

## Infinite dimensional forward-backward stochastic differential equations and the KPZ equation\*

Sergio A. Almada Monter<sup>†</sup>      Amarjit Budhiraja<sup>‡</sup>

### Abstract

Kardar-Parisi-Zhang (KPZ) equation is a quasilinear stochastic partial differential equation (SPDE) driven by a space-time white noise. In recent years there have been several works directed towards giving a rigorous meaning to a solution of this equation. Bertini, Cancrini and Giacomin [2, 3] have proposed a notion of a solution through a limiting procedure and a certain renormalization of the nonlinearity. In this work we study connections between the KPZ equation and certain infinite dimensional forward-backward stochastic differential equations. Forward-backward equations with a finite dimensional noise have been studied extensively, mainly motivated by problems in mathematical finance. Equations considered here differ from the classical works in that, in addition to having an infinite dimensional driving noise, the associated SPDE involves a non-Lipschitz (specifically, a quadratic) function of the gradient. Existence and uniqueness of solutions of such infinite dimensional forward-backward equations is established and the terminal values of the solutions are then used to give a new probabilistic representation for the solution of the KPZ equation.

**Keywords:** Kardar-Parisi-Zhang equation; infinite dimensional noise; Backward stochastic differential equations; nonlinear stochastic partial differential equations; probabilistic representations.

**AMS MSC 2010:** Primary 60H15; 60H30, Secondary 35R60.

Submitted to EJP on March 29, 2013, final version accepted on March 19, 2014.

## 1 Introduction

In [16] probabilistic representations for solutions of certain quasilinear stochastic partial differential equations (SPDE) in terms of finite dimensional forward-backward stochastic differential equations have been studied. The driving noise in the SPDE of [16] is a finite dimensional Brownian motion. The paper shows that such representations can be used to prove existence and uniqueness results for the associated quasilinear equation. In recent years there have been many works that have established similar probabilistic representations for more general partial differential equations; see for

---

\*Research supported in part by the National Science Foundation (DMS-1004418, DMS-1016441), the Army Research Office (W911NF-10-1-0158) and the US-Israel Binational Science Foundation (Grant 2008466).

<sup>†</sup>Department of Statistics and Operations Research, University of North Carolina, Chapel Hill, USA.  
E-mail: salmada3@email.unc.edu

<sup>‡</sup>Department of Statistics and Operations Research, University of North Carolina, Chapel Hill, USA.  
E-mail: budhiraj@email.unc.edu

example [23] for such results for fully nonlinear equations. Backward stochastic differential equations have a long history of applications in financial mathematics; see [13] for a survey of the field; also see [12] or [17] for some modern applications. They have also motivated numerical methods for nonlinear partial differential equations; see for example [6] or [21].

In this work we are interested in a quasilinear SPDE driven by a space-time white noise of the following form

$$\partial_t h(t, x) = \partial_x^2 h(t, x)/2 - (\partial_x h(t, x))^2/2 + \mathscr{W}_t(x), \quad t \in [0, T], \quad x \in \mathbb{R}. \quad (1.1)$$

Here  $\mathscr{W}$  is the formal white noise field:

$$\mathbf{E} \mathscr{W}_t(x) \mathscr{W}_s(y) = \delta(t - s) \delta(x - y),$$

where  $\delta$  is the Dirac delta function. The above equation, known as the Kardar-Parisi-Zhang (KPZ) equation has been proposed in [10] to describe the long scale behavior of interface fluctuations in certain random polymer growth models. The solution  $h(t, x) \equiv h_t(x)$  represents the height of the interface at time  $t$  and location  $x$ . The equation (1.1) is ill posed in that, due to lack of spatial regularity of the noise, differentiable solutions do not exist and the nonlinear term on the right side of the equation does not allow a weak sense formulation of a solution. In [2] and [3] an interpretation of a solution of (1.1) is proposed through a limiting procedure and a certain ‘Wick renormalization’ of the nonlinear term. The Bertini-Cancrini-Giacomin (BCG) solution of (1.1) was shown in [3] to arise as a scaling limit of the fluctuation field for a microscopic interface model known as the ‘weakly asymmetric single step solid on solid process’. In recent years there have been several interesting papers that have studied scaling limits for similar particle models characterizing them through equations of the form (1.1); see for example [1], [7], [4] and references therein. In a different direction, a recent paper [8] has proposed a notion of a solution of (1.1) using rough path theory.

In this work we give a probabilistic representation of the BCG solution of (1.1) through solutions of certain infinite dimensional forward-backward stochastic differential equations. Although it is not clear whether the formulation of a solution to (1.1) given in this work enables one to prove new scaling limit theorems for stochastic particle systems, we believe that the probabilistic representations obtained here are natural for the study of nonlinear SPDE of the form (1.1). They can be regarded as extensions of classical Feynman-Kac formulae for solutions of linear SPDE [14, 15, 19]. One use of such probabilistic representations is in proving existence and uniqueness of solutions of quasilinear SPDE. Indeed the classical work of Pardoux and Peng [16] proves well-posedness of certain nonlinear SPDE driven by a finite dimensional noise using probabilistic representations of the form (1.1). Furthermore, such representations can be used to obtain numerical schemes to solve nonlinear equations. There is an extensive literature (cf. [13, 5] and references therein) that takes probabilistic representations for solutions of partial differential equations as a starting point to develop numerical schemes and to show their convergence properties. The current work extends the models studied in [16] to a setting where the driving noise is infinite dimensional. Another significant difference from the setting considered in [16] is that the equation (1.1) involves a quadratic function of the gradient while [16] considers the case of a Lipschitz non-linearity. Quadratic non-linearity has been studied by several authors in the setting of finite dimensional backward stochastic differential equations; see for example [11], [22]. However, none of these works treat equations involving both forward and backward stochastic integrals or the setting of an infinite dimensional noise.

A precise description of the representation obtained in this work requires some mathematical notation and background, which is given in Section 2, however below

we proceed formally in order to describe the basic idea. Denote by  $G_t$  the standard heat kernel on  $\mathbb{R} \times \mathbb{R}$ . Then a solution of (1.1) can formally be expressed as

$$h(S, x) = G_S \star h_0(x) - \frac{1}{2} \int_0^S G_{S-r} \star (\partial_x h(r, \cdot))^2(x) dr + \int_{\mathbb{R} \times [0, S]} G_{S-r}(x, y) \mathscr{W}(dy, dr), \quad (1.2)$$

for  $S \in [0, T]$  and  $x \in \mathbb{R}$ , and where  $\star$  denotes the convolution in space and  $h_0(x) = h(0, x)$  is the initial condition for (1.1). Fix  $S \in [0, T]$  and  $x \in \mathbb{R}$ . Let  $W$  be a standard Brownian Motion independent of  $\mathscr{W}$  and  $X_r^S(x) = x + W(S) - W(r)$ , for  $r \in [0, S]$ . Then one can rewrite the expression in the above display as

$$h(S, x) = \mathbf{E} \left[ h_0(X_0^S(x)) - \frac{1}{2} \int_0^S (\partial_x h(r, X_r^S(x)))^2 dr + \int_0^S \mathscr{W}(X_r^S(x), dr) \middle| \mathcal{F}^{\mathscr{W}} \right], \quad (1.3)$$

where  $\mathcal{F}^{\mathscr{W}}$  denotes the  $\sigma$ -field generated by  $\mathscr{W}$ . The stochastic integral on the right side above is of course entirely formal (as is much of this description). Define stochastic processes

$$z_S(r, x) = \partial_x h(r, X_r^S(x)), \quad y_S(r, x) = h(r, X_r^S(x)), \quad r \in [0, S]. \quad (1.4)$$

Note that the values of these processes at time  $r$  depend on the past values (i.e. values over  $[0, r]$ ) of  $\mathscr{W}$  and the future increments (those over  $[r, S]$ ) of  $W$ . Also note that  $y_S(S, x) = h(S, x)$ . Let, for  $0 \leq r \leq S$ ,  $\mathcal{F}_r^{\mathscr{W}} \vee \mathcal{F}_{r,S}^W$  be the  $\sigma$ -field generated by  $\{\mathscr{W}(s, x), s \leq r, x \in \mathbb{R}\}$  and  $\{W(S) - W(u), 0 \leq r \leq u \leq S\}$ . Then (1.3) can be written as

$$\begin{aligned} h(S, x) &= y_S(S, x) \\ &= \mathbf{E} \left[ h_0(X_0^S(x)) - \frac{1}{2} \int_0^S z_S(r, x)^2 dr + \int_0^S \mathscr{W}(X_r^S(x), dr) \middle| \mathcal{F}_S^{\mathscr{W}} \vee \mathcal{F}_{S,S}^W \right]. \end{aligned}$$

The above formula suggests an evolution equation for  $y_S(u, x)$ ,  $0 \leq u \leq S$  of the following form

$$y_S(u, x) = h_0(X_0^S(x)) - \frac{1}{2} \int_0^u z_S(r, x)^2 dr + \int_0^u \mathscr{W}(X_r^S(x), dr) + M_S(u, x), \quad 0 \leq u \leq S, \quad (1.5)$$

where the process  $M_S(u, x)$ ,  $u \in [0, S]$ , is such that

$$\mathbf{E} [M_S(u, x) \mid \mathcal{F}_u^{\mathscr{W}} \vee \mathcal{F}_{u,S}^W] = 0,$$

Thus if one can make (1.5) rigorous, one can then obtain a solution  $h(S, x)$  of (1.1) by evaluating the solution of (1.5) at  $u = S$ . The goal of this work is to show that after a suitable mollification of the infinite dimensional noise, the above equation can indeed be interpreted in a rigorous manner. The stochastic process  $M_S(u, x)$  in the mollified equation (see (2.5)) is given as a backward stochastic integral with respect to  $W$ . Our main result (Theorem 2.2) says that there is a unique pair of processes  $(z_S(\cdot, x), y_S(\cdot, x))$  (in a suitable class) that satisfy equation (2.5). This equation is a forward-backward SDE with a quadratic nonlinearity on the right side. The non-Lipschitz feature of the nonlinearity makes the uniqueness proof somewhat challenging. Our proof relies on certain truncation arguments along with Tanaka's formula and properties of local times of backward semimartingales. Denoting the solution of (2.5) as  $(y_S^k, z_S^k)$ , where  $k \in \mathbb{N}$  is a mollification parameter, we also give a Feynman-Kac formula for  $y_S^k(t, x)$  that involves the two noise processes  $W$  and  $\mathscr{W}$  (see Lemma 3.1). When  $t = S$ , this formula after some simplification reduces to the formula for the logarithm of the solution of the mollified stochastic heat equation given in [2] (see (2.17) in [2] and also the proof of

Theorem 2.2 in the current work). More generally, for  $0 \leq t \leq S$ , we have that  $y_S^k(t, x) = -\log \psi^k(t, x + W(S) - W(t))$ , where  $\psi^k$  is the solution of the mollified stochastic heat equation with the mollification parameter  $k$  (see (4.2)). Let  $\mathcal{S}_T = \{(t, S) : 0 \leq t \leq S \leq T\}$ . Then, letting  $y^k = \{y_S^k(t, x) : (t, S) \in \mathcal{S}_T, x \in \mathbb{R}\}$  we have that  $\{y^k\}_{k \geq 1}$  is a sequence of random variables with values in  $C(\mathcal{S}_T \times \mathbb{R} : \mathbb{R})$  (the space of continuous functions from  $\mathcal{S}_T \times \mathbb{R}$  to  $\mathbb{R}$  equipped with the usual topology of local uniform convergence). It then follows, using Theorem 2.1 of [3], that  $y^k$  converges in distribution to a limit  $y$  and, letting  $h(S, x) = y_S(S, x)$ ,  $S \in [0, T]$ ,  $h$  is the solution of (1.1) (see Theorem 2.3). In contrast to the Feynman-Kac formula, equation (2.5) gives a stochastic differential equation which can in principle be (numerically) solved in a dynamic fashion to yield an approximation for the solution of the KPZ equation.

