s))/\tilde{\omega}(s)$. The jumps Y_n in Theorem 5.2 have Fourier symbol $\psi_A(k) = \lambda(1 - \hat{\mu}(k))$ and then (5.10) implies

$$\bar{Q}(k, s) = \frac{1}{s} \frac{\psi(s)}{\psi_A(k) + \psi(s)} = \frac{1}{s} \frac{\frac{1 - \tilde{\omega}(s)}{\tilde{\omega}(s)}}{\frac{1 - \tilde{\omega}(s)}{\tilde{\omega}(s)} + (1 - \hat{\mu}(k))} = \frac{1}{s} \frac{1 - \tilde{\omega}(s)}{1 - \tilde{\omega}(s)\hat{\mu}(k)}$$

which provides a different proof that the CTRW equals $A(E_\psi(t))$ in this case. To simulate the sample paths of the non-Markovian process $A(E_\psi(t))$, it is sufficient to simulate the CTRW. In particular, the renewal process $N_\psi(t)$ gives the exact jump times of the inverse subordinator $E_\psi(t)$.

Remark 5.6. In the general case, where $A(t)$ is not compound Poisson, Theorem 5.2 provides a useful approximation. Given a Lévy process $A(t)$, take $Y_n = A(n) - A(n-1)$, so that $S(n) = A(n)$. Take $N(t)$ a Poisson process with rate 1, so that $S(\lambda^{-1}N(\lambda t))$ is compound Poisson with Fourier symbol

$$\lambda(1 - e^{-\lambda^{-1}\psi_A(k)}) \rightarrow \psi_A(k), \quad \text{as } \lambda \rightarrow \infty.$$

Then $S(\lambda^{-1}N(\lambda t)) \Rightarrow A(t)$ as $\lambda \rightarrow \infty$, and the CTRW with IID waiting times (4.4) and these compound Poisson jumps converges to $A(E_\psi(t))$ as $\lambda \rightarrow \infty$. As in Remark 5.5, this fact can be used to simulate sample paths of the process $A(E_\psi(t))$. This fact has been exploited by Fulger, Scalas and Germano [20] to develop fast simulation methods for space-time fractional diffusion equations.

Example 5.7. Tempered stable subordinators are theoretically interesting [5, 43] and practically useful [18, 38]. Take $D_a(t)$ tempered stable with Laplace symbol $\psi(s) = (s+a)^\beta - a^\beta$ for $a > 0$ and $0 < \beta < 1$, and let $E_a(t)$ be its inverse (1.7). Theorem 4.1 shows that $N_1(E_a(t))$ is a renewal process. Let (τ_n) denote the arrival times of this renewal process, and use (4.5) to get

$$\mathbb{E}(e^{-s\tau_1}) = \frac{\lambda}{\lambda + (s+a)^\beta - a^\beta}.$$

This tempered fractional Poisson process $N_1(E_a(t))$ has tempered Mittag-Leffler waiting times, but with a different rate parameter: Use (2.4) to see that the Mittag-Leffler PDF $f(t) = \partial_t [1 - E_\beta(-\eta t^\beta)]$ has Laplace transform $\eta/(\eta + s^\beta)$, and so

$$\int_0^\infty e^{-st} f(t) e^{-at} dt = \frac{\eta}{\eta + (s+a)^\beta}.$$

Of course $f(t)e^{-at}$ is not a PDF, and in fact we have (set $s = 0$ above)

$$\int_0^\infty f(t) e^{-at} dt = \frac{\eta}{\eta + a^\beta}.$$

Then the tempered Mittag-Leffler PDF $f_a(t) = f(t)e^{-at}(\eta + a^\beta)/\eta$ has LT

$$\int_0^\infty e^{-st} f_a(t) dt = \frac{\eta + a^\beta}{\eta + (s + a)^\beta} = \frac{\lambda}{\lambda + (s + a)^\beta - a^\beta} = \mathbb{E}(e^{-s\tau_1})$$

