

Ricci solitons on 3-dimensional C_{12} -manifolds

Bayour Benaoumeur, Gherici Beldjilali

Abstract. In the present paper we study 3-dimensional C_{12} -manifolds admitting Ricci solitons and generalized Ricci solitons and then we introduce a new generalization of η -Ricci soliton. We give a class of examples.

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1 Introduction

In the classification of D. Chinea and C. Gonzalez [4] of almost contact metric manifolds there is a class C_{12} -manifolds which can be integrable but never normal. Recently, in [7], The authors have developed a systematic study of the curvature of the Chinea-Gonzalez class $C_5 \oplus C_{12}$ and obtain some classification theorems for those manifolds that satisfy suitable curvature conditions. This class is defined by using a certain function α and when this function vanishes the class $C_5 \oplus C_{12}$ reduces to class C_{12} .

Recently, in [2], the authors have study some properties of three dimensional C_{12} -manifolds and construct some relations between class C_{12} and other classes as C_6 and $C_2 \oplus C_9$ or $|C|$.

Here, we investigate these manifolds to construct Ricci soliton and generalized Ricci soliton. It is shown that if in a 3-dimensional C_{12} -manifolds the metric is Ricci soliton, where potential vector field V is collinear with the characteristic vector field ξ , then the manifold is η -Einstein. We also prove that an η -Einstein 3-dimensional C_{12} -manifold with

$$S = \mu g + \sigma \eta \otimes \eta \quad \mu + \sigma = -\operatorname{div} \psi \quad V = \beta \xi \quad \text{and} \quad \operatorname{grad} \beta = \beta \psi - \sigma \xi$$

admits a Ricci soliton. On the other hand, it is shown that any 3-dimensional C_{12} -manifold with $|\psi|^2 - 2\operatorname{div} \psi - \frac{r}{2} = 0$ satisfies the generalized Ricci soliton equation.

This paper is organized in the following way:

Section 2, is devoted to some basic definitions for 3-dimensional C_{12} -manifold. In Section 3, we obtain some results for a 3-dimensional C_{12} -manifold admitting Ricci soliton. In the last section, we present a study on 3-dimensional C_{12} -manifold which

satisfies the generalized Ricci soliton equation and we give concrete examples. Finally, we introduce a generalization of η -Ricci soliton and we prove the existence through several examples.

2 Preliminaries

The notion of Ricci soliton was introduced by Hamilton [10] in 1982. A Ricci soliton is a natural generalization of an Einstein metric. A pseudo-Riemannian manifold (M, g) is called a Ricci soliton if it admits a smooth vector field V (potential vector field) on M such that

$$(2.1) \quad (\mathcal{L}_V g)(X, Y) + 2S(X, Y) + 2\lambda g(X, Y) = 0,$$

where $\mathcal{L}_X g$ is the Lie-derivative of g along X given by:

$$(2.2) \quad (\mathcal{L}_X g)(Y, Z) = g(\nabla_Y X, Z) + g(\nabla_Z X, Y),$$

λ is a constant and X, Y are arbitrary vector fields on M .

A Ricci soliton is said to be shrinking, steady or expanding according to λ being negative, zero or positive, respectively. It is obvious that a trivial Ricci soliton is an Einstein manifold with V zero or Killing.

The generalized Ricci soliton equation in Riemannian manifold (M, g) is defined by (see [12]):

$$(2.3) \quad \mathcal{L}_X g = -2c_1 X^\flat \odot X^\flat + 2c_2 S + 2\lambda g,$$

where $X^\flat(Y) = g(X, Y)$ and $c_1, c_2, \lambda \in \mathbb{R}$.

Equation (2.3), is a generalization of Killing's equation ($c_1 = c_2 = \lambda = 0$), Equation for homotheties ($c_1 = c_2 = 0$), Ricci soliton ($c_1 = 0, c_2 = -1$), Cases of Einstein-Weyl ($c_1 = 1, c_2 = \frac{-1}{n-2}$), Metric projective structures with skew-symmetric Ricci tensor in projective class ($c_1 = 1, c_2 = \frac{-1}{n-1}, \lambda = 0$), Vacuum near-horizon geometry equation ($c_1 = 1, c_2 = \frac{1}{2}$), and is also a generalization of Einstein manifolds (For more details, see [1], [5], [8], [9], [12]).