In the next section, we give a precise formulation and present our main results.

## 2 Mathematical Preliminaries and Main Results.

In order to state our precise representation for the BCG solution of (1.1) we need some notation. Let  $\mathbf{H} = L^2(\mathbb{R}, dx)$ , i.e. the Hilbert space of square integrable (with respect to the Lebesgue measure) functions on the real line. We will denote the inner product and the norm on  $\mathbf{H}$  by  $\langle \cdot, \cdot \rangle$  and  $\| \cdot \|$ , respectively. Let  $(\Omega, \mathcal{F}, \mathbf{P})$  be a complete probability space on which is given a collection of continuous real stochastic processes  $\{B_t(h); t \geq 0\}_{h \in \mathbf{H}}$  that defines a cylindrical Brownian motion (c.B.m) on  $\mathbf{H}$ . Namely,

- $B_t(0) = 0$  and for each nonzero  $h \in \mathbf{H}$ ,  $B_t(h)\langle h, h \rangle^{-1/2}$  is a one dimensional standard Wiener process.
- For each  $h \in \mathbf{H}$ ,  $\{B_t(h)\}_{t \geq 0}$  is a  $\mathcal{F}_t^B$  martingale, where  $\mathcal{F}_t^B = \sigma\{B_s(v) : s \leq t, v \in \mathbf{H}\} \vee \mathcal{N}$  and  $\mathcal{N}$  is the collection of all  $\mathbf{P}$  null sets.

Next, following [3], we consider a regularized version of (1.1). Let  $\zeta \in C_0^\infty(\mathbb{R})$  [space of smooth functions on the real line with compact support] be a nonnegative even function such that  $\int_{\mathbb{R}} \zeta(x) dx = 1$ . For  $k \in \mathbb{N}$ , let  $\zeta^k(y) = k\zeta(ky)$ ,  $y \in \mathbb{R}$ . For  $x \in \mathbb{R}$ , define  $\zeta_x^k \in C_0^\infty(\mathbb{R})$  as  $\zeta_x^k(y) = \zeta^k(x - y)$ ,  $y \in \mathbb{R}$ . Consider the Gaussian random field

$$B^k(t, x) = B_t(\zeta_x^k), \quad t \geq 0, x \in \mathbb{R}$$

with covariance

$$\mathbf{E}B^k(t, x)B^{k'}(s, y) = (t \wedge s)C^k(x - y), \quad x, y \in \mathbb{R}, t, s \in [0, \infty),$$

where

$$C^k(x) = \zeta^k \star \zeta^k(x) \equiv \int_{\mathbb{R}} \zeta_x^k(y)\zeta^k(y)dy.$$

Note that  $C^k(0) = k\|\zeta\|^2$ .

The mollified KPZ equation (see [3]) is given as follows.

$$h^k(t, x) = h_0(x) + \frac{1}{2} \int_0^t (\partial_x^2 h^k(s, x) - ((\partial_x h^k(s, x))^2 - C^k(0))) ds + B^k(t, x). \quad (2.1)$$

The initial condition  $h_0$  is a  $C(\mathbb{R})$  valued random variable, independent of  $B$ , satisfying the following integrability condition

$$\text{for every } p > 0 \text{ there exist } a_p > 0 \text{ such that } \sup_{x \in \mathbb{R}} e^{-a_p|x|} \mathbf{E}e^{p|h_0(x)|} \equiv b_p < \infty. \quad (2.2)$$

The hypothesis imposed on the initial condition in (2.4) of [3] is weaker than the integrability condition in (2.2), but the condition we impose covers all classical cases, in

particular the combinations of the so called Brownian and Flat geometries (see last column of Table 5 in [4]). Note also that the assumption in (2.2) (and also condition (2.4) of [3]) exclude the model setting considered in [1], and [20] where the initial condition is a distribution.

Solution of (2.1) over any fixed time interval  $[0, T]$  is understood in the weak sense, namely, it is a  $\{\mathcal{F}_t^B\}$ - adapted stochastic process  $\{h^k(t, \cdot)\}_{0 \leq t \leq T} \equiv \{h_t^k\}_{0 \leq t \leq T}$  with sample paths in  $C([0, T] : C(\mathbb{R})) \cap C((0, T] : C^1(\mathbb{R}))$ , such that for every smooth function  $\varphi$  on  $\mathbb{R}$  with a compact support

$$h_t^k(\varphi) = h_0(\varphi) + \frac{1}{2} \int_0^t [h_s^k(\varphi'') - ((\partial_x h_s^k)^2 - C^k(0))(\varphi)] ds + B_t^k(\varphi)$$

where for a function  $g$  on  $\mathbb{R}$  (with suitable integrability properties), and a smooth function  $\varphi$ ,  $g(\varphi) = \int_{\mathbb{R}} g(x)\varphi(x)dx$ . Here  $C(\mathbb{R})$  [resp.  $C^1(\mathbb{R})$ ] is the space of continuous [resp. continuously differentiable] functions on the real line.

The paper [3] shows that there is a unique solution of (2.1) in the class of processes that satisfy

$$\sup_{t \in [0, T], r \in \mathbb{R}} e^{-a|r|} \mathbf{E} \left[ e^{-2h_t^k(r)} \right] < \infty \text{ for some } a \in (0, \infty).$$

Furthermore, the paper [3] shows that as  $k \rightarrow \infty$ ,  $h^k$  converges in distribution (as a  $C([0, T] : C(\mathbb{R}))$  valued random variable) to a limit process  $h$ , which is *defined* to be the solution of (1.1). Throughout this work, this process (strictly speaking – its probability law on  $C([0, T] : C(\mathbb{R}))$ ) will be referred to as *the BCG solution of the KPZ equation*.

We will now introduce a forward - backward stochastic differential equation associated with (2.2). Assume, without loss of generality, that we are given on  $(\Omega, \mathcal{F}, \mathbf{P})$  another real standard Brownian motion  $W$  that is independent of  $(B, h_0)$ . For  $S > 0$  and  $0 \leq t \leq S$ , we denote

$$\mathcal{F}_{t,S}^W = \sigma\{W(s) - W(t) : s \in [t, S]\} \vee \mathcal{N}, \text{ and } \mathcal{F}_t^S = \mathcal{F}_{t,S}^W \vee \mathcal{F}_t^B \vee \sigma\{h_0\}.$$

Note that  $\mathcal{F}^S \equiv \{\mathcal{F}_t^S : t \in [0, S]\}$  is not a filtration since the  $\sigma$ -fields in this collection are neither increasing nor decreasing in  $t$ . However, abusing terminology, we will say a stochastic process  $\{V_t\}_{t \in [0, S]}$  is  $\mathcal{F}^S$  adapted if  $V_t$  is  $\mathcal{F}_t^S$  measurable for every  $t \in [0, S]$ .

Throughout we will fix a complete orthonormal system  $\{\gamma_m\}_{m \in \mathbb{N}}$  in  $\mathbf{H}$  and denote  $B(\gamma_m) = \beta_m$ . Note that  $\{\beta_m\}_{m \in \mathbb{N}}$  is a sequence of independent standard Brownian motions, independent of  $W$ . For a  $\mathcal{F}^S$  adapted  $\mathbf{H}$ -valued process  $\{\varphi(t)\}_{0 \leq t \leq S}$  satisfying  $\mathbf{E} \int_0^S \|\varphi(t)\|^2 dt < \infty$ , the Itô stochastic integral  $\int_0^t \langle \varphi(r), dB_r \rangle$  for  $t \in [0, S]$  is well defined and is given as

$$\int_0^t \langle \varphi(r), dB_r \rangle = \sum_{m \in \mathbb{N}} \int_0^t \langle \varphi(r), \gamma_m \rangle d\beta_m(r),$$

where the series on the right converges in  $L^2(\mathbf{P})$ .

For a family of sigma algebras  $\{\mathcal{G}_t; 0 \leq t \leq S\}$ , let  $\mathcal{H}_S^p(\mathcal{G})$  [ resp.  $\mathcal{H}_S^\infty(\mathcal{G})$ ] be the space of measurable [resp. continuous] processes  $\{\phi(t) : t \in [0, S]\}$  such that  $\phi(t)$  is  $\mathcal{G}_t$  measurable for every  $t$ , and

$$\mathbf{E} \int_0^S |\phi(t)|^p dt < \infty \text{ [ resp. } \mathbf{E} \sup_{t \in [0, S]} |\phi(t)|^2 < \infty.]$$

For  $H \in \mathcal{H}_S^2(\mathcal{F}^S)$ , we denote the backward stochastic integral of  $H$  with respect to  $W$  by  $\int_t^S H(r) \downarrow dW$ . See Appendix for a brief review of such stochastic integrals.

Let  $X_t^S(x) \equiv X^S(t, x) = x + W(S) - W(t)$ , for  $0 \leq t \leq S$  and  $x \in \mathbb{R}$ . Define

$$Z^k(t, x) = \int_0^t \langle \zeta_{X^S(r,x)}^k, dB_r \rangle, \quad (t, x) \in [0, S] \times \mathbb{R}. \quad (2.3)$$

Note that  $\|\zeta_x^k\|^2 = \|\zeta_0^k\|^2 = C^k(0)$  for all  $x$  and consequently

$$\int_0^t \|\zeta_{X^S(r,x)}^k\|^2 dr = tC^k(0) \text{ for all } t \in [0, S]. \quad (2.4)$$

Also,  $\{\zeta_{X^S(t,x)}^k\}_{0 \leq t \leq S}$  is  $\mathcal{F}^S$  adapted and so the stochastic integral in (2.3) is well defined and has the representation

$$Z^k(t, x) = \sum_{m \in \mathbb{N}} \int_0^t \langle \zeta_{X^S(r,x)}^k, \gamma_m \rangle d\beta_m(r)$$

with the series converging in  $L^2(\mathbf{P})$ . We now consider the following doubly backward SDE

$$y_S^k(t, x) = h_0(X_0^S(x)) - \frac{1}{2} \int_0^t (z_S^k(r, x)^2 - C^k(0)) dr + Z^k(t, x) - \int_0^t z_S^k(r, x) \downarrow dW(r). \quad (2.5)$$

**Definition 2.1.** We say the collection  $\{y_S^k(t, x), z_S^k(t, x), 0 \leq t \leq S, x \in \mathbb{R}\}$  of real random variables is a solution of (2.5) if for every  $x \in \mathbb{R}$ ,  $(y_S^k(\cdot, x), z_S^k(\cdot, x)) \in \mathcal{H}_S^\infty(\mathcal{F}^S) \times \mathcal{H}_S^2(\mathcal{F}^S)$  and the equation is satisfied for all  $(t, x) \in [0, S] \times \mathbb{R}$ , almost surely. We say that the equation has a unique solution if  $\{\tilde{y}_S^k(t, x), \tilde{z}_S^k(t, x), 0 \leq t \leq S, x \in \mathbb{R}\}$  is another such collection then  $(y_S^k(t, x), z_S^k(t, x)) = (\tilde{y}_S^k(t, x), \tilde{z}_S^k(t, x))$  a.s., for all  $(t, x) \in [0, S] \times \mathbb{R}$ .

