

Classification of cylindrically symmetric static space-times according to their proper homothetic vector fields

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Abstract. A complete classification of cylindrically symmetric static space-times according to their proper homothetic vector field is given by using direct integration technique. Using the above mentioned technique we have shown that a very special class of the above space-times admit proper homothetic vector field. The dimensions of homothetic Lie algebras are 4, 5, 7 and 11.

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§1. Introduction

The aim of this paper is to study all the possibilities when the cylindrically symmetric static space-times admit proper homothetic vector fields by using direct integration technique. Throughout M is representing the four dimensional, connected, hausdorff space-time manifold with Lorentz metric g of signature $(-, +, +, +)$. The curvature tensor associated with g_{ab} , through Levi-Civita connection, is denoted in component form by R^a_{bcd} , and the Ricci tensor components are $R_{ab} = R^c_{acb}$. The usual covariant, partial and Lie derivatives are denoted by a semicolon, a comma and the symbol L , respectively. Round and square brackets denote the usual symmetrization and skew-symmetrization, respectively.

Any vector field X on M can be decomposed as

$$(1.1) \quad X_{a;b} = \frac{1}{2}h_{ab} + F_{ab}$$

where $h_{ab} = h_{ba}$ and $F_{ab} = -F_{ba}$ are symmetric and skew-symmetric tensor on M , respectively. If

$$h_{ab} = \alpha g_{ab}, \quad \alpha \in R$$

equivalent,

$$(1.2) \quad g_{ab,c}X^c + g_{bc}X^c_{,a} + g_{ac}X^c_{,b} = \alpha g_{ab}$$

then X is called a homothetic vector field on M . If X is homothetic and $\alpha \neq 0$ then it is called proper homothetic while $\alpha = 0$ it is Killing [2, 3]. Further consequences and geometrical interpretations of (1.2) are explored in [1, 4]. It also follows from (1.2) that [1, 3]

$$L_X R^a{}_{bcd} = 0 \quad L_X R_{ab} = 0.$$

The Lie algebra of a set of vector fields on a manifold is completely characterized by the structure constants C_{bc}^a given in term of the Lie brackets by

$$(1.3) \quad [X_b, X_c] = C_{bc}^a X_a, \quad C_{bc}^a = -C_{cb}^a,$$

where X_a are the generators. The Lie algebras for homothetic vector fields in term of structure constants are also given.

§2. Main results

Consider a cylindrically symmetric static space time in the usual coordinate system (labeled as (x^0, x^1, x^2, x^3)) with line element [7]

$$(2.4) \quad ds^2 = -e^{v(r)} dt^2 + dr^2 + e^{u(r)} d\theta^2 + e^{w(r)} d\phi^2,$$

where v , u and w are some functions of r only. The linearly independent Killing vector fields are [5, 6] $\frac{\partial}{\partial t}$, $\frac{\partial}{\partial \theta}$, $\frac{\partial}{\partial \phi}$. A vector field X is said to be a homothetic vector field if it satisfies the equation (1.2). One can write (1.2) explicitly using (2.4) as

$$(2.5) \quad v' X^1 + 2X_{,0}^0 = 2c,$$

$$(2.6) \quad X_{,0}^1 - e^v X_{,1}^0 = 0,$$

$$(2.7) \quad e^u X_{,0}^2 - e^v X_{,2}^0 = 0,$$

$$(2.8) \quad e^w X_{,0}^3 - e^v X_{,3}^0 = 0,$$

$$(2.9) \quad X_{,1}^1 = c,$$

$$(2.10) \quad e^u X_{,1}^2 + X_{,2}^1 = 0,$$

$$(2.11) \quad e^w X_{,1}^3 + X_{,3}^1 = 0,$$

$$(2.12) \quad u' X^0 + 2X_{,2}^2 = 2c,$$

$$(2.13) \quad e^w X_{,2}^3 + e^u X_{,3}^2 = 0,$$

$$(2.14) \quad w' X^0 + 2X_{,3}^3 = 2c.$$

Equations (2.6), (2.9), (2.10) and (2.11) give

$$(2.15) \quad \left. \begin{aligned} X^0 &= A_t^1(t, \theta, \phi) \int e^{-v} dr + A^2(t, \theta, \phi) \\ X^1 &= cr + A^1(t, \theta, \phi) \\ X^2 &= A_\theta^1(t, \theta, \phi) \int e^{-u} dr + A^3(t, \theta, \phi) \\ X^3 &= A_\phi^1(t, \theta, \phi) \int e^{-w} dr + A^4(t, \theta, \phi) \end{aligned} \right\},$$