An odd-dimensional Riemannian manifold (M^{2n+1}, g) is said to be an almost contact metric manifold if there exist on M a $(1, 1)$ -tensor field φ , a vector field ξ (called the structure vector field) and a 1-form η such that

$$(2.4) \quad \eta(\xi) = 1, \varphi^2(X) = -X + \eta(X)\xi \quad \text{and} \quad g(\varphi X, \varphi Y) = g(X, Y) - \eta(X)\eta(Y),$$

for any vector fields X, Y on M . In particular, in an almost contact metric manifold we also have $\varphi\xi = 0$ and $\eta \circ \varphi = 0$.

The fundamental 2-form ϕ is defined by $\phi(X, Y) = g(X, \varphi Y)$. It is known that the almost contact structure (φ, ξ, η) is said to be normal if and only if

$$(2.5) \quad N^{(1)}(X, Y) = N_\varphi(X, Y) + 2d\eta(X, Y)\xi = 0,$$

for any X, Y on M , where N_φ denotes the Nijenhuis torsion of φ , given by

$$(2.6) \quad N_\varphi(X, Y) = \varphi^2[X, Y] + [\varphi X, \varphi Y] - \varphi[\varphi X, Y] - \varphi[X, \varphi Y].$$

Given an almost contact structure, one can associate in a natural manner an almost CR-structure $(\mathcal{D}, \varphi|_{\mathcal{D}})$, where $\mathcal{D} := Ker(\eta) = Im(\varphi)$ is the distribution of rank $2n$ transversal to the characteristic vector field ξ . If this almost CR-structure is integrable (i.e., $N_{\varphi} = 0$) the manifold M^{2n+1} is said to be CR-integrable. It is known that normal almost contact manifolds are CR-manifolds.

In the classification of D. Chinea and C. Gonzalez [4], the almost contact metric structures have been completely classified. The $C_5 \oplus C_{12}$ class was recently discussed by S. de Candia and M. Falcitelli [7]. We just recall the defining relations of $C_5 \oplus C_{12}$ class, which will be used in this study.

The $C_5 \oplus C_{12}$ -manifolds can be characterized by:

$$(2.7) \quad \begin{aligned} (\nabla_X \varphi)Y &= \alpha(g(\varphi X, Y)\xi - \eta(Y)\varphi X) \\ &\quad - \eta(X)((\nabla_{\xi}\eta)(\varphi Y)\xi + \eta(Y)\varphi\nabla_{\xi}\xi). \end{aligned}$$

It is known that any almost contact metric manifold (φ, ξ, η, g) from $C_5 \oplus C_{12}$ class satisfies (see [7])

$$(2.8) \quad \begin{cases} \nabla_X \xi = -\alpha\varphi^2 X + \eta(X)\nabla_{\xi}\xi, \\ d\eta = \eta \wedge \nabla_{\xi}\eta, \\ d(\nabla_{\xi}\eta) = -(\alpha\nabla_{\xi}\eta + \nabla_{\xi}(\nabla_{\xi}\eta)) \wedge \eta, \end{cases}$$

where $\dim M = 2n + 1$ and $\alpha = -\frac{1}{2n}\delta\eta$. Furthermore, if $\dim M \geq 5$, the Lee form of M is $\omega = -\alpha\eta$ and it is closed. Applying (2.8), one has

$$(2.9) \quad d\alpha = \xi(\alpha)\eta + \alpha\nabla_{\xi}\eta.$$

In this paper, we will focus on the class C_{12} . So, putting $\alpha = 0$, $\omega = -(\nabla_{\xi}\xi)^{\flat} = -\nabla_{\xi}\eta$ and if ψ is the vector field given by $\omega(X) = g(X, \psi)$ for all X vector field on M , from formula (2.7) M is of class C_{12} if and only if

$$(2.10) \quad (\nabla_X \varphi)Y = \eta(X)(\omega(\varphi Y)\xi + \eta(Y)\varphi\psi).$$

Moreover, from (2.8) it follow,

$$(2.11) \quad \begin{cases} \nabla_X \xi = -\eta(X)\psi, \\ d\eta = \omega \wedge \eta, \\ d\omega = 0. \end{cases}$$

Notice that $\nabla_{\xi}\xi = -\psi$.