Frequently, when clear from the context, we will suppress  $x$  and denote the solution as  $\{y_S^k(t), z_S^k(t)\}_{0 \leq t \leq S}$  or merely as  $(y_S^k, z_S^k)$ . The first result in this work establishes wellposedness of the above equation.

**Theorem 2.2.** Fix  $k \in \mathbb{N}$  and  $S > 0$ . Suppose that  $h_0$  is a  $C(\mathbb{R})$  valued random variable, independent of  $(B, W)$ , satisfying (2.2). Then there exists a unique solution to (2.5).

Our second result concerns the asymptotic behavior of  $y_S^k$ , as  $k \rightarrow \infty$ , and relation with the KPZ equation. Let  $y^k = \{y_S^k(t, x) : (t, S) \in \mathcal{S}_T, x \in \mathbb{R}\}$ . Then  $\{y^k\}_{k \geq 1}$  is a sequence of random variables with values in  $C(\mathcal{S}_T \times \mathbb{R} : \mathbb{R})$  (the space of continuous functions from  $\mathcal{S}_T \times \mathbb{R}$  to  $\mathbb{R}$  equipped with the usual topology of local uniform convergence).

**Theorem 2.3.** Fix  $x \in \mathbb{R}$  and  $k \in \mathbb{N}$ . Let  $h_0$  be as in Theorem 2.2 and, for  $k \geq 1$  and  $S > 0$ ,  $(y_S^k, z_S^k)$  be as obtained from Theorem 2.2. Then, the sequence  $y^k$  converges in distribution as  $k \rightarrow \infty$  to a  $C(\mathcal{S}_T \times \mathbb{R} : \mathbb{R})$  valued random variable  $y \equiv \{y_S(t, x), (t, S) \in \mathcal{S}_T, x \in \mathbb{R}\}$ . Furthermore, letting  $h(t, x) = y_t(t, x)$ ,  $(t, x) \in [0, T] \times \mathbb{R}$ ,  $h$  is a solution of (1.1).

We remark that from the formal discussion in the introduction (see (1.4)) one expects that there is a modification of the random field  $(y_S^k, z_S^k)$  (denoted once more by the same symbols) such that  $y_S^k$  is differentiable in  $x$ , the derivative is jointly continuous in  $(t, x)$  and  $\partial_x y_S^k(t, x) = z_S^k(t, x)$ . Although, this does not follow from the proof of Theorem 2.2, one can see from the arguments in the proof of Theorem 2.3 that in fact such a modification is available. In fact as (4.3) shows that such a modification of  $y_S^k$  can be obtained by solving a regularized stochastic heat equation.

The rest of this work is devoted to the proof of the above two results.

### 3 Proof of Theorem 2.2.

Fix  $k \in \mathbb{N}$  and  $S > 0$ . Suppose that  $(y_S^k, z_S^k)$  solves equation (2.5). Define

$$u_S^k(t, x) = \exp(-y_S^k(t, x)), \quad v_S^k(t, x) = u_S^k(t, x)z_S^k(t, x), \quad 0 \leq t \leq S, \quad x \in \mathbb{R}. \quad (3.1)$$

A formal application of Itô's formula (Lemma 5.4) yields the following equation for  $(u_S^k, v_S^k)$ .

$$u_S^k(t, x) = u_0(X^S(0, x)) - \int_0^t u_S^k(r, x) dZ^k(r, x) + \int_0^t v_S^k(r, x) \downarrow dW(r), \quad (3.2)$$

where  $u_0(x) = \exp(-h_0(x))$ . The transformation in (3.1) thus motivates the study of equation (3.2), and as a first step we will now establish the wellposedness of (3.2). Namely, we first prove the existence and uniqueness of a pair  $(u_S^k, v_S^k)$ , with appropriate integrability and measurability properties, which satisfies (3.2). Note that the integrals on the right side of (3.2) are well defined if  $(u_S^k, v_S^k) \in \mathcal{H}_S^\infty(\mathcal{F}^S) \times \mathcal{H}_S^2(\mathcal{F}^S)$ .

**Lemma 3.1.** *Fix  $x \in \mathbb{R}, k \in \mathbb{N}$  and  $S > 0$ . Then there is a unique pair  $(u_S^k, v_S^k) \in \mathcal{H}_S^\infty(\mathcal{F}^S) \times \mathcal{H}_S^2(\mathcal{F}^S)$  that satisfies equation (3.2). Furthermore,*

$$u_S^k(t, x) = \mathbf{E} \left[ u_0(X^S(0, x)) \exp \left\{ -Z^k(t, x) - \frac{1}{2} C^k(0)t \right\} \middle| \mathcal{F}_t^S \right]. \quad (3.3)$$

Finally, for any  $p \geq 2$ , there is a  $C(p, k) \in (0, \infty)$  such that, for all  $x \in \mathbb{R}$ ,

$$\mathbf{E} \sup_{t \leq S} u_S^k(t, x)^p + \mathbf{E} \left( \int_0^t v_S^k(r, x)^2 dr \right)^{p/2} \leq C(p, k)(1 + \mathbf{E} u_0(X^S(0, x))^{4p}). \quad (3.4)$$

**Proof.** Since  $x, k$  and  $S$  are fixed, we omit them from the notation throughout this proof. In particular we write  $Z^k(t, x)$  and  $X^S(t, x)$  as  $Z(t)$  and  $X(t)$  respectively. For a stochastic process  $H = \{H(t)\}_{0 \leq t \leq S}$ , we define its time reversed path

$$\tilde{H}(t) = H(S - t) - H(S), \quad 0 \leq t \leq S.$$

In particular,

$$\tilde{B}_t(f) = B_{S-t}(f) - B_S(f), \quad f \in \mathbf{H}, \quad \text{and} \quad \tilde{W}(t) = W(S - t) - W(S).$$

Define  $\hat{X}(r) = x - \tilde{W}(r)$ . Then note that

$$\begin{aligned} \tilde{Z}(t) &= Z(S - t) - Z(S) = - \int_{S-t}^S \langle \zeta_{X(r)}^k, dB_r \rangle \\ &= - \sum_m \int_{S-t}^S \langle \zeta_{X(r)}^k, \gamma_m \rangle d\beta_m(r) = \sum_m \int_0^t \langle \zeta_{\hat{X}(r)}^k, \gamma_m \rangle d\tilde{\beta}_m(r) \\ &= \int_0^t \langle \zeta_{\hat{X}(r)}^k, d\tilde{B}_r \rangle, \end{aligned} \quad (3.5)$$

where the next to last equality follows on noting that

$$\hat{X}(S - r) = x - \tilde{W}(S - r) = x - W(r) + W(S) = X(r).$$

Let, for  $0 \leq t \leq s \leq S$ ,

$$\mathcal{F}_{t,s}^{\tilde{B}} = \sigma\{\tilde{B}_s(v) - \tilde{B}_u(v) : u \in [t, s], v \in \mathbf{H}\} \vee \mathcal{N}$$

and

$$\mathcal{F}_t^{\tilde{W}} = \sigma\{\tilde{W}(s) : 0 \leq s \leq t\} \vee \mathcal{N} = \mathcal{F}_{S-t}^W. \tag{3.6}$$

Note in particular that

$$\begin{aligned} \mathcal{F}_{t,S}^{\tilde{B}} &= \sigma\{\tilde{B}_S(v) - \tilde{B}_u(v) : u \in [t, S], v \in \mathbf{H}\} \vee \mathcal{N} \\ &= \sigma\{B_u(v) : u \in [0, S-t], v \in \mathbf{H}\} \vee \mathcal{N} = \mathcal{F}_{S-t}^B \end{aligned} \tag{3.7}$$

Also, let

$$\tilde{\mathcal{F}}_t^S = \mathcal{F}_t^{\tilde{W}} \vee \mathcal{F}_{t,S}^{\tilde{B}} \vee \sigma\{h_0\} = \mathcal{F}_{S-t}^W \vee \mathcal{F}_{S-t}^B \vee \sigma\{h_0\} = \mathcal{F}_{S-t}^S. \tag{3.8}$$

From Corollary 5.2 in the Appendix it follows that in order to prove the first statement of the lemma it suffices to show that there exists a unique pair  $(\hat{u}, \hat{v}) \in \mathcal{H}_S^\infty(\tilde{\mathcal{F}}^S) \times \mathcal{H}_S^2(\tilde{\mathcal{F}}^S)$  that solves the time reversed equation

$$\hat{u}(t) = u_0(\hat{X}(S)) + \int_t^S \hat{u}(r) \downarrow d\tilde{Z}(r) - \int_t^S \hat{v}(r) d\tilde{W}(r). \tag{3.9}$$

The unique solution  $(u, v)$  of (3.2) can then be obtained on taking  $(u(t), v(t)) = (\hat{u}(S-t), \hat{v}(S-t))$ . We now consider the unique solvability of (3.9). Let  $\mathcal{G}_t = \mathcal{F}_t^{\tilde{W}} \vee \mathcal{F}_{0,S}^{\tilde{B}} \vee \sigma\{h_0\}$ . Recalling that  $\|\zeta_x^k\|^2 = C_k(0)$  for all  $x \in \mathbb{R}$  and elementary properties of Brownian motions, we see that

$$\mathbf{E} \sup_{0 \leq t \leq S} \exp\{m|Z(t)|\} < \infty, \quad \mathbf{E} \exp\{m|\hat{X}(S)|\} < \infty, \quad \text{for all } m \in \mathbb{N}. \tag{3.10}$$

Combining this with the integrability condition (2.2) and an application of Cauchy-Schwarz inequality we have

$$\mathbf{E} \left[ u_0(\hat{X}(S)) \exp\{\tilde{Z}(S)\} \right]^2 < \infty.$$

Consequently,

$$M(t) = \mathbf{E} \left[ u_0(\hat{X}(S)) \exp\{\tilde{Z}(S) - \frac{1}{2}C^k(0)S\} \mid \mathcal{G}_t \right], \quad 0 \leq t \leq S, \tag{3.11}$$

is a square integrable  $\mathcal{G}_t$ -martingale. From a straightforward extension of the classical martingale representation theorem, there is a  $\mathcal{G}_t$ -progressively measurable process  $\{J(t); 0 \leq t \leq S\}$  such that  $\mathbf{E} \int_0^S J(r)^2 dt < \infty$ , and

$$M(t) = M(0) + \int_0^t J(r) d\tilde{W}(r), \quad 0 \leq t \leq S. \tag{3.12}$$

Define, for  $0 \leq t \leq S$ ,

$$E(t) = \exp\{-\tilde{Z}(t) + \frac{1}{2}C^k(0)t\}, \quad V(t) = E(t)J(t), \quad U(t) = E(t)M(t). \tag{3.13}$$