when $\eta + a^\beta = \lambda$. Cartea and Del-Castillo [14] show that the tempered fractional derivative $\psi(\partial_t)g(t) = e^{-at} \partial_t^\beta [e^{at} g(t)] - a^\beta g(t)$. It is also known (e.g., see [5]) that the corresponding Lévy measure is exponentially tempered: $\phi(dt) = e^{-at} \phi_\beta(dt)$, where $\phi_\beta(t, \infty) = t^{-\beta}/\Gamma(1 - \beta)$ is the Lévy measure of the standard β -stable subordinator. Then Theorem 5.2 shows that the CTRW with tempered Mittag-Leffler waiting times and compound Poisson jumps solves a tempered fractional Cauchy problem

$$e^{-at} \partial_t^\beta [e^{at} P(x, t)] - a^\beta P(x, t) = \psi_A(-i\partial_x)P(x, t) + H(x)\phi(t, \infty)$$

with $\psi_A(-i\partial_x)$ given by (5.15) and $\phi(t, \infty) = \beta \int_t^\infty e^{-at} t^{-\beta-1} dt / \Gamma(1 - \beta)$. More generally, Theorem 5.1 shows that the distribution function of the CTRW scaling limit $A(E_\psi(t))$ is governed by this equation, with the corresponding operator $\psi_A(-i\partial_x)$. Apply ∂_x on both sides of (5.18) to see that the PDF of the renewal process with tempered Mittag-Leffler waiting times solves

$$e^{-at} \partial_t^\beta [e^{at} p(x, t)] - a^\beta p(x, t) = -\lambda[p(x, t) - p(x - 1, t)] + \delta(x)\phi(t, \infty).$$

A wide variety of tempered stable models in \mathbb{R}^d are discussed in Rosiński [43]. Random walks in \mathbb{R}^d with tempered stable scaling limit are developed in [15]. For exponentially tempered stable waiting times in \mathbb{R}^1 , a renewal process with tempered Mittag-Leffler waiting times gives the same process exactly, without taking limits. This can be useful for simulating sample paths.

Example 5.8. Checkin et al. [17, 16] used distributed order fractional derivatives to model multi-scale anomalous subdiffusion, where a different power law pertains at short and long time scales, and ultraslow diffusion, for a plume of particles spreading at a logarithmic rate. Given a finite Borel measure ν on $(0, 1)$, the distributed order fractional derivative is defined by

$$\mathbb{D}_t^\nu g(t) = \int_0^1 \partial_t^\beta g(t) \nu(d\beta), \quad (5.19)$$

where ∂_t^β is the Caputo fractional derivative (1.5). If ν is discrete, this is a linear combination of fractional derivatives. Let $D_\nu(t)$ be the distributed order stable subordinator with Laplace symbol $\psi(s) = \int s^\beta \nu(d\beta)$ and $E_\nu(t)$ its inverse (4.3). Let $\nu(d\beta) = p(\beta)d\beta$ for some $p \in C^1(0, 1)$, then by (2.19) in Kochubei [27]

$$P(J_n > t) = \mathbb{E}(e^{-\lambda E_\nu(t)}) = \frac{\lambda}{\pi} \int_0^\infty r^{-1} e^{-tr} \Phi(r, 1) dr \quad (5.20)$$

where

$$\Phi(r, 1) = \frac{\int_0^1 r^\beta \sin(\beta\pi) \Gamma(1 - \beta) p(\beta) d\beta}{[\int_0^1 r^\beta \cos(\beta\pi) \Gamma(1 - \beta) p(\beta) d\beta + \lambda]^2 + [\int_0^1 r^\beta \sin(\beta\pi) \Gamma(1 - \beta) p(\beta) d\beta]^2}.$$