In [2], we have given a characterization of class C_{12} as follows:

Theorem 2.1. *An almost contact metric manifold is of class C_{12} if and only if there exists a 1-form ω such that*

$$(2.12) \quad d\eta = \omega \wedge \eta \quad d\phi = 0 \quad \text{and} \quad N_{\varphi} = 0.$$

Now, we denote by R, S, r the curvature tensor, the Ricci curvature and the scalar curvature respectively, which are defined for all $X, Y, Z \in \mathfrak{X}(M)$ by

$$(2.13) \quad R(X, Y)Z = \nabla_X \nabla_Y Z - \nabla_Y \nabla_X Z - \nabla_{[X, Y]}Z,$$

$$(2.14) \quad S(X, Y) = \sum_{i=1}^{2n+1} g(R(e_i, X)Y, e_i),$$

$$(2.15) \quad r = \sum_{i=1}^{2n+1} S(e_i, e_i),$$

with $\{e_1, \dots, e_{2n+1}\}$ is a local orthonormal basis . The divergence of a vector field X on M is defined by:

$$(2.16) \quad \operatorname{div}\psi = \sum_{i=1}^{2n+1} g(\nabla_{e_i}\psi, e_i).$$

(For more details of previous definitions, see for example [11]). Then, from Corollary 3.1 of [7] we have,

$$(2.17) \quad R(X, Y)\xi = -2d\eta(X, Y)\psi - \eta(Y)\nabla_X\psi + \eta(X)\nabla_Y\psi,$$

$$(2.18) \quad S(X, \xi) = -\eta(X)\operatorname{div}\psi.$$

Proposition 2.2. *In a 3-dimensional C_{12} -manifold, Ricci tensor and curvature tensor are given respectively by*

$$(2.19) \quad \begin{aligned} S(X, Y) &= \left(\frac{r}{2} + \operatorname{div}\psi\right)g(X, Y) + (|\psi|^2 - 2\operatorname{div}\psi - \frac{r}{2})\eta(X)\eta(Y) \\ &- \omega(X)\omega(Y) - g(\nabla_X\psi, Y), \end{aligned}$$

and

$$(2.20) \quad \begin{aligned} R(X, Y)Z &= (|\psi|^2 - 2\operatorname{div}\psi - \frac{r}{2})\eta(Z)(\eta(Y)X - \eta(X)Y) \\ &- g(Y, Z)\left(\omega(X)\psi + \nabla_X\psi - (2\operatorname{div}\psi + \frac{r}{2})X\right) \\ &+ g(X, Z)\left(\omega(Y)\psi + \nabla_Y\psi - (2\operatorname{div}\psi + \frac{r}{2})Y\right) \\ &+ (|\psi|^2 + 2\operatorname{div}\psi - \frac{r}{2})\left(g(Y, Z)\eta(X) - g(X, Z)\eta(Y)\right)\xi \\ &- \omega(Z)(\omega(Y)X - \omega(X)Y) + g(\nabla_X\psi, Z)Y - g(\nabla_Y\psi, Z)X. \end{aligned}$$

Proof. Suppose that $(M, \varphi, \xi, \psi, \eta, \omega, g)$ is a 3-dimensional C_{12} -manifold. Setting $Y = Z = \xi$ in the well known formula (which holds for any 3-dimensional Riemannian manifold [3]):

$$(2.21) \quad \begin{aligned} R(X, Y)Z &= g(Y, Z)QX - g(X, Z)QY + S(Y, Z)X - S(X, Z)Y \\ &- \frac{r}{2}(g(Y, Z)X - g(X, Z)Y), \end{aligned}$$

where Q is the Ricci operator defined by

$$(2.22) \quad S(X, Y) = g(QX, Y).$$

We get

$$(2.23) \quad R(X, \xi)\xi = QX - (\operatorname{div}\psi)X + 2(\operatorname{div}\psi)\eta(X)\xi + \frac{r}{2}\varphi^2 X.$$

Again, Setting $Y = \xi$ in formula (2.17), we obtain

$$(2.24) \quad R(X, \xi)\xi = -g(\nabla_\xi \xi, X)\psi - \nabla_X \psi + \eta(X)\nabla_\xi \psi.$$