We now show that  $(U, V) \in \mathcal{H}_S^\infty(\tilde{\mathcal{F}}^S) \times \mathcal{H}_S^2(\tilde{\mathcal{F}}^S)$ . More precisely, for the process  $V$  we will show that there is a  $\tilde{V} \in \mathcal{H}_S^2(\tilde{\mathcal{F}}^S)$  such that  $\tilde{V}(t) = V(t)$  for a.e.  $t$ , a.s. Consider  $U$  first. Note that  $U$  can be rewritten as

$$\begin{aligned} U(t) &= E(t) \left( M(S) - \int_t^S J(r) d\tilde{W}(r) \right) \\ &= u_0(\hat{X}(S)) \frac{E(t)}{E(S)} - E(t) \int_t^S J(r) d\tilde{W}(r). \end{aligned} \tag{3.14}$$



From (3.13) and (3.5) it follows that  $U$  is  $\{\mathcal{G}_t\}$  adapted. Next, using the independence between  $\tilde{Z}(S) - \tilde{Z}(t)$  and  $\mathcal{F}_{0,t}^{\tilde{B}}$  and that  $\{M(t)\}$  is a  $\{\mathcal{G}_t\}$  martingale, we see from the above display that

$$\begin{aligned} U(t) &= \mathbf{E}(U(t) \mid \mathcal{G}_t) \\ &= \mathbf{E} \left[ u_0(\hat{X}(S)) \frac{E(t)}{E(S)} \mid \mathcal{G}_t \right] \end{aligned} \tag{3.15}$$

$$= \mathbf{E} \left[ u_0(\hat{X}(S)) \frac{E(t)}{E(S)} \mid \mathcal{F}_t^{\tilde{W}} \vee \mathcal{F}_{t,S}^{\tilde{B}} \vee \sigma\{h_0\} \right]. \tag{3.16}$$

Thus  $U$  is  $\tilde{\mathcal{F}}^S$  - adapted. We now argue that there is a  $\tilde{V}$  that is  $\tilde{\mathcal{F}}^S$  - adapted and such that  $V(t) = \tilde{V}(t)$ , for a.e.  $t$ , a.s. For  $c \in (0, \infty)$ , define  $J^c(r) = J(r)1_{|J(r)| \leq c}$ . Let, for  $\varepsilon > 0$

$$F_\varepsilon^c = \frac{1}{\sqrt{\varepsilon}} \int_t^{t+\varepsilon} J^c(r) d\tilde{W}(r).$$

By Itô's isometry, we have that

$$\mathbf{E} \left[ F_\varepsilon^c \frac{\tilde{W}(t+\varepsilon) - \tilde{W}(t)}{\sqrt{\varepsilon}} \mid \mathcal{G}_t \right] = \mathbf{E} \left[ \frac{1}{\varepsilon} \int_t^{t+\varepsilon} J^c(r) dr \mid \mathcal{G}_t \right].$$

Sending  $\varepsilon \rightarrow 0$ ,  $c \rightarrow \infty$ , and recalling that  $J(t)$  is  $\mathcal{G}_t$  measurable, we have that

$$\limsup_{c \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \mathbf{E} \left[ F_\varepsilon^c \frac{\tilde{W}(t+\varepsilon) - \tilde{W}(t)}{\sqrt{\varepsilon}} \mid \mathcal{G}_t \right] = J(t) = \frac{V(t)}{E(t)}, \text{ a.e. } t, \text{ a.s.}$$

and therefore since  $E(t)$  is  $\mathcal{G}_t$  measurable

$$V(t) = \limsup_{c \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \mathbf{E} \left[ E(t) F_\varepsilon^c \frac{\tilde{W}(t+\varepsilon) - \tilde{W}(t)}{\sqrt{\varepsilon}} \mid \mathcal{G}_t \right] \text{ a.e. } t, \text{ a.s.} \tag{3.17}$$

Define  $U^c$  by replacing  $J$  with  $J^c$  on the right side of (3.14). Then a calculation similar to the one leading to (3.16) shows that  $U^c$  is  $\tilde{\mathcal{F}}^S$  - adapted.

Also note that

$$\sqrt{\varepsilon} E(t) F_\varepsilon^c = \exp\{\tilde{Z}(t+\varepsilon) - \tilde{Z}(t) - \frac{1}{2} C^k(0)\varepsilon\} U^c(t+\varepsilon) - U^c(t)$$

and consequently  $\sqrt{\varepsilon} E(t) F_\varepsilon^c$  is independent of  $\mathcal{F}_{0,t}^{\tilde{B}}$ . Thus the right side of (3.17) equals

$$\limsup_{c \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \mathbf{E} \left[ E(t) F_\varepsilon^c \frac{\tilde{W}(t+\varepsilon) - \tilde{W}(t)}{\sqrt{\varepsilon}} \mid \tilde{\mathcal{F}}_t^S \right]$$

and so  $\tilde{V}$  defined by the right side of (3.17) is  $\tilde{\mathcal{F}}^S$  - adapted.

We now prove the stated integrability properties of  $(U, V)$ . From (3.15), for  $m \in \mathbb{N}$ ,

$$\mathbf{E} \sup_{0 \leq t \leq S} U(t)^m \leq \mathbf{E} \left[ \sup_{0 \leq t \leq S} \mathbf{E} \left[ \left( \frac{u_0(\hat{X}(S))^m}{E(S)^m} \sup_{0 \leq r \leq S} E(r)^m \right) \mid \mathcal{G}_t \right] \right]. \tag{3.18}$$

From (3.10) it follows that, for any  $m \geq 1$ ,

$$\mathbf{E} \left[ \frac{u_0(\hat{X}(S))^m}{E(S)^m} \sup_{0 \leq r \leq S} E(r)^m \right] < \infty.$$

Standard martingale inequalities now show that the right side of (3.18) is finite, indeed we have that, for every  $m \in \mathbb{N}$  there are  $C_1(m), C_2(m) \in (0, \infty)$  such that

$$\begin{aligned} \mathbf{E} \sup_{0 \leq t \leq S} U(t)^m &\leq C_1(m) \left( 1 + \mathbf{E} \left[ \frac{u_0(\hat{X}(S))^{2m}}{E(S)^{2m}} \sup_{0 \leq r \leq S} E(r)^{2m} \right] \right) \\ &\leq C_2(m) \left( 1 + \mathbf{E} u_0((\hat{X}(S))^{4m}) \right) \\ &< \infty. \end{aligned} \tag{3.19}$$

Next consider  $V$ . By classical martingale inequalities (cf. Proposition 3.3.26 of [9]), for every  $m \in \mathbb{N}$  there is a  $b_m \in (0, \infty)$  such that

$$\mathbf{E} \left( \int_0^S J(r)^2 dr \right)^m \leq b_m (\mathbf{E} M(S)^{2m} + \mathbf{E} M(0)^{2m}).$$

Thus, recalling the definition of  $\{M(t)\}$  (see (3.11) ) and using (3.10) once more, we have that for every  $m \in \mathbb{N}$  there is a  $C_3(m) \in (0, \infty)$ , such that

$$\begin{aligned} \mathbf{E} \left( \int_0^S J(r)^2 dr \right)^m &\leq C_3(m) \left( 1 + \mathbf{E}(u_0(\hat{X}(S))^{4m}) \right) \\ &< \infty. \end{aligned} \tag{3.20}$$

Next,

$$\begin{aligned} \mathbf{E} \left( \int_0^S V(r)^2 dr \right)^m &= \mathbf{E} \left( \int_0^S E(r)^2 J(r)^2 dr \right)^m \\ &\leq \mathbf{E} \left[ \left( \sup_{0 \leq r \leq S} E(r)^{2m} \right) \left( \int_0^S J(r)^2 dr \right)^m \right]. \end{aligned}$$

Finiteness of the last term is immediate from (3.10) and (3.20). In fact we have that for every  $m \in \mathbb{N}$  there is a  $C_4(m) \in (0, \infty)$ , such that

$$\mathbf{E} \left( \int_0^S V(r)^2 dr \right)^m \leq C_4(m) \left( 1 + \mathbf{E}(u_0(\hat{X}(S))^{8m}) \right). \tag{3.21}$$

Combining (3.19), (3.21) and the  $\tilde{\mathcal{F}}^S$  adaptedness of  $(U, V)$  we have in particular that  $(U, V) \in \mathcal{H}_S^\infty(\tilde{\mathcal{F}}^S) \times \mathcal{H}_S^2(\tilde{\mathcal{F}}^S)$ . By an application of Itô's formula, we now see that, for  $t \in [0, S]$

$$U(t) = U(S) + \int_t^S U(r) \downarrow d\tilde{Z}(r) - \int_t^S V(r) d\tilde{W}(r). \tag{3.22}$$

For completeness, we give a proof of the above equality in the Appendix.

Thus we have shown that  $(U, V)$  is a solution of (3.9) and therefore, as noted earlier

$$(u_S^k(t), v_S^k(t)) \equiv (U(S-t), V(S-t)) \tag{3.23}$$

defines a solution of (3.2). Representation (3.3) for the solution  $u_S^k$  is immediate from the definition of  $\tilde{Z}$  and (3.16). Also, it follows from (3.19) and (3.21) that the solution satisfies (3.4) for any  $p \geq 2$ .

We now prove uniqueness. Let  $(u, v), (u', v') \in \mathcal{H}_S^\infty(\tilde{\mathcal{F}}^S) \times \mathcal{H}_S^2(\tilde{\mathcal{F}}^S)$  be two solutions of (3.9). Then, the differences  $\xi = u - u'$ , and  $\eta = v - v'$  satisfy

$$\xi(t) = \int_t^T \xi(r) \downarrow d\tilde{Z}(r) - \int_t^T \eta(r) d\tilde{W}(r).$$

Using Lemma 5.4 (ii) in the Appendix, we get that

$$\xi(t)^2 + \int_t^S \eta(r)^2 dr = 2 \int_t^S \xi(r)^2 \downarrow d\tilde{Z}(r) - 2 \int_t^S \eta(r)\xi(r)d\tilde{W}(r) + C^k(0) \int_t^S \xi(r)^2 dr.$$

Taking expectations and using Gronwall's inequality it follows that

$$\mathbf{E}\xi(t)^2 + \mathbf{E} \int_t^S \eta(r)^2 dr = 0,$$

The unique solvability of (3.9), and consequently that of (3.2) follows. This completes the proof of the lemma.  $\square$