where $A^1(t, \theta, \phi)$, $A^2(t, \theta, \phi)$, $A^3(t, \theta, \phi)$ and $A^4(t, \theta, \phi)$ are functions of integration. In order to determine $A^1(t, \theta, \phi)$, $A^2(t, \theta, \phi)$, $A^3(t, \theta, \phi)$ and $A^4(t, \theta, \phi)$ we need to integrate the remaining six equations. To avoid lengthy calculations here we will only present results for full details see [8].

Case (1) Four independent homothetic vector fields:

In this case the space-time (2.4) takes the form

$$(2.16) \quad ds^2 = -(cr + d_9)^{2(1-\frac{d_5}{c})} dt^2 + dr^2 + (cr + d_9)^{2(1-\frac{d_{10}}{c})} d\theta^2 + (cr + d_9)^{2(1-\frac{d_{12}}{c})} d\phi^2,$$

and homothetic vector fields in this case are

$$(2.17) \quad \left. \begin{aligned} X^0 &= td_5 + d_6 \\ X^1 &= rc + d_9 \\ X^2 &= \theta d_{10} + d_{11} \\ X^3 &= \phi d_{12} + d_{13} \end{aligned} \right\},$$

where $d_5, d_6, d_9, d_{10}, d_{11}, d_{12}, d_{13} \in R$. The above space-time (2.16) admits four independent homothetic vector fields in which three are Killing vector fields which are given in (1.3) and one is proper homothetic vector field which is

$$(2.18) \quad Y^1 = (0, r, 0, 0).$$

The generators in this case are:

$$X_1 = \frac{\partial}{\partial t}, X_2 = \frac{\partial}{\partial \theta}, X_3 = \frac{\partial}{\partial \phi} \text{ and } X_4 = r \frac{\partial}{\partial r}.$$

Here, all the structure constants are zero.

Case (2) Four independent homothetic vector fields:

In this case the space-time (2.4) takes the form

$$(2.19) \quad ds^2 = -(cr + d_9)^{2(1-\frac{d_5}{c})} dt^2 + dr^2 + d\theta^2 + (cr + d_9)^{2(1-\frac{d_{12}}{c})} d\phi^2.$$

Homothetic vector fields in this case are

$$(2.20) \quad \left. \begin{aligned} X^0 &= td_5 + d_6 \\ X^1 &= rc + d_9 \\ X^2 &= \theta c + d_{11} \\ X^3 &= \phi d_{12} + d_{13} \end{aligned} \right\},$$

where $d_5, d_6, d_9, d_{11}, d_{12}, d_{13} \in R$. The above space-time (2.19) admits four independent homothetic vector fields in which three are Killing vector fields and one is proper homothetic vector field which is

$$(2.21) \quad Y^2 = (0, r, \theta, 0).$$

The generators in this case are:

$$X_1 = \frac{\partial}{\partial t}, X_2 = \frac{\partial}{\partial \theta}, X_3 = \frac{\partial}{\partial \phi} \text{ and } X_4 = r \frac{\partial}{\partial r} + \theta \frac{\partial}{\partial \theta}.$$

The only non zero structure constant is: $C_{24}^2 = 1$

Case (3) Five independent homothetic vector fields:

In this case the space-time (2.4) takes the form

$$(2.22) \quad ds^2 = (cr + d_9)^{2(1-\frac{d_8}{c})}(-dt^2 + d\phi^2) + dr^2 + (cr + d_9)^{2(1-\frac{d_{10}}{c})}d\theta^2$$

and homothetic vector fields in this case are

$$(2.23) \quad \left. \begin{aligned} X^0 &= td_8 + \phi d_{13} + d_{15} \\ X^1 &= rc + d_9 \\ X^2 &= \theta d_{10} + d_{11} \\ X^3 &= \phi d_8 + td_{13} + d_{16} \end{aligned} \right\},$$

where $d_8, d_9, d_{10}, d_{11}, d_{13}, d_{15}, d_{16} \in R$. The above space-time (2.22) admits five independent homothetic vector fields in which four are Killing vector fields and one is proper homothetic vector field which is

$$(2.24) \quad Y^3 = (0, r, 0, 0).$$

The generators in this case are:

$$X_1 = \frac{\partial}{\partial t}, X_2 = \frac{\partial}{\partial \theta}, X_3 = \frac{\partial}{\partial \phi}, X_4 = t \frac{\partial}{\partial \phi} + \phi \frac{\partial}{\partial t} \text{ and } X_5 = r \frac{\partial}{\partial r}.$$

The only non zero structure constant are: $C_{14}^3 = C_{34}^1 = 1$.

Case (4) Five independent homothetic vector fields:

In this case the space-time (2.4) takes the form

$$(2.25) \quad ds^2 = (cr + d_9)^{2(1-\frac{d_8}{c})}(-dt^2 + d\phi^2) + dr^2 + d\theta^2$$

and homothetic vector fields in this case are

$$(2.26) \quad \left. \begin{aligned} X^0 &= td_8 + \phi d_{13} + d_{14} \\ X^1 &= rc + d_9 \\ X^2 &= \theta c + d_{17} \\ X^3 &= \phi d_8 + td_{13} + d_{10} \end{aligned} \right\},$$

where $d_8, d_9, d_{10}, d_{13}, d_{14}, d_{17} \in R$. The above space-time (2.25) admits five independent homothetic vector fields in which four are Killing vector fields and one is proper homothetic vector field which is

$$(2.27) \quad Y^4 = (0, r, \theta, 0).$$

The generators in this case are:

$$X_1 = \frac{\partial}{\partial t}, X_2 = \frac{\partial}{\partial \theta}, X_3 = \frac{\partial}{\partial \phi}, X_4 = t \frac{\partial}{\partial \phi} + \phi \frac{\partial}{\partial t} \text{ and } X_5 = r \frac{\partial}{\partial r} + \theta \frac{\partial}{\partial \theta}.$$

The only non zero structure constant are: $C_{14}^3 = C_{25}^2 = C_{34}^1 = 1$.

Case (5) Seven independent homothetic vector fields:

In this case the space-time (2.4) takes the form

$$(2.28) \quad ds^2 = (cr + d_7)^{2(1-\frac{d_6}{c})}(-dt^2 + d\phi^2 + d\theta^2) + dr^2.$$

Homothetic vector fields in this case are

$$(2.29) \quad \left. \begin{aligned} X^0 &= td_6 + \phi d_{11} + \theta d_{10} + d_{12} \\ X^1 &= rc + d_7 \\ X^2 &= \theta d_6 + td_{10} - \phi d_{13} + d_{15} \\ X^3 &= \phi d_6 + td_{11} + \theta d_{13} + d_{14} \end{aligned} \right\},$$

where $d_6, d_7, d_{10}, d_{11}, d_{12}, d_{13}, d_{14}, d_{15} \in R$. The above space-time (2.28) admits seven independent homothetic vector fields in which six are Killing vector fields and one is proper homothetic vector field which is

$$(2.30) \quad Y^5 = (0, r, 0, 0).$$

The generators in this case are:

$$X_1 = \frac{\partial}{\partial t}, X_2 = \frac{\partial}{\partial \theta}, X_3 = \frac{\partial}{\partial \phi}, X_4 = t \frac{\partial}{\partial \phi} + \phi \frac{\partial}{\partial t}, X_5 = t \frac{\partial}{\partial \theta} + \theta \frac{\partial}{\partial t}, X_6 = \theta \frac{\partial}{\partial \phi} - \phi \frac{\partial}{\partial \theta}$$

and

$$X_7 = r \frac{\partial}{\partial r}.$$

The only non zero structure constant are: $C_{14}^3 = C_{15}^2 = C_{25}^1 = C_{26}^3 = C_{34}^1 = C_{56}^4 = 1$ and $C_{36}^2 = C_{45}^6 = C_{46}^5 = -1$.