On the other hand, we have

$$\begin{aligned} 2d\omega(\xi, X) = 0 &\Leftrightarrow g(\nabla_\xi \psi, X) = g(\nabla_X \psi, \xi) \\ &= -g(\psi, \nabla_X \xi) \\ &= \omega(\psi)\eta(X), \end{aligned}$$

which gives

$$(2.25) \quad \nabla_\xi \psi = \omega(\psi)\xi.$$

So, using (2.11) and (2.25) in formula (2.24) we get

$$(2.26) \quad R(X, \xi)\xi = -\omega(X)\psi - \nabla_X \psi + |\psi|^2 \eta(X)\xi.$$

In view of (2.23) and (2.26), we obtain

$$(2.27) \quad QX = -\omega(X)\psi - \nabla_X \psi + (\operatorname{div}\psi + \frac{r}{2})X + (|\psi|^2 - 2\operatorname{div}\psi - \frac{r}{2})\eta(X)\xi.$$

Finally, equation (2.19) follows from (2.27) and (2.22). Using (2.22) and (2.27) in (2.21), the curvature tensor in a 3-dimensional C_{12} -manifold is given by (2.20). \square

Example 2.1. We denote the Cartesian coordinates in a 3-dimensional Euclidean space \mathbb{R}^3 by (x, y, z) and define a symmetric tensor field g by

$$g = e^{2f} \begin{pmatrix} \rho^2 + \tau^2 & 0 & -\tau \\ 0 & \rho^2 & 0 \\ -\tau & 0 & 1 \end{pmatrix},$$

where $f = f(y)$, $\tau = \tau(x)$ and $\rho = \rho(x, y)$ are functions on \mathbb{R}^3 with $f' = \frac{\partial f}{\partial y}$. Further, we define an almost contact metric (φ, ξ, η) on \mathbb{R}^3 by

$$\varphi = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & -\tau & 0 \end{pmatrix}, \quad \xi = e^{-f} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \quad \eta = e^f(-\tau, 0, 1).$$

The fundamental 1-form η and the 2-form ϕ have the forms,

$$\eta = e^f(dz - \tau dx) \quad \text{and} \quad \phi = -2\rho^2 e^{2f} dx \wedge dy,$$

and hence

$$d\eta = f' e^f (\tau dx \wedge dy + dy \wedge dz),$$

$$d\phi = 0.$$

By a direct computation the non trivial components of $N_{kj}^{(1)i}$ are given by

$$N_{12}^{(1)3} = \tau f', \quad N_{23}^{(1)3} = f'.$$

But, $\forall i, j, k \in \{1, 2, 3\}$

$$(N_\varphi)_{kj}^i = 0,$$

implying that the structure (φ, ξ, η, g) is CR-integrable.

Therefore, to continue studying this example, it suffices to take $f' \neq 0$ to ensure that the structure is CR-integrable not normal.

In order to define the closed 1-form ω , putting $\omega = adx + bdy + cdz$ where a, b and c are functions on \mathbb{R}^3 , and using formulas $d\eta = \omega \wedge \eta$ and $\omega(\xi) = 0$, we can check that is very simply as follows:

$$(2.28) \quad \omega = f' dy,$$

notice that $d\omega = 0$.

Knowing that ω is the g -dual of ψ i.e. $\omega(X) = g(X, \psi)$, we have immediately that

$$(2.29) \quad \psi = \frac{f'}{\rho^2} e^{-2f} \frac{\partial}{\partial y}.$$

Thus, $(\varphi, \xi, \psi, \eta, \omega, g)$ becomes a C_{12} structure on \mathbb{R}^3 .