*Proof of Theorem 2.2.* As in the proof of Lemma 3.1, we will suppress  $n, k, S$  from the notation, unless necessary. Let  $(u, v) \in \mathcal{H}_S^\infty(\mathcal{F}^S) \times \mathcal{H}_S^2(\mathcal{F}^S)$  be the solution of (3.2). We will obtain a solution of (2.5) by taking the logarithmic transform of  $u$ . We begin by showing that

$$\inf_{0 \leq t \leq S} u(t) > 0, \text{ a.s.} \tag{3.24}$$

Recall from (3.13) that  $u(t) = E(S-t)M(S-t)$ ,  $0 \leq t \leq S$ . Clearly  $\inf_{0 \leq t \leq S} E(S-t) > 0$ . Also, from the expression of  $M(t)$  given in (3.11) we see that, for each  $t$ ,  $M(t) > 0$  a.s., since the random variable inside the conditional expectation is strictly positive a.s. Also, since  $M$  is continuous, we have that  $\inf_{0 \leq t \leq S} M(t) > 0$  a.s. Combining these observations we see that (3.24) holds. Define

$$y(t) = -\log u(t), \text{ and } z(t) = \frac{v(t)}{u(t)}. \tag{3.25}$$

We now argue that  $(y, z) \in \mathcal{H}_S^\infty(\mathcal{F}^S) \times \mathcal{H}_S^2(\mathcal{F}^S)$ . For  $y$  note that

$$\begin{aligned} \mathbf{E} \sup_{t \in [0, S]} y(t)^2 &\leq \mathbf{E} \sup_{t \in [0, S]} y(t)^2 \mathbf{1}_{u(t) \leq 1} + \mathbf{E} \sup_{t \in [0, S]} y(t)^2 \mathbf{1}_{u(t) > 1} \\ &\equiv T_1 + T_2, \end{aligned} \tag{3.26}$$

Using the inequality  $0 < \log \theta < \theta$  for all  $\theta > 1$ ,

$$T_2 \leq \mathbf{E} \sup_{t \in [0, S]} u(t)^2 < \infty.$$

Next consider  $T_1$ . From (3.15), (3.23) and an application of Jensen's inequality we have that

$$\begin{aligned} |y(S-t)| \mathbf{1}_{u(S-t) \leq 1} &= -\log(U(t) \mathbf{1}_{U(t) \leq 1} + \mathbf{1}_{U(t) > 1}) \\ &= -\log \mathbf{E} \left[ u_0(\hat{X}(S)) \frac{E(t)}{E(S)} \mathbf{1}_{U(t) \leq 1} + \mathbf{1}_{U(t) > 1} \mid \mathcal{G}_t \right] \\ &\leq -\mathbf{E} \left[ \log \left( u_0(\hat{X}(S)) \frac{E(t)}{E(S)} \mathbf{1}_{U(t) \leq 1} + \mathbf{1}_{U(t) > 1} \right) \mid \mathcal{G}_t \right] \\ &= -\mathbf{E} \left[ \log \left( u_0(\hat{X}(S)) \frac{E(t)}{E(S)} \right) \mid \mathcal{G}_t \right] \mathbf{1}_{U(t) \leq 1}. \end{aligned}$$

Recalling that  $u_0 = \exp\{-h_0\}$ , we have

$$|y(S-t)| \mathbf{1}_{u(S-t) \leq 1} \leq \mathbf{E} \left[ |h_0(\hat{X}(S))| \mid \mathcal{G}_t \right] + \mathbf{E} \left[ |\tilde{Z}(t) - \tilde{Z}(S)| \mid \mathcal{G}_t \right] + \frac{1}{2} C^k(0)(S-t).$$

Recalling that  $\{\mathcal{G}_t\}$  is a filtration and that from (3.10) and (2.2)

$$\mathbf{E} \left( \sup_{0 \leq t \leq S} |\tilde{Z}(t)|^2 + |h_0(\hat{X}(S))|^2 \right) < \infty,$$

we have by an application of Doob's inequality that for some  $C_1 \in (0, \infty)$

$$T_1 = \mathbf{E} \sup_{t \in [0, S]} y(t)^2 \mathbf{1}_{u(t) \leq 1} = \mathbf{E} \sup_{t \in [0, S]} y(S-t)^2 \mathbf{1}_{u(S-t) \leq 1} < \infty.$$

Using the above estimates on  $T_1$  and  $T_2$  in (3.26) we see that  $y \in \mathcal{H}_S^\infty(\mathcal{F}^S)$ .

We now consider  $z$ . Let  $p \geq 2$  and  $q$  be such that  $p^{-1} + q^{-1} = 1$ . Then using Holder's inequality

$$\begin{aligned} \mathbf{E} \int_0^S z(r)^2 dr &= \mathbf{E} \int_0^S \left( \frac{v(r)}{u(r)} \right)^2 dr \\ &\leq \mathbf{E} \sup_{t \in [0, S]} u(t)^{-2} \int_0^S v(r)^2 dr \\ &\leq \left( \mathbf{E} \sup_{t \in [0, S]} u(t)^{-2p} \right)^{p^{-1}} \left( \mathbf{E} \left( \int_0^S v(r)^2 dr \right)^q \right)^{q^{-1}}. \end{aligned} \tag{3.27}$$

From (3.15) and Jensen's inequality

$$\begin{aligned} U(t)^{-2p} &= (\mathbf{E}[U(t) \mid \mathcal{G}_t])^{-2p} \\ &= \left( \mathbf{E} \left[ u_0(\hat{X}(S)) \frac{E(t)}{E(S)} \mid \mathcal{G}_t \right] \right)^{-2p} \\ &\leq \mathbf{E} \left[ \left( u_0(\hat{X}(S)) \right)^{-2p} \frac{E(S)^{2p}}{E(t)^{2p}} \mid \mathcal{G}_t \right]. \end{aligned} \tag{3.28}$$

Recalling (3.23), we have that

$$\begin{aligned} \mathbf{E} \sup_{0 \leq t \leq S} u(t)^{-2p} &= \mathbf{E} \sup_{0 \leq t \leq S} U(t)^{-2p} \\ &\leq \mathbf{E} \sup_{0 \leq t \leq S} \mathbf{E} \left[ \left( u_0(\hat{X}(S)) \right)^{-2p} E(S)^{2p} \sup_{0 \leq r \leq S} E(r)^{-2p} \mid \mathcal{G}_t \right]. \end{aligned} \tag{3.29}$$

Also, from (3.10) and (2.2)

$$\mathbf{E} \left[ \left( u_0(\hat{X}(S)) \right)^{-4p} E(S)^{4p} \sup_{0 \leq r \leq S} E(r)^{-4p} \right] < \infty.$$

Since  $\{\mathcal{G}_t\}$  is a filtration, we have that the conditional expectation in (3.29) is a martingale and so by Doob's maximal inequality it follows that

$$\mathbf{E} \sup_{0 \leq t \leq S} u(t)^{-2p} < \infty.$$

Combining this estimate with (3.21), (3.27) and recalling (3.23), we have that  $z \in \mathcal{H}_S^2(\mathcal{F}^S)$ .

To finish the proof of existence of solutions, we now verify that  $(y, z)$  defined in (3.25) satisfy (2.5). We will apply Lemma 5.4 (i) with  $\alpha = u, \beta = 0, \gamma = -u, \delta = v$ , and  $\phi(x) = -\log(x)$ . Note that although  $\phi$  is only  $C^2$  on  $(0, \infty)$ , (3.24) guarantees the applicability of Itô's formula. Representation (3.2), and Lemma 5.4 imply that

$$\begin{aligned} y(t) &= y(0) + \int_0^t \frac{u(r)}{u(r)} dZ(r) - \int_0^t \frac{v(r)}{u(r)} \downarrow dW(r) + \int_0^t \frac{u(r)^2}{2u(r)^2} C^k(0) dr - \int_0^t \frac{v(r)^2}{2u(r)^2} dr \\ &= y(0) + Z(t) - \int_0^t \frac{v(r)}{u(r)} \downarrow dW(r) - \frac{1}{2} \int_0^t \left( \left( \frac{v(r)}{u(r)} \right)^2 - C^k(0) \right) dr. \end{aligned}$$

From the fact that  $z(t) = v(t)/u(t)$ , and  $y(0) = h(X^S(0, x))$  we see that this equation is the same as (2.5). This completes the proof of existence.

We now prove uniqueness. Suppose  $(y, z), (\tilde{y}, \tilde{z}) \in \mathcal{H}_S^\infty(\mathcal{F}^S) \times \mathcal{H}_S^2(\mathcal{F}^S)$  are two solutions of (2.5). Let  $(\bar{y}, \bar{z}) = (y - \tilde{y}, z - \tilde{z})$ . For  $M \in (0, \infty)$ , define  $\psi_M : \mathbb{R} \times \mathbb{R} \rightarrow [-M, M]$  as  $\psi_M(a, b) = \frac{1}{2}(2a - b)\mathbf{1}_{|2a - b| \leq M}$ . Let  $\varphi_M(r) = \psi_M(z(r), \bar{z}(r))$ ,  $r \in [0, S]$  and let  $y_M$  be a continuous process defined as

$$y_M(t) = - \int_0^t \bar{z}(r)\varphi_M(r)dr - \int_0^t \bar{z}(r) \downarrow dW(r). \tag{3.30}$$

We will now show that

$$y_M(t) = 0, \text{ a.s. for all } t \in [0, S] \text{ and } M \in (0, \infty). \tag{3.31}$$

Note that if (3.31) holds, we have on sending  $M \rightarrow \infty$ , and observing that  $y_M(t) \rightarrow \bar{y}(t)$  in probability, for every  $t \in [0, S]$ , that  $y$  and  $\tilde{y}$  are indistinguishable. Moreover, an application of Itô's formula (see Lemma 5.4(i)) shows that

$$y_M^2(t) = -2 \int_0^t y_M(r)\bar{z}(r)\varphi_M(r)dr - 2 \int_0^t y_M(r)\bar{z}(r) \downarrow dW(r) - \int_0^t \bar{z}^2(r)dr$$

and so if (3.31) holds, we have that  $z(t) = \tilde{z}(t)$ , a.e.  $t \in [0, S]$ , a.s. Combining the above observations we see that in order to prove uniqueness, it suffices to verify (3.31).

From Tanaka's formula (cf. Theorem IV.68 in [18]) it follows that

$$\begin{aligned} (y_M(t))_+ &= - \int_0^t \mathbf{1}_{\{y_M(r) > 0\}} \bar{z}(r)\varphi_M(r)dr \\ &\quad - \int_0^t \mathbf{1}_{\{y_M(r) > 0\}} \bar{z}(r) \downarrow dW(r) - \frac{1}{2}L^0(t), \end{aligned}$$

where  $y_+ = \max\{y, 0\}$ , and  $L^0$  is the local time at 0 process for  $y_M$  (see Chapter IV of [18]). In particular,  $L^0$  is non-decreasing, non-negative process such that

$$\int_{[0, \infty)} \mathbf{1}_{\{y_M(t) > 0\}} dL^0(t) = 0. \tag{3.32}$$

We remark that the cited theorem establishes the above formula for equations with forward stochastic integrals, however the version with backward integrals used here follows by straightforward modifications of the proof.