Case (6) Eleven independent homothetic vector fields:

In this case the above space-time (2.4) becomes Minkowski space-time

$$(2.31) \quad ds^2 = -dt^2 + dr^2 + d\theta^2 + d\phi^2$$

homothetic vector fields in this case are

$$(2.32) \quad \left. \begin{aligned} X^0 &= tc + rd_{10} + \phi d_{15} + \theta d_{16} + d_{17} \\ X^1 &= rc + td_{10} + \phi d_{11} + \theta d_{12} + d_{13} \\ X^2 &= \theta c - rd_{12} + td_{16} - \phi d_{18} + d_{14} \\ X^3 &= \phi c - rd_{11} + td_{15} + \theta d_{18} + d_{19} \end{aligned} \right\},$$

where $d_{10}, d_{11}, d_{12}, d_{13}, d_{14}, d_{15}, d_{16}, d_{17}, d_{18}, d_{19} \in R$. The above space-time (2.31) admits eleven independent homothetic vector fields in which ten are Killing vector fields and one is proper homothetic vector field which is

$$(2.33) \quad Y^6 = (t, r, \theta, \phi).$$

The generators in this case are:

$$\begin{aligned} X_1 &= \frac{\partial}{\partial t}, X_2 = \frac{\partial}{\partial r}, X_3 = \frac{\partial}{\partial \theta}, X_4 = \frac{\partial}{\partial \phi}, X_5 = t \frac{\partial}{\partial r} + r \frac{\partial}{\partial t}, \\ X_6 &= t \frac{\partial}{\partial \phi} + \phi \frac{\partial}{\partial t}, X_7 = \theta \frac{\partial}{\partial t} + t \frac{\partial}{\partial \theta}, X_8 = \phi \frac{\partial}{\partial r} - r \frac{\partial}{\partial \phi}, X_9 = \theta \frac{\partial}{\partial r} - r \frac{\partial}{\partial \theta}, \\ X_{10} &= \theta \frac{\partial}{\partial \phi} - \phi \frac{\partial}{\partial \theta} \text{ and } X_{11} = t \frac{\partial}{\partial t} + r \frac{\partial}{\partial r} + \theta \frac{\partial}{\partial \theta} + \phi \frac{\partial}{\partial \phi}. \end{aligned}$$

The only non zero structure constant are: $C_{15}^1 = C_{16}^4 = C_{17}^3 = C_{11}^1 = C_{25}^1 = C_{21}^2 = C_{37}^1 = C_{39}^2 = C_{31}^4 = C_{31}^3 = C_{46}^1 = C_{48}^2 = C_{41}^4 = C_{68}^5 = C_{79}^5 = C_{71}^6 = C_{89}^1 = C_{91}^8 = 1$ and $C_{28}^4 = C_{29}^3 = C_{41}^3 = C_{56}^8 = C_{57}^9 = C_{58}^6 = C_{59}^7 = C_{67}^1 = C_{61}^7 = C_{81}^9 = -1$.

SUMMARY

In this paper a study of cylindrically symmetric static space-times according to their proper homothetic vector fields is given by using the direct integration technique. From the above study we obtain the following:

- (i) The space-times which admit four independent homothetic vector fields are given in equations (2.16) and (2.19) (see for details cases (1) and (2)).
- (ii) The space-times which admit five independent homothetic vector fields are given in equations (2.22) and (2.25) (see for details cases (3) and (4)).
- (iii) The space-time which admits seven independent homothetic vector fields are given in equation (2.28) (see for details case (5)).
- (iv) The space-time which admits eleven independent homothetic vector fields are given in equation (2.31) (see for details case (6)).

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