Substitute (5.20) into (4.10) to obtain an explicit formula for the probability mass function of the distributed order Poisson process. In the special case where $\psi(s) = s^{\beta_0}$ for some $0 < \beta_0 < 1$, we

have $p(\beta) = \Gamma(1 - \beta)^{-1} \delta(\beta - \beta_0)$ and $D_\nu(t)$ is a stable subordinator. In this case (5.20) reduces to the following formula for the Mittag-Leffler function [21, Equation (3.25)]:

$$E_{\beta_0}(-\lambda t^{\beta_0}) = \frac{\lambda}{\pi} \int_0^\infty r^{-1} e^{-tr} \frac{r^{\beta_0} \sin(\beta_0 \pi)}{[r^{\beta_0} \cos(\beta_0 \pi) + \lambda]^2 + [r^{\beta_0} \sin(\beta_0 \pi)]^2} dr$$

If $\nu(d\beta) = p(\beta)d\beta$, where $p(\beta)$ is regularly varying at $\beta = 0$ with index $\alpha - 1$ for some $\alpha > 0$, then $\psi(s) = R(\log s)$ and R is regularly varying at infinity with index $-\alpha$, see [36, Lemma 3.1]. Then $E_\nu(t)$ is “ultraslow” in that $\mathbb{E}(E_\nu(t)^\gamma) = S(\log t)$, where S varies regularly with index $\gamma\alpha$ at infinity, by [36, Theorem 3.9]. Take an IID sequence of mixing variables (B_i) with distribution μ concentrated on $(0, 1)$, and assume $\mathbb{P}(J_i^c > u | B_i = \beta) = c^{-1}u^{-\beta}$ for $u \geq c^{-1/\beta}$, so that the waiting times are conditionally Pareto. Then [36, Theorem 3.4] implies that the distributed order stable subordinator is a random walk limit $\sum_{i=1}^{\lfloor ct \rfloor} J_i^c \Rightarrow D_\nu(t)$. This requires $\int (1 - \beta)^{-1} \mu(d\beta) < \infty$ so that $\nu(d\beta) = \Gamma(1 - \beta)\mu(d\beta)$ is a finite measure. An easy computation shows that the Lévy measure $\phi(t, \infty) = \int_0^1 t^{-\beta} \nu(d\beta) / \Gamma(1 - \beta)$. Then Theorem 5.1 implies that a CTRW with these conditionally Pareto waiting times has a scaling limit $A(E_\nu(t))$ whose distribution $Q(x, t)$ solves the distributed-order fractional diffusion equation

$$\mathbb{D}_t^\nu Q(x, t) = -\psi_A(-i\partial_x)Q(x, t).$$

If $A(t)$ is compound Poisson, Theorem 5.2 shows that the distribution function $P(x, t)$ of a CTRW with waiting times (4.4) solves

$$\mathbb{D}_t^\nu P(x, t) = -\lambda P(x, t) + \lambda \int P(x - y, t) \mu(dy),$$

without passing to the limit. Then the PDF $p(x, t)$ of the renewal process with waiting times (4.4) solves

$$\mathbb{D}_t^\nu p(x, t) = -\lambda[p(x, t) - p(x - 1, t)].$$

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References

- [1] H. Allouba, W. Zheng. Brownian-time processes: The PDE connection and the half-derivative generator. *Ann. Prob.* **29** (2009), 1780–1795. MR1880242
- [2] D. Applebaum. *Lévy Processes and Stochastic Calculus*. Second Edition, Cambridge University Press, New York, 2009. MR2512800
- [3] B. Baeumer, M.M. Meerschaert. Stochastic solutions for fractional Cauchy problems. *Fractional Calculus and Applied Analysis* **4** (2001), 481–500. MR1874479
- [4] B. Baeumer, M.M. Meerschaert, E. Nane. Brownian subordinators and fractional Cauchy problems. *Trans. Amer. Math. Soc.* **361** (2009), 3915–3930. MR2491905