Define for  $n \in \mathbb{N}$ ,  $\xi_n : [0, \infty) \rightarrow [0, \infty)$  as

$$\xi_n(u) = (u \wedge n)^2 + 2n(u - n)_+, \quad u \in [0, \infty).$$

Then  $\xi_n$  is a  $C^1$ -convex function with

$$\xi_n'(u) = 2(u \wedge n), \quad u \in [0, \infty). \tag{3.33}$$

By Meyer-Itô formula (cf. Theorem IV.70 in [18])

$$\begin{aligned} \xi_n((y_M(t))_+) &= - \int_0^t \xi_n'((y_M(s))_+) \mathbf{1}_{\{y_M(s) > 0\}} \bar{z}(r)\varphi_M(r)dr \\ &\quad - \int_0^t \xi_n'((y_M(s))_+) \mathbf{1}_{\{y_M(s) > 0\}} \bar{z}(r) \downarrow dW(r) \\ &\quad - \frac{1}{2} \int_0^t \xi_n'((y_M(s))_+) dL^0(s) - \frac{1}{2} \int_{-\infty}^\infty L^a(t) \mu(da). \end{aligned} \tag{3.34}$$

where  $L^a$  is the local time process of  $(y_M)_+$  at level  $a$  and  $\mu$  represents the second derivative of  $\xi_n$  in the generalized function sense. Since  $\xi_n$  is convex,  $\mu$  is a (non-negative) measure and in fact equals

$$\mu(da) = 2\mathbf{1}_{[0,n]}(a)da.$$

Thus (cf. Corollary IV.1 of [18])

$$\int_{-\infty}^{\infty} L^a(t)\mu(da) = 2 \int_0^n L^a(t)da = 2 \int_0^t \mathbf{1}_{(0,n]}((y_M(s))_+) \bar{z}^2(s) ds.$$

Combining this with the fact that the third term on the right side of (3.34) is zero, we have from (3.33) that

$$\begin{aligned} \xi_n(y_M(t)_+) &+ \int_0^t \mathbf{1}_{(0,n]}((y_M(r))_+) \bar{z}^2(r) dr \\ &= -2 \int_0^t \mathbf{1}_{(0,n]}((y_M(r))_+) (y_M(r))_+ \bar{z}(r) \varphi_M(r) dr \\ &\quad - 2n \int_0^t \mathbf{1}_{(n,\infty)}((y_M(r))_+) \bar{z}(r) \varphi_M(r) dr \\ &\quad - \int_0^t \mathbf{1}_{\{y_M(r)>0\}} \xi'_n((y_M(r))_+) \bar{z}(r) \downarrow dW(r). \end{aligned} \tag{3.35}$$

Using Young's inequality we have that, for any  $\alpha > 0$ ,

$$\begin{aligned} \int_0^t \mathbf{1}_{(0,n]}((y_M(r))_+) (y_M(r))_+ |\bar{z}(r)| |\varphi_M(r)| dr &\leq \frac{\alpha}{2} \int_0^t \mathbf{1}_{(0,n]}((y_M(r))_+) |\bar{z}(r)|^2 dr \\ &\quad + \frac{1}{2\alpha} \int_0^t (y_M(r))_+^2 |\varphi_M(r)|^2 dr. \end{aligned}$$

Using the above estimate with  $\alpha < 1$  in (3.35), we have

$$\begin{aligned} \xi_n(y_M(t)_+) &\leq \frac{M^2}{\alpha} \int_0^t (y_M(r))_+^2 dr \\ &\quad + 2nM \int_0^t \mathbf{1}_{(n,\infty)}((y_M(r))_+) |\bar{z}(r)| dr \\ &\quad - \int_0^t \mathbf{1}_{\{y_M(r)>0\}} \xi'_n((y_M(r))_+) \bar{z}(r) \downarrow dW(r). \end{aligned} \tag{3.36}$$

Next, from (3.30), using that  $|\varphi_M(r)| \leq M$  and Doob's inequality, we have

$$\mathbf{E} \sup_{t \in [0,T]} y_M^2(t) \leq 2M^2 \mathbf{S} \mathbf{E} \int_0^S \bar{z}^2(r) dr + 8 \mathbf{E} \int_0^S \bar{z}^2(r) dr \equiv C_1 < \infty. \tag{3.37}$$

Let

$$\tau_{M,n} = \inf \{t \in [0, S] : y_M(t) \geq n\}, \quad n \in \mathbf{N},$$

where infimum over an empty set, by convention, is taken to be  $S$ . Then

$$\begin{aligned} n \mathbf{E} \int_0^t \mathbf{1}_{(n,\infty)}((y_M(r))_+) |\bar{z}(r)| dr &\leq n \mathbf{E} \mathbf{1}_{\{\tau_{M,n} < S\}} \int_{\tau_{M,n} \wedge t}^t |\bar{z}(r)| dr \\ &\leq n \left( \mathbf{P} \left( \sup_{t \in [0,S]} y_M(t) \geq n \right) \right)^{1/2} \left( \mathbf{E} \left( \int_{\tau_{M,n} \wedge t}^t \bar{z}(r) dr \right)^2 \right)^{1/2} \\ &\leq C_1^{1/2} \left( \mathbf{E} \left[ (t - \tau_{M,n} \wedge t) \int_0^S \bar{z}^2(r) dr \right] \right)^{1/2}, \end{aligned} \tag{3.38}$$

where the third inequality is a consequence of (3.37). Since  $(t - \tau_{M,n} \wedge t)$  converges to 0 as  $n \rightarrow \infty$  and  $\mathbf{E} \int_0^S \bar{z}^2(r) dr < \infty$ , we have that the expression on the last line of the above display converges to 0 as  $n \rightarrow \infty$ . Thus we have shown that

$$\lim_{n \rightarrow \infty} n \mathbf{E} \int_0^t \mathbf{1}_{(n, \infty)}((y_M(r))_+) |\bar{z}(r)| dr = 0. \tag{3.39}$$

Taking expectations in (3.36) and noting that since  $\xi'_n$  is bounded, the expectation of the third term on the right side of (3.36) is zero, we have

$$\limsup_{n \rightarrow \infty} \mathbf{E} \xi_n(y_M(t)_+) \leq \frac{M^2}{\alpha} \int_0^t \mathbf{E}(y_M(s)_+)^2 ds.$$

Finally, noting that  $\xi_n(u) \rightarrow u^2$  as  $n \rightarrow \infty$ , for all  $u \in [0, \infty)$ , we have by Fatou's lemma that

$$\mathbf{E}(y_M(t)_+)^2 \leq \frac{M^2}{\alpha} \int_0^t \mathbf{E}(y_M(s)_+)^2 ds.$$

Gronwall's lemma now yields that  $(y_M(t))_+ = 0$  for all  $t \in [0, S]$ . A similar argument shows that  $(y_M(t))_-$  and consequently (3.31) follows. As argued earlier, this proves the desired uniqueness.  $\square$

### 4 Proof of Theorem 2.3

Fix  $0 \leq t \leq S \leq T$ . The representation in Lemma 3.1 and (3.25) give

$$\begin{aligned} y_S^k(t, x) &= -\log \mathbf{E} \left[ e^{-h_0(X^S(0, x))} \exp\{-Z^k(t, x) - \frac{1}{2}C^k(0)t\} \middle| \mathcal{F}_t^S \right] \\ &= -\log \mathbf{E} \left[ \mathbf{E} \left[ e^{-h_0(X^S(0, x))} \exp\{-Z^k(t, x) - \frac{1}{2}C^k(0)t\} \middle| \mathcal{F}_t^S \vee \sigma\{W(t)\} \right] \middle| \mathcal{F}_t^S \right]. \end{aligned} \tag{4.1}$$

Define a  $C([0, t] : \mathbb{R})$  valued random variable  $X^{S,t}$  as

$$X^{S,t}(r) = X^S(r, x), \quad r \in [0, t]$$

and a  $C([0, S] : \mathbb{R}^\infty)$  valued random variable  $\beta$  as

$$\beta(r) = (\beta_m(r))_{m \geq 1}, \quad r \in [0, S].$$

Then there is a measurable map

$$\Psi : C([0, t] : \mathbb{R}) \times C([0, S] : \mathbb{R}^\infty) \rightarrow \mathbb{R}_+$$

such that

$$\Psi(X^{S,t}, \beta) = \exp\{-Z^k(t, x) - \frac{1}{2}C^k(0)t\}.$$

In fact one has the following characterization of  $\Psi$ . For  $\omega \in C([0, t] : \mathbb{R})$  define

$$M_\omega^k(t) = \sum_{m \in \mathbb{N}} \int_0^t \langle \zeta_{\omega(r)}^k, \gamma_m \rangle d\beta_m(r).$$

Then  $\Psi$  satisfies

$$\Psi(\omega, \beta) = \exp\{-M_\omega^k(t) - \frac{1}{2}C^k(0)t\}, \text{ for all } \omega \in C([0, t] : \mathbb{R}), \text{ a.s.}$$

Let  $\mathbf{P}_t^{\mu,\nu}$  denote the Brownian bridge measure on  $C([0, t] : \mathbb{R})$  with starting point  $\mu$  and ending point  $\nu$ . Define  $\Psi_0 : [0, S] \times \mathbb{R} \times \mathbb{R} \times C([0, S] : \mathbb{R}^\infty) \rightarrow \mathbb{R}_+$  as

$$\Psi_0(t, \mu, \nu, \vartheta) = \int_{C([0, t] : \mathbb{R})} \Psi(\omega, \vartheta) d\mathbf{P}_t^{\mu,\nu}(\omega).$$

In particular

$$\begin{aligned} \Psi_0(t, \mu, \nu, \beta) &= \int_{C([0, t] : \mathbb{R})} \Psi(\omega, \beta) d\mathbf{P}_t^{\mu,\nu}(\omega) \\ &= \int_{C([0, t] : \mathbb{R})} \exp\{-M_\omega^k(t) - \frac{1}{2}C^k(0)t\} d\mathbf{P}_t^{\mu,\nu}(\omega) \\ &\equiv \mathbf{E}_t^{\mu,\nu} \left[ \exp\{-M_\bullet^k(t) - \frac{1}{2}C^k(0)t\} \right]. \end{aligned}$$

Next, using the independence of  $W(t)$  and  $\mathcal{F}_t^S$  we have

$$\begin{aligned} \mathbf{E} \left[ \exp\{-Z^k(t, x) - \frac{1}{2}C^k(0)t\} \mid \mathcal{F}_t^S \vee \sigma\{W(t)\} \right] &= \mathbf{E} [\Psi(X^{S,t}, \beta) \mid \mathcal{F}_t^S \vee \sigma\{W(t)\}] \\ &= \Psi_0(t, \gamma + W_t, \gamma, \beta) \\ &= \mathbf{E}_t^{\gamma+W_t, \gamma} \left[ \exp\{-M_\bullet^k(t) - \frac{1}{2}C^k(0)t\} \right], \end{aligned}$$

where  $\gamma = x + W(S) - W(t)$ . Therefore

$$\begin{aligned} &\mathbf{E} \left[ e^{-h_0(X^S(0,x))} \exp\{-Z^k(t, x) - \frac{1}{2}C^k(0)t\} \mid \mathcal{F}_t^S \right] \\ &= \mathbf{E} \left[ e^{-h_0(\gamma+W(t))} \mathbf{E}_t^{\gamma+W(t), \gamma} \left[ \exp\{-M_\bullet^k(t) - \frac{1}{2}C^k(0)t\} \right] \mid \mathcal{F}_t^S \right] \\ &= \int_{\mathbb{R}} e^{-h(y)} G_t(\gamma - y) \mathbf{E}_t^{y, \gamma} \left[ \exp\{-M_\bullet^k(t) - \frac{1}{2}C^k(0)t\} \right] dy, \end{aligned}$$

where  $G_t$  is the standard Heat Kernel. The last expression is seen from expression (2.17) of [2] to be same as  $\psi_t^k(\gamma)$ , where  $\psi_t^k$  is the solution of the regularized stochastic heat equation.

$$\psi_t^k(x) = G_t \star \psi_0(x) + \int_0^t \langle G_{t-s} \star \psi_s^k, dB_s^k \rangle. \tag{4.2}$$

(See Section 2.2 of [2].) Therefore

$$\mathcal{G}_S^k(t, x) = -\log \psi_t^k(x + W(S) - W(t)), \quad 0 \leq t \leq S \leq T. \tag{4.3}$$

The result now follows from Theorem 2.2 of [2].  $\square$

### 5 Appendix.

In this section we collect some basic results on forward-backward stochastic integrals that are used at various places in this work. Most of the statements follow by minor modifications of classical results (eg. [16]) and thus only partial sketches are provided. Throughout this section we will fix  $S \in (0, \infty)$ ,  $x \in \mathbb{R}$  and  $k \in \mathbb{N}$ . As previously, we will suppress  $k$  and  $z$  from the notation when writing  $Z^k(t, x)$ ,  $\tilde{Z}^k(t, x)$  etc.