Now we have

$$\left\{ e_1 = \frac{e^{-f}}{\rho} \left(\frac{\partial}{\partial x} + \tau \frac{\partial}{\partial z} \right), \quad e_2 = \frac{e^{-f}}{\rho} \frac{\partial}{\partial y}, \quad e_3 = \xi = e^{-f} \frac{\partial}{\partial z} \right\}$$

form an orthonormal basis. To verify result in formula (2.10), the non zero components of the Levi-Civita connection corresponding to g are given by:

$$\nabla_{e_1} e_1 = -\frac{(f' \rho + \rho_2)}{\rho^2 e^f} e_2, \quad \nabla_{e_1} e_2 = \frac{(f' \rho + \rho_2)}{\rho^2 e^f} e_1,$$

$$\nabla_{e_2} e_1 = \frac{\rho_1}{\rho^2 e^f} e_2, \quad \nabla_{e_2} e_2 = -\rho_1 \frac{1}{\rho^2 e^f} e_1,$$

$$\nabla_{e_3} e_2 = \frac{f'}{\rho e^f} e_3, \quad \nabla_{e_3} e_3 = -\frac{f'}{\rho e^f} e_2.$$

Then, one can easily check that for all $i, j \in \{1, 2, 3\}$

$$\begin{aligned} (\nabla_{e_i} \varphi) e_j &= \nabla_{e_i} \varphi e_j - \varphi \nabla_{e_i} e_j \\ &= \eta(e_i) (\omega(\varphi e_j) \xi + \eta(e_j) \varphi \psi). \end{aligned}$$

3 Ricci soliton

In this section, we consider a 3-dimensional C_{12} -manifold M admitting a Ricci soliton defined by (2.1). Let V be a pointwise collinear vector field with the structure vector field ξ , that is $V = \beta\xi$, where β is a function on M . From (2.1) we write

$$(3.1) \quad g(\nabla_X \beta \xi, Y) + g(\nabla_Y \beta \xi, X) + 2S(X, Y) + 2\lambda g(X, Y) = 0,$$

for all X and Y vector fields on M . Then, we have

$$\begin{aligned} X(\beta)\eta(Y) + \beta g(\nabla_X \xi, Y) + Y(\beta)\eta(X) \\ + \beta g(\nabla_Y \xi, X) + 2S(X, Y) + 2\lambda g(X, Y) = 0, \end{aligned}$$

which implies

$$(3.2) \quad \begin{aligned} X(\beta)\eta(Y) - \beta\eta(X)\omega(Y) + Y(\beta)\eta(X) \\ - \beta\eta(Y)\omega(X) + 2S(X, Y) + 2\lambda g(X, Y) = 0, \end{aligned}$$

by virtue of (2.11). By putting $Y = \xi$ in (3.2) and using (2.18) we obtain

$$(3.3) \quad X(\beta) - \beta\omega(X) + (\xi(\beta) - 2\operatorname{div}\psi + 2\lambda)\eta(X) = 0.$$

Taking $X = \xi$ in the previous equation gives

$$(3.4) \quad \xi(\beta) = \operatorname{div}\psi - \lambda.$$

If we replace (3.4) in (3.3), we get

$$(3.5) \quad X(\beta) = \beta\omega(X) + (\operatorname{div}\psi - \lambda)\eta(X),$$

again, if we replace (3.5) in (3.2), we obtain

$$(3.6) \quad S(X, Y) = -\lambda g(X, Y) + (\lambda - \operatorname{div}\psi)\eta(X)\eta(Y),$$

for all X and Y vector fields on M . Hence we have

Theorem 3.1. *Let $(M, \varphi, \xi, \psi, \eta, \omega, g)$ be a 3-dimensional C_{12} -manifold. If M admits a Ricci soliton and V is pointwise collinear with the structure vector field ξ , then M is an η -Einstein manifold.*

In addition, if $\lambda = \operatorname{div}\psi = \text{constant}$ then M is an Einstein manifold.

Let assume the converse, that is, let M be a 3-dimensional η -Einstein C_{12} -manifold with $V = \beta\xi$. Then we can write

$$(3.7) \quad S(X, Y) = ag(X, Y) + b\eta(X)\eta(Y),$$

where a and b are scalars and X, Y are vector fields on M . From (2.2) we have

$$\begin{aligned} (\mathcal{L}_V g)(Y, Y) &= g(\nabla_X V, Y) + g(\nabla_Y V, X) \\ &= X(\beta)\eta(Y) + Y(\beta)\eta(X) - \beta\eta(X)\omega(Y) - \beta\eta(Y)\omega(X), \end{aligned}$$

which implies that

$$\begin{aligned} (\mathcal{L}_V g)(X, Y) + 2S(X, Y) + 2\lambda g(X, Y) &= 2(a + \lambda)g(X, Y) \\ &+ \eta(X)(b\eta(Y) - \beta\omega(Y) + Y(\beta)) \\ &+ \eta(Y)(b\eta(X) - \beta\omega(X) + X(\beta)). \end{aligned}$$