- [5] B. Baeumer, M.M. Meerschaert. Tempered stable Lévy motion and transient super-diffusion. *J. Comput. Appl. Math.* **233** (2010), 2438–2448. MR2577834
- [6] L. Beghin, E. Orsingher. Fractional Poisson processes and related random motions. *Electron. J. Probab.*, **14** (2009), 1790–1826. MR2535014
- [7] L. Beghin, E. Orsingher. Poisson-type processes governed by fractional and higher-order recursive differential equations. *Electron. J. Probab.*, **15** (2010), 684–709. MR2650778
- [8] P. Billingsley. *Convergence of Probability Measures*. John Wiley, New York, 1968. MR0233396
- [9] N.H. Bingham. Limit theorems for occupation times of Markov processes. *Z. Wahrsch. Verw. Gebiete* **17** (1971), 1–22. MR0281255
- [10] N.H. Bingham. Maxima of sums of random variables and suprema of stable processes. *Z. Wahrsch. Verw. Gebiete* **26** (1973), 273–296. MR0415780
- [11] D.O. Cahoy. Fractional Poisson process in terms of alpha-stable densities. Doctoral Thesis, Case Western Reserve University. 2007. 106 pp. MR2710115
- [12] D.O. Cahoy, V.V. Uchaikin, W.A. Woyczynski. Parameter estimation for fractional Poisson processes. *J. Statist. Plann. Inf.* **140** (2010), 3106–3120.
- [13] M. Caputo. Linear models of dissipation whose Q is almost frequency independent, Part II. *Geophys. J. R. Astr. Soc.* **13** (1967), 529–539.
- [14] A. Cartea, D. Del-Castillo-Negrete. Fluid limit of the continuous-time random walk with general Lévy jump distribution functions. *Phys. Rev. E* **76** (2007), 041105.
- [15] A. Chakrabarty, M.M. Meerschaert. Tempered stable laws as random walk limits. *Statist. Probab. Lett.* **81**(8) (2011), 989–997.
- [16] A.V. Chechkin, R. Gorenflo, I.M. Sokolov. Retarding subdiffusion and accelerating superdiffusion governed by distributed-order fractional diffusion equations. *Phys. Rev. E* **66** (2002), 046129–046135.
- [17] A.V. Chechkin, J. Klafter, I.M. Sokolov. Fractional Fokker-Plank equation for ultraslow kinetics. *Europhys. Lett.* **63** (2003), 326–332.
- [18] R. Cont, P. Tankov. *Financial modelling with jump processes*. Chapman & Hall/CRC, Boca Raton, Florida, 2004. MR2042661
- [19] W. Feller. *An Introduction to Probability Theory and Its Applications*. Vol. II, 2nd Ed., Wiley, New York, 1971. MR0270403
- [20] D. Fulger, E. Scalas, G. Germano. Monte Carlo simulation of uncoupled continuous-time random walks yielding a stochastic solution of the space-time fractional diffusion equation. *Phys Rev E* **77** (2008), 021122.
- [21] R. Gorenflo, F. Mainardi. Fractional Calculus: Integral and Differential Equations of Fractional Order *Fractals and Fractional Calculus in Continuum Mechanics*. A. Carpinteri and F. Mainardi (Eds), Springer Verlag, New York, (1997), 223–276. MR1611585