Define  $\sigma$ -fields

$$\mathcal{G}_r^S = \mathcal{F}_{r,S}^W \vee \mathcal{F}_S^B, \quad \mathcal{H}_r^S = \mathcal{F}_{0,S}^W \vee \mathcal{F}_r^B, \quad \tilde{\mathcal{G}}_r^S = \mathcal{F}_S^{\tilde{W}} \vee \mathcal{F}_{r,S}^{\tilde{B}}, \quad \tilde{\mathcal{H}}_r^S = \mathcal{F}_r^{\tilde{W}} \vee \mathcal{F}_{0,S}^{\tilde{B}}.$$



Abusing terminology, we say a stochastic process  $\{A(r)\}_{0 \leq r \leq S}$  is adapted to a collection of  $\sigma$ -fields  $\{\mathcal{U}_r\}_{0 \leq r \leq S}$  if  $A(r)$  is  $\mathcal{U}_r$  measurable for every  $r \in [0, S]$ . For such a family of  $\sigma$ -fields we denote by  $\mathcal{A}^2(\mathcal{U})$  the collection of all adapted processes  $\{A(r)\}$  such that  $\int_0^S |A(r)|^2 dr < \infty$  a.s. Then the following stochastic integrals are well defined:

$$\int_0^t A(r) \downarrow dW(r), A \in \mathcal{A}^2(\mathcal{G}^S); \int_0^t A(r) dZ(r), A \in \mathcal{A}^2(\mathcal{H}^S),$$

$$\int_0^t A(r) \downarrow d\tilde{Z}(r), A \in \mathcal{A}^2(\tilde{\mathcal{G}}^S); \int_0^t A(r) d\tilde{W}(r), A \in \mathcal{A}^2(\tilde{\mathcal{H}}^S), t \in [0, S].$$

Indeed, consider for example the first stochastic integral. If  $A$  is of the form  $A(r) = \zeta \mathbf{1}_{[a,b)}(r)$ , where  $\zeta$  is a bounded  $\mathcal{G}_b^S$  measurable random variable and  $0 \leq a < b \leq S$ , then

$$\int_0^t A(r) \downarrow dW(r) \equiv \zeta (W(b \wedge t) - W(a \wedge t)).$$

The integral is extended to linear combinations of such elementary processes by linearity, and then by denseness and  $L^2$ -isometry to all  $A \in \mathcal{A}^2(\mathcal{G}^S)$  satisfying  $\mathbf{E} \int_0^S |A(r)|^2 dr < \infty$ ; and finally by localization to all  $A \in \mathcal{A}^2(\mathcal{G}^S)$ .

The following elementary lemma gives a basic relation between forward and backward integrals.

**Lemma 5.1.** *Let  $K \in \mathcal{A}^2(\mathcal{H}^S)$  and  $H \in \mathcal{A}^2(\tilde{\mathcal{H}}^S)$ . Let*

$$\tilde{K}(t) = K(S - t), \quad \tilde{H}(t) = H(S - t), \quad t \in [0, S].$$

*Then  $\tilde{K} \in \mathcal{A}^2(\tilde{\mathcal{G}}^S)$  and  $\tilde{H} \in \mathcal{A}^2(\mathcal{G}^S)$ . Furthermore, for  $t \in [0, S]$ ,*

$$\int_0^t H(r) d\tilde{W}(r) = - \int_{S-t}^S \tilde{H}(r) \downarrow dW(r),$$

$$\int_0^t K(r) dZ(r) = - \int_{S-t}^S \tilde{K}(r) \downarrow d\tilde{Z}(r)$$

*Proof.* The first statement in the lemma is an immediate consequence of (3.6) and (3.7). Of the two equalities in the above display, we only prove the first one. The proof of the second identity follows by a similar argument. Consider first the case where  $H(t) = \zeta \mathbf{1}_{(a,b]}(t)$ , where  $\zeta$  is a bounded  $\tilde{\mathcal{H}}_a^S$  measurable random variable, and  $0 \leq a < b \leq S$ . In that case, note that

$$\tilde{H}(r) = H(S - r) = \zeta \mathbf{1}_{(a,b]}(S - r) = \zeta \mathbf{1}_{[S-b, S-a)}(r),$$

and,

$$\int_{S-t}^S \tilde{H}(r) \downarrow dW(r)$$

$$= \zeta [(W((S - a) \vee (S - t))) - (W((S - b) \vee (S - t)))]$$

$$= \zeta [(W((S - a) \vee (S - t)) - W(S)) - (W((S - b) \vee (S - t)) - W(S))]$$

$$= \zeta (\tilde{W}(t \wedge a) - \tilde{W}(t \wedge b))$$

$$= - \int_0^t H(r) d\tilde{W}(r).$$

The general case follows by linearity, denseness (along with  $L^2$  isometry) and a localization argument. Details are omitted.  $\square$

As an immediate consequence of the lemma we have the following corollary.

**Corollary 5.2.** *A pair of processes  $(\hat{u}, \hat{v}) \in \mathcal{H}_S^\infty(\tilde{\mathcal{F}}^S) \times \mathcal{H}_S^2(\tilde{\mathcal{F}}^S)$  solves (3.9) if and only if  $(u, v)$ , defined as  $(u(t), v(t)) = (u(S - t), v(S - t))$ ,  $t \in [0, S]$ , solves (3.2).*

*Proof.* The proof is immediate from Lemma 5.1. □

The following elementary lemma will be used in the proof of (3.22).

**Lemma 5.3.** *Let  $\varphi$  be a  $C^1$  function on  $\mathbb{R}$  and  $\psi : [0, S] \rightarrow \mathbb{R}$  be a continuous function. Suppose that for all  $t \in [0, T]$ ,  $\varphi(\tilde{Z}(t)) = \varphi_1(\tilde{Z}(t) - \tilde{Z}(T))\varphi_2(\tilde{Z}(T))$ , a.s., for some continuous functions  $\varphi_1, \varphi_2$ . Then for all  $t \in [0, S]$ ,*

$$\int_0^t \varphi(\tilde{Z}(r))\psi(r)d\tilde{Z}(r) = \varphi_2(\tilde{Z}(T)) \int_0^t \varphi_1(\tilde{Z}(r) - \tilde{Z}(T))\psi(r) \downarrow d\tilde{Z}(r) - C^k(0) \int_0^t \varphi'(\tilde{Z}(r))\psi(r)dr. \tag{5.1}$$

*Proof.* Fix  $t \in [0, S]$  and let  $\Pi_n = \{0 = t_0^{(n)} < t_1^{(n)} < t_2^{(n)} \dots < t_k^{(n)} = t\}$  be a partition of  $[0, t]$  such that  $|\Pi_n| \rightarrow 0$  as  $n \rightarrow \infty$ . Then (suppressing  $n$ ) letting  $\Delta_i \tilde{Z} = \tilde{Z}(t_{i+1}) - \tilde{Z}(t_i)$ , we see that  $\varphi_2(\tilde{Z}(T)) \int_0^t \varphi_1(\tilde{Z}(r) - \tilde{Z}(T))\psi(r) \downarrow d\tilde{Z}(r)$  is the limit in probability of

$$\begin{aligned} & \varphi_2(\tilde{Z}(T)) \sum_{i=0}^{k-1} \varphi_1(\tilde{Z}(t_{i+1}) - \tilde{Z}(T))\psi(t_{i+1})\Delta_i \tilde{Z} \\ &= \sum_{i=0}^{k-1} \varphi(\tilde{Z}(t_{i+1}))\psi(t_{i+1})\Delta_i \tilde{Z} \\ &= \sum_{i=0}^{k-1} \varphi(\tilde{Z}(t_i))\psi(t_{i+1})\Delta_i \tilde{Z} \\ &+ \sum_{i=0}^{k-1} (\varphi(\tilde{Z}(t_{i+1})) - \varphi(\tilde{Z}(t_i))) \psi(t_{i+1})\Delta_i \tilde{Z}. \end{aligned} \tag{5.2}$$

From standard arguments it follows that, in probability,

$$\lim_{n \rightarrow \infty} \sum_{i=0}^{k-1} (\varphi(\tilde{Z}(t_{i+1})) - \varphi(\tilde{Z}(t_i))) \psi(t_{i+1})\Delta_i \tilde{Z} = C^k(0) \int_0^t \varphi'(\tilde{Z}(r))\psi(r)dr.$$

Likewise, it is easily seen that

$$\lim_{n \rightarrow \infty} \sum_{i=0}^{k-1} \varphi(\tilde{Z}(t_i))\psi(t_{i+1})\Delta_i \tilde{Z} = \int_0^t \varphi(\tilde{Z}(r))\psi(r)d\tilde{Z}(r),$$

in probability. These two identities combined with (5.2) give the result. □

We now present a variation of Itô's formula that is used in our work.