From the previous equation it is obvious that M admits a Ricci soliton (g, V, λ) if

$$a + \lambda = 0 \quad \text{and} \quad b\eta(Y) - \beta\omega(Y) + Y(\beta) = 0.$$

Equating the right hand sides of (3.7) and (2.18) and taking $X = Y = \xi$ gives

$$a + b = -\text{div}\psi,$$

Thus, we get

Theorem 3.2. *Let $(M, \varphi, \xi, \psi, \eta, \omega, g)$ be a 3-dimensional C_{12} -manifold with $\text{div}\psi$ is constant. If M is an η -Einstein manifold with $S = ag + b\eta \otimes \eta$ and $a + b = -\text{div}\psi$, then the manifold admits a Ricci soliton $(g, \beta\xi, a)$ with $\text{grad}\beta = \beta\psi - b\xi$.*

4 Generalized Ricci soliton

In this section we will study the generalized Ricci soliton equation (2.3) on a C_{12} -manifold of dimension three. let's start with our main result

Theorem 4.1. *Any three-dimensional C_{12} -manifold satisfies the generalized Ricci soliton equation (2.3) with $X = \psi$, $c_1 = 1$, $c_2 = -1$ and $\lambda = |\psi|^2 - \text{div}\psi$ if and only if*

$$(4.1) \quad |\psi|^2 - 2\text{div}\psi - \frac{r}{2} = 0.$$

Proof. Suppose that $(M, \varphi, \xi, \psi, \eta, \omega, g)$ is a C_{12} -manifold of dimension three which satisfies the generalized Ricci soliton equation (2.3) with $X = \psi$, that is, for all $Y, Z \in \Gamma(TM)$

$$(4.2) \quad (\mathcal{L}_\psi g)(Y, Z) = -2c_1\omega(Y)\omega(Z) + 2c_2S(Y, Z) + 2\lambda g(Y, Z).$$

Since ω is closed then $g(\nabla_Y \psi, Z) = g(\nabla_Z \psi, Y)$. Therefore, we can express the generalized soliton equation as

$$(4.3) \quad \nabla_Y \psi = -c_1\omega(Y)\psi + c_2QY + \lambda Y.$$

Now, from (2.27) we get

$$(4.4) \quad \nabla_Y \psi = -\omega(Y)\psi - QY + (\text{div}\psi + \frac{r}{2})Y + (|\psi|^2 - 2\text{div}\psi - \frac{r}{2})\eta(Y)\xi.$$

In view of (4.4) and (4.3) the proof is complete. \square

Proposition 4.2. *Let $(M, \varphi, \xi, \psi, \eta, \omega, g)$ be a C_{12} -manifold of dimension three which satisfies the generalized Ricci soliton equation (2.3) with $X = \psi$. If $|\psi| = 1$ then $r = \text{constant}$.*

Proof. The proof is direct, it suffices to use Theorem 4.1. \square

Example 4.1. Let's go back to the class of the previous examples . With simple but long calculations, we can get the following:

$$|\psi|^2 = \frac{f'^2}{\rho^2} e^{-2f}, \quad \operatorname{div}\psi = \frac{e^{-2f}}{\rho^2} (f'^2 + f''), \quad \lambda = \frac{f'' e^{-2f}}{\rho^2},$$

$$r = \frac{2e^{-2f}}{\rho^4} (\rho_1^2 - \rho\rho_{11} + \rho_2^2 - \rho\rho_{22} - 2f''\rho^2 - f'^2\rho^2),$$

where $\rho_i = \frac{\partial\rho}{\partial x_i}$. Then, the condition (4.1) gives the following differential equation

$$(4.5) \quad \rho_1^2 - \rho\rho_{11} + \rho_2^2 - \rho\rho_{22} = 0.$$

Henceforth, we can construct a non-trivial generalized Ricci soliton.