- [22] R. Hilfer, L. Anton. Fractional master equations and fractal time random walks, *Phys. Rev. E* **51** (1995), R848-ÚR851.
- [23] E. Hille, R.S. Phillips. *Functional Analysis and Semi-Groups*. Amer. Math. Soc. Coll. Publ. **31**, American Mathematical Society, Providence, 1957. MR0089373
- [24] T. Huillet. On Linnik's continuous-time random walk. *J. Phys. A* **33** (2000), 2631–2652. MR1761635
- [25] G. Jumarie. Fractional master equation: non-standard analysis and Liouville-Riemann derivative. *Chaos Solitons Fractals*, **12** (2001), 2577-Ú2587.
- [26] V. Keyantuo,C. Lizama. On a connection between powers of operators and fractional Cauchy problems. (2009). Preprint available at [netlizama.usach.cl/Keyantuo-Lizama\(AMPA\)\(2009\).PDF](http://netlizama.usach.cl/Keyantuo-Lizama(AMPA)(2009).PDF).
- [27] A.N. Kochubei. Distributed order calculus and equations of ultraslow diffusion. *J. Math. Anal. Appl.* **340** (2008), 252–281. MR2376152
- [28] T.J. Kozubowski. The inner characterization of geometric stable laws. *Statist. Decisions* **12** (1994), 307–321. MR1309670
- [29] A.N. Lageras. A renewal-process-type expression for the moments of inverse subordinators. *J. Appl. Probab.*, **42** (2005), 1134–1144. MR2203828
- [30] N. Laskin. Fractional Poisson process. *Commun. Nonlinear Sci. Numer. Simul.*, **8** (2003), 201-Ú213.
- [31] F. Mainardi. The fundamental solutions for the fractional diffusion-wave equation. *Appl. Math. Lett.* **9** (1996), 23–28. MR1419811
- [32] F. Mainardi, R. Gorenflo, E. Scalas. A fractional generalization of the Poisson processes. *Vietnam Journ. Math.* **32** (2004), 53–64. MR2120631
- [33] F. Mainardi, R. Gorenflo,A. Vivoli. Beyond the Poisson renewal process: A tutorial survey. *J. Comput. Appl. Math.*, **205** (2007), 725–735. MR2329648
- [34] M.M. Meerschaert, H.P. Scheffler. *Limit Distributions for Sums of Independent Random Vectors: Heavy Tails in Theory and Practice*. Wiley Interscience, New York, 2001. MR1840531
- [35] M.M. Meerschaert, H.P. Scheffler. Limit theorems for continuous time random walks with infinite mean waiting times. *J. Appl. Probab.* **41** (2004), 623–638. MR2074812
- [36] M.M. Meerschaert, H.P. Scheffler. Stochastic model for ultraslow diffusion. *Stochastic Process Appl.* **116** (2006), 1215–1235. MR2251542
- [37] M.M. Meerschaert, H.P. Scheffler. Triangular array limits for continuous time random walks. *Stochastic Process Appl.* **118** (2008), 1606–1633. MR2442372
- [38] M.M. Meerschaert, Y. Zhang, B. Baeumer. Tempered anomalous diffusion in heterogeneous systems. *Geophys. Res. Lett.* **35** (2008), L17403.

- [39] M.M. Meerschaert, E. Nane, P Vellaisamy. Fractional Cauchy problems on bounded domains. *Ann. Probab.* **37** (2009), 979–1007. MR2537547
- [40] E. Nane. Stochastic solutions of a class of higher order Cauchy problems in \mathbb{R}^d . *Stochastics and Dynamics* **10** (2010), 341–366. MR2671380
- [41] R.N. Pillai. On Mittag-Leffer functions and related distributions. *Ann. Inst. Statist. Math.* **42** (1990), 157–161. MR1054728
- [42] O.N. Repin, A.I. Saichev. Fractional Poisson law. *Radiophys. and Quantum Electronics*, **43** (2000), 738–741.
- [43] J. Rosiński. Tempering stable processes. *Stochastic Process Appl.* **117** (2007), 677–707. MR2327834
- [44] H. Scher, M. Lax. Stochastic transport in a disordered solid. I. Theory. *Phys. Rev. B* **7** (1973), 4491–4502. MR0391854
- [45] E. Seneta. *Regularly Varying Functions*. Lecture Notes in Mathematics **508** (1976), Springer-Verlag, Berlin. MR0453936
- [46] V.V. Uchaikin, D.O. Cahoy, R.T. Sibatov. Fractional processes: from Poisson to branching one. *Internat. J. Bifur. Chaos Appl. Sci. Engrg.*, **18** (2008), 2717–2725. MR2479327
- [47] W. Whitt. *Stochastic-Process Limits*. Springer, New York, 2002. MR1876437
- [48] G. Zaslavsky. Fractional kinetic equation for Hamiltonian chaos. Chaotic advection, tracer dynamics and turbulent dispersion. *Phys. D* **76** (1994), 110–122. MR1295881