**Lemma 5.4.** *Let  $\phi \in C^2(\mathbb{R})$ .*

(i) *Let processes  $\alpha \in \mathcal{H}_S^\infty(\mathcal{F}^S), \beta, \gamma, \delta \in \mathcal{H}_S^2(\mathcal{F}^S)$  be such that*

$$\alpha(t) = \alpha(0) + \int_0^t \beta(r)dr + \int_0^t \gamma(r)dZ^k(r) + \int_0^t \delta(r) \downarrow dW(r), \quad 0 \leq t \leq T.$$

Then, for all  $t \in [0, S]$

$$\begin{aligned} \phi(\alpha(t)) &= \phi(\alpha(0)) + \int_0^t \phi'(\alpha(r))\beta(r)dr + \int_0^t \phi'(\alpha(r))\gamma(r)dZ^k(r) \\ &\quad + \int_0^t \phi'(\alpha(r))\delta(r) \downarrow dW(r) + \frac{C^k(0)}{2} \int_0^t \phi''(\alpha(r))\gamma(r)^2 dr \\ &\quad - \frac{1}{2} \int_0^t \phi''(\alpha(r))\delta(r)^2 dr. \end{aligned}$$

(ii) Let processes  $\alpha \in \mathcal{H}_S^\infty(\tilde{\mathcal{F}}^S), \beta, \gamma, \delta \in \mathcal{H}_S^2(\tilde{\mathcal{F}}^S)$  be such that

$$\alpha(t) = \alpha(0) + \int_0^t \beta(r)dr + \int_0^t \gamma(r) \downarrow d\tilde{Z}^k(r) + \int_0^t \delta(r)d\tilde{W}(r), \quad 0 \leq t \leq T.$$

Then, for all  $t \in [0, S]$

$$\begin{aligned} \phi(\alpha(t)) &= \phi(\alpha(0)) + \int_0^t \phi'(\alpha(r))\beta(r)dr + \int_0^t \phi'(\alpha(r))\gamma(r) \downarrow d\tilde{Z}^k(r) \\ &\quad + \int_0^t \phi'(\alpha(r))\delta(r)d\tilde{W}(r) - \frac{C^k(0)}{2} \int_0^t \phi''(\alpha(r))\gamma(r)^2 dr \\ &\quad + \frac{1}{2} \int_0^t \phi''(\alpha(r))\delta(r)^2 dr. \end{aligned}$$

*Proof.* We will only consider (i). The statement in (ii) follows similarly. The proof follows using standard arguments (cf. Theorem 3.3.3 in [9]). We merely comment on one key point. Suppose that  $\phi''$  is bounded. (The general case can be reduced to such a setting by localization.) Fix  $t \in [0, S]$  and let  $\Pi_n = \{0 = t_0^{(n)} < t_1^{(n)} < t_2^{(n)} \dots < t_k^{(n)} = t\}$  be a partition of  $[0, t]$  such that  $|\Pi_n| \rightarrow 0$  as  $n \rightarrow \infty$ . Then the only change to standard proofs is in the treatment of the term

$$\sum_{i=1}^k \phi''(\alpha_{i-1})\Delta_i W \Delta_i Z, \tag{5.3}$$

where for a process  $\zeta$ , we write  $\Delta_i \zeta = \zeta(t_i) - \zeta(t_{i-1})$ . One needs to argue that the expression in (5.3) approaches 0 as  $n \rightarrow \infty$ , which follows on noting that

$$\begin{aligned} \mathbf{E} \left[ \sum_{i=1}^k \phi''(\alpha_{i-1})\Delta_i W \Delta_i Z \right]^2 &= \mathbf{E} \left[ \sum_{i=1}^k (\phi''(\alpha_{i-1}))^2 (\Delta_i W)^2 (\Delta_i Z)^2 \right] \\ &\leq \sup_x |\phi''(x)|^2 \mathbf{E} \left[ \sum_{i=1}^k (\Delta_i W)^2 \mathbf{E} \left[ (\Delta_i Z)^2 \mid \mathcal{F}_{t_{i-1}}^B \vee \mathcal{F}_S^W \right] \right] \\ &= \sup_x |\phi''(x)|^2 C^k(0) \sum_{i=1}^k (t_i - t_{i-1}) \mathbf{E}(\Delta_i W)^2 \\ &= \sup_x |\phi''(x)|^2 C^k(0) \sum_{i=1}^k (t_i - t_{i-1})^2, \end{aligned}$$

where the first equality follows on noting that by a conditioning argument the cross-product terms do not contribute while the next to last equality follows from (2.4).  $\square$

Finally, we give the proof of (3.22).

**Proof of (3.22).** Note that  $\{\tilde{Z}(t)\}_{t \in [0, S]}$  is a martingale with respect to the filtration  $\tilde{\mathcal{G}}_t = \mathcal{F}_{0,t}^{\tilde{B}} \vee \mathcal{F}_S^{\tilde{W}}$ , with quadratic variation given as  $\langle Z \rangle_t = C^k(0)t$ . Thus, by an application of Itô's formula, we have that

$$\begin{aligned} E(S) - E(t) &= - \left[ \int_t^S E(r) d\tilde{Z}(r) - C^k(0) \int_t^S E(r) dr \right] \\ &= -E(T) \int_t^S \frac{E(r)}{E(T)} \downarrow d\tilde{Z}(r), \end{aligned} \tag{5.4}$$

where the second equality is a consequence of Lemma 5.3 on taking  $\varphi(x) = \varphi_1(x) = \varphi_2(x) = e^{-x}$  and  $\psi(t) = \exp\{\frac{1}{2}C^k(0)t\}$ . Also, recall from (3.12) that

$$M(t) = M(S) - \int_t^S J(r) d\tilde{W}(r), \quad 0 \leq t \leq S. \tag{5.5}$$

Let  $\Pi_n = \{t = t_0^{(n)} < t_1^{(n)} < t_2^{(n)} \dots < t_k^{(n)} = S\}$  be a partition of  $[t, S]$  such that  $|\Pi_n| \rightarrow 0$  as  $n \rightarrow \infty$ . Then, suppressing  $n$  in the notation

$$\begin{aligned} U(t) - U(S) &= - \sum_{i=1}^k (U(t_i) - U(t_{i-1})) \\ &= - \sum_{i=1}^k (M(t_i)E(t_i) - M(t_{i-1})E(t_{i-1})) \\ &= - \sum_{i=1}^k M(t_i)(E(t_i) - E(t_{i-1})) - \sum_{i=1}^k E(t_{i-1})(M(t_i) - M(t_{i-1})) \end{aligned}$$

The equality in (3.22) now follows from (5.4) and (5.5) on taking limit as  $n \rightarrow \infty$  in the last line.

## References

- [1] Gideon Amir, Ivan Corwin, and Jeremy Quastel. Probability distribution of the free energy of the continuum directed random polymer in 1 + 1 dimensions. *Comm. Pure Appl. Math.*, 64(4):466–537, 2011. MR-2796514
- [2] Lorenzo Bertini and Nicoletta Cancrini. The stochastic heat equation: Feynman-Kac formula and intermittence. *J. Statist. Phys.*, 78(5-6):1377–1401, 1995. MR-1316109
- [3] Lorenzo Bertini and Giambattista Giacomin. Stochastic Burgers and KPZ equations from particle systems. *Comm. Math. Phys.*, 183(3):571–607, 1997. MR-1462228
- [4] I. Corwin. The Kardar-Parisi-Zhang equation and universality class. *Random Matrices: Theory and Applications*, 1, 2012. MR-2930377
- [5] P. Del Moral. Pierre Feynman-Kac formulae. Genealogical and interacting particle systems with applications. *Probability and its Applications (New York)*. Springer-Verlag, New York, 2004. MR-2044973
- [6] Arash Fahim, Nizar Touzi, and Xavier Warin. A probabilistic numerical method for fully nonlinear parabolic PDEs. *Ann. Appl. Probab.*, 21(4):1322–1364, 2011. MR-2857450
- [7] Patrícia Gonçalves and Milton Jara. Scaling limits of a tagged particle in the exclusion process with variable diffusion coefficient. *J. Stat. Phys.*, 132(6):1135–1143, 2008. MR-2430777
- [8] Martin Hairer. Solving the KPZ equation. *Preprint*. MR-3071506

- [9] Ioannis Karatzas and Steven E. Shreve. *Brownian motion and stochastic calculus*, volume 113 of *Graduate Texts in Mathematics*. Springer-Verlag, New York, second edition, 1991. MR-1121940
- [10] Mehran Kardar, Giorgio Parisi, and Yi-Cheng Zhang. Dynamic scaling of growing interfaces. *Phys. Rev. Lett.*, 56:889–892, Mar 1986.
- [11] Magdalena Kobylanski. Backward stochastic differential equations and partial differential equations with quadratic growth. *Ann. Probab.*, 28(2):558–602, 2000. MR-1782267
- [12] Yuping Liu and Jin Ma. Optimal reinsurance/investment problems for general insurance models. *Ann. Appl. Probab.*, 19(4):1495–1528, 2009. MR-2538078
- [13] Jin Ma and Jiongmin Yong. *Forward-backward stochastic differential equations and their applications*, volume 1702 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 1999. MR-1704232
- [14] E. Pardoux. Un résultat sur les équations aux dérivées partielles stochastiques et filtrage des processus de diffusion. *C. R. Acad. Sci. Paris Sér. A-B*, 287(16):1065–1068, 1978. MR-0520405
- [15] E. Pardoux. Stochastic PDEs and filtering of diffusion processes. *Stochastics*, 3:127–167, 1979. MR-0553909
- [16] Etienne Pardoux and Shige Peng. Backward doubly stochastic differential equations and systems of quasilinear SPDEs. *Probability Theory and Related Fields*, 98:209–227, 1994. 10.1007/BF01192514. MR-1258986
- [17] Arnaud Porchet, Nizar Touzi, and Xavier Warin. Valuation of power plants by utility indifference and numerical computation. *Math. Methods Oper. Res.*, 70(1):47–75, 2009. MR-2529425
- [18] Philip E. Protter. *Stochastic integration and differential equations*, volume 21 of *Stochastic Modelling and Applied Probability*. Springer-Verlag, Berlin, 2005. Second edition. Version 2.1, Corrected third printing. MR-2273672
- [19] B. Rozovskii. *Stochastic Evolution Equations*. Dordrecht:Reidel, 1991.
- [20] Tomohiro Sasamoto and Herbert Spohn. One-dimensional Kardar-Parisi-Zhang equation: An exact solution and its universality. *Phys. Rev. Lett.*, 104:230602, Jun 2010.
- [21] H. Mete Soner and Nizar Touzi. The dynamic programming equation for second order stochastic target problems. *SIAM J. Control Optim.*, 48(4):2344–2365, 2009. MR-2556347
- [22] Revaz Tevzadze. Solvability of backward stochastic differential equations with quadratic growth. *Stochastic Processes and their Applications*, 118(3):503 – 515, 2008. MR-2389055
- [23] Nizar Touzi. Second order backward SDEs, fully nonlinear PDEs, and applications in finance. In *Proceedings of the International Congress of Mathematicians. Volume IV*, pages 3132–3150, New Delhi, 2010. Hindustan Book Agency. MR-2828009

---

# Electronic Journal of Probability

## Electronic Communications in Probability

---

### Advantages of publishing in EJP-ECP

- Very high standards
- Free for authors, free for readers
- Quick publication (no backlog)

### Economical model of EJP-ECP

- Low cost, based on free software (OJS<sup>1</sup>)
- Non profit, sponsored by IMS<sup>2</sup>, BS<sup>3</sup>, PKP<sup>4</sup>
- Purely electronic and secure (LOCKSS<sup>5</sup>)

### Help keep the journal free and vigorous

- Donate to the IMS open access fund<sup>6</sup> (click here to donate!)
- Submit your best articles to EJP-ECP
- Choose EJP-ECP over for-profit journals

---

<sup>1</sup>OJS: Open Journal Systems <http://pkp.sfu.ca/ojs/>

<sup>2</sup>IMS: Institute of Mathematical Statistics <http://www.imstat.org/>

<sup>3</sup>BS: Bernoulli Society <http://www.bernoulli-society.org/>

<sup>4</sup>PK: Public Knowledge Project <http://pkp.sfu.ca/>

<sup>5</sup>LOCKSS: Lots of Copies Keep Stuff Safe <http://www.lockss.org/>

<sup>6</sup>IMS Open Access Fund: <http://www.imstat.org/publications/open.htm>