For example:

$$(1): \rho = 1, \quad \lambda = f'' e^{-2f},$$

$$\begin{aligned} \nabla_{e_1}\psi &= e^{-2f} f'^2 e_1, & Qe_1 &= -e^{-2f} (f'' + f'^2) e_1, \\ \nabla_{e_2}\psi &= e^{-2f} (-f'^2 + f'') e_2, & Qe_2 &= -2e^{-2f} f'' e_2, \\ \nabla_{e_3}\psi &= e^{-2f} f'^2 e_3, & Qe_3 &= -e^{-2f} (f'^2 + f'') e_3. \end{aligned}$$

$$(2): \rho = e^y, \quad \lambda = f'' e^{-2(y+f)},$$

$$\begin{aligned} \nabla_{e_1}\psi &= e^{-2(y+f)} (f'(f' + 1)) e_1, & Qe_1 &= -e^{-2(y+f)} (f'' + f'^2 + f') e_1, \\ \nabla_{e_2}\psi &= e^{-2(y+f)} (-f'^2 + f'' - f') e_2, & Qe_2 &= -2e^{-2(y+f)} (f'' + f'^2) e_2, \\ \nabla_{e_3}\psi &= e^{-2(y+f)} f'^2 e_3, & Qe_3 &= -e^{-2(y+f)} (f'^2 + f'') e_3. \end{aligned}$$

$$(3): \rho = e^{-f}, \quad \lambda = 0,$$

$$\begin{aligned} \nabla_{e_1}\psi &= 0, & Qe_1 &= 0, \\ \nabla_{e_2}\psi &= f'' e_2, & Qe_2 &= -(f'^2 + f'') e_2, \\ \nabla_{e_3}\psi &= f'^2 e_3, & Qe_3 &= -(f'^2 + f'') e_3. \end{aligned}$$

Of course, we must choose f so that λ is constant. We can construct further examples of generalized Ricci soliton on a 3-dimensional C_{12} -manifold by the similar way.

At the end of this section, we present the concept of the generalized η -Ricci soliton as a generalization of the η -Ricci soliton given by Cho-Kimura in [6] by the following equation:

$$(4.6) \quad \mathcal{L}_V g + 2S + 2\lambda g + \mu\eta \otimes \eta = 0,$$

where the tensor product notation $(\eta \otimes \eta)(X, Y) = \eta(X)\eta(Y)$ is used and λ, μ are real constants.

The generalized η -Ricci soliton equation in Riemannian manifold (M, g) is defined by:

$$(4.7) \quad \mathcal{L}_X g = -2c_1 X^\flat \odot X^\flat + 2c_2 S + 2\lambda g + \mu\eta \otimes \eta,$$

where $c_1, c_2, \lambda, \mu \in \mathbb{R}$.

With the same reasoning above, we can express formula (4.7) as follows

$$(4.8) \quad \nabla_X \psi = -c_1 \omega(X) \psi + c_2 QX + \lambda X + \mu \eta \otimes \xi.$$

Now, based on equation (4.4), we declare the following result

Theorem 4.3. *Any 3-dimensional C_{12} -manifold satisfies the generalized η -Ricci soliton equation with*

$$c_1 = 1 \quad c_2 = -1 \quad \lambda = |\psi|^2 - \operatorname{div} \psi \quad \text{and} \quad \mu = |\psi|^2 - 2 \operatorname{div} \psi - \frac{r}{2}.$$

Example 4.2. From Example 4.1, we can construct several non-trivial cases, namely:

1) If $f = y$ and $\rho = \frac{4}{e^{2y}-c}$ with $c \in \mathbb{R}$, then we get

$$c_1 = 1, \quad c_2 = -1, \quad \lambda = 0, \quad \text{and} \quad \mu = c.$$

2) If $f = \ln\left(\frac{1}{\sin^2 y}\right)$ and $\rho = c \sin y$, then we get

$$c_1 = 1, \quad c_2 = -1, \quad \lambda = -\frac{2}{c^2}, \quad \mu = -\frac{1}{c^2}.$$

Of course, while taking into account the necessary conditions on f and ρ .

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Authors' address:

Bayour Benaoumeur and Beldjilali Gherici
Laboratory of Quantum Physics and Mathematical Modeling (LPQ3M),
University of Mascara, Algeria.
E-mail addresses: b.bayour@univ-mascara.dz , gherici.beldjilali@univ-mascara.dz