

Global and Blowup Solutions of Quasilinear Parabolic Equation with Critical Sobolev Exponent and Lower Energy Initial Value

ZHONG TAN^{a,*} and ZHENG-AN YAO^b

^a*Department of Mathematics, Xiamen University, Fujian Xiamen 361005, China;*

^b*Department of Mathematics, Zhongshan University Guangdong Guangzhou 510275, China*

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In this paper, by means of the energy method, we first study the existence and asymptotic estimates of global solution of quasilinear parabolic equations involving p -Laplacian ($p > 2$) and critical Sobolev exponent and *lower energy* initial value in a bounded domain in R^N ($N \geq 3$), and also study the sufficient conditions of finite time blowup of local solution by the classical concave method. Finally, we study the asymptotic behavior of any global solutions $u(x, t; u_0)$ which may possess *high energy* initial value function $u_0(x)$. We can prove that there exists a time subsequence $\{t_n\}$ such that the asymptotic behavior of $u(x, t_n; u_0)$ as $t_n \rightarrow \infty$ is similar to the Palais–Smale sequence of stationary equation of the above parabolic problem.

Keywords: Quasilinear parabolic equation; Critical Sobolev exponent; Lower energy initial value; Asymptotic estimates; Finite time blow up

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1 INTRODUCTION

In this paper we are concerned with the asymptotic estimates of global solutions, and blow-up of local solutions of quasilinear parabolic

* Corresponding author.

equations of the following form:

$$\begin{aligned} u_t - \Delta_p u &= |u|^{q-2}u, & (x, t) \in \Omega \times (0, T), \\ u(x, t) &= 0, & (x, t) \in \partial\Omega \times (0, T), \\ u(x, 0) &= u_0(x), & u_0(x) \geq 0, u_0(x) \neq 0, \end{aligned} \quad (1.1)$$

with *lower energy* initial value, and asymptotic behavior of any global solutions which may possess *high energy* initial value function. Here $\Delta_p u = \operatorname{div}(|\nabla u|^{p-2} \nabla u)$, $2 < p < N$, $q = p^* = pN/(N-p)$ is the critical Sobolev exponent. Ω is a bounded domain in $R^N (N \geq 3)$ with smooth boundary $\partial\Omega$.

Equation (1.1) is a class of degenerate parabolic equations and appears in the relevant theory of nonNewtonian fluids [1]. For the case of $p = 2$, various authors have derived sufficient conditions for the existence and asymptotic behavior of global solutions of (1.1) [2,6,10,11]. For the case of $p \neq 2$, Tsutsumi [19], Ishii [9], Otani [16], Nakao [14] have studied the existence and the asymptotic behavior of global solution with $q < p^*$. In [15], Nakao considered the problem with critical or supercritical nonlinear and the condition imposed on the initial data is $u_0 \in W_0^{1,p}(\Omega) \cap L^\infty(\Omega)$ or $u_0 \in W_0^{1,p}(\Omega) \cap L^{p_0}(\Omega) (p_0 > p^*)$, and obtained precise estimates about the asymptotic behavior as $t \rightarrow \infty$. The first object of this paper is to relax this additional condition of $u_0 \in L^\infty(\Omega)$ or $u \in L^{p_0}(\Omega)$ and to study the time-asymptotic estimates and finite time blow up of (1.1) with *lower energy* initial value. The second object of this paper is to consider the asymptotic behavior of any global solution which may possess *high energy* initial value function. We can prove that there exists a subsequence $\{t_n\}$ such that the asymptotic behavior of $u(t_n)$ as $t_n \rightarrow \infty$ is similar to the Palais-Smale sequence of stationary equation of (1.1).

To state the main idea, we first give some useful definitions and notations.

Denotes the usual Sobolev space by $W_0^{1,p}(\Omega)$, endowed with the norm $\|\nabla u\|_p = (\int_\Omega |\nabla u|^p dx)^{1/p}$, denote the norm of $L^r(\Omega)$ by $\|\cdot\|_r$. Denote $\Omega \times (0, T)$ by Q_T .

DEFINITION 1.1 *We say that a function u is a solution of (1.1) in Q_T iff*

$$\begin{aligned} u &\in L^\infty(0, T; W_0^{1,p}(\Omega)), \\ u_t &\in L^2(Q_T) = L^2(0, T; L^2(\Omega)), \end{aligned}$$

and satisfies (1.1) in the distribution sense. If $T = \infty$, u is called global solution. We always denote by $u(x, t; u_0)$ the solution with initial value $u_0(x)$.

Let S be the best constant for the Sobolev embedding $W_0^{1,p}(\Omega) \subset L^{p^*}(\Omega)$ which defined as follows:

$$S = \inf_{\substack{u \in W_0^{1,p}(\Omega) \\ \|u\|_{p^*} = 1}} \|\nabla u\|_p^p.$$

Remark 1.1 Let S be the best constant for the Sobolev embedding of $W_0^{1,p}(\Omega) \subset L^{p^*}(\Omega)$. Then

- (a) S is independent of Ω and depends only on N .
- (b) The infimum S is never achieved when Ω is a bounded domain.

The proof can be found in Talenti [18].

Denote the energy function of (1.1) by

$$J(u) = \frac{1}{p} \int_{\Omega} |\nabla u|^p \, dx - \frac{1}{p^*} \int_{\Omega} |u|^{p^*} \, dx.$$

DEFINITION 1.2 We say that a function $u_0(x)$ possesses lower energy if

$$J(u_0) < \frac{1}{N} S^{N/p}.$$

where S is the best Sobolev constant.

Remark 1.2 The value $(1/N)S^{N/p}$ is the energy of the unique positive radial solution of the quasilinear elliptic equation

$$\begin{aligned} -\Delta_p u &= |u|^{p^*-2} u, \quad x \in \mathbb{R}^N, \\ u(|x|) &\rightarrow 0, \quad \text{as } |x| \rightarrow \infty, \end{aligned} \tag{1.2}$$

to the energy functional

$$J^\infty(u) = \frac{1}{p} \int_{\mathbb{R}^N} |\nabla u|^p \, dx - \frac{1}{p^*} \int_{\mathbb{R}^N} |u|^{p^*} \, dx.$$

A number of authors have studied the perturbation problem of (1.2) (or in $p=2$ and bounded domain) comparing the energy functional of perturbation problem with $(1/N)S^{N/p}$ (e.g. [3,7,8,13,20]).

Now we can state the main results of the first object. First we consider the case: $J(u_0) \leq 0$, we have

THEOREM 1.1 *Let $u_0(x)$ be a lower energy initial value and $J(u_0) \leq 0$. Then $u(x, t; u_0)$ blowup in finite time.*

Now we consider the case of positive energy, i.e. $0 < J(u_0) < (1/N)S^{N/p}$, we have

THEOREM 1.2 *Let $u_0(x) (\neq 0)$ be a lower energy initial value.*

(1) *If $\int_{\Omega} |u_0|^{p^*} dx < S^{N/p}$, then (1.1) has a global solution $u(x, t; u_0)$.*

Moreover

$$\|\nabla u(t)\|_p^p = O(t^{-2/(p-2)}), \quad \text{as } t \rightarrow \infty. \quad (1.3)$$

and

$$\|u\|_2^2 = O(t^{-2/(p-2)}), \quad \text{as } t \rightarrow \infty, \quad (1.4)$$

(2) *If $\int_{\Omega} |u_0|^{p^*} dx \geq S^{N/p}$, then the local solution blows up in finite time.*

Remark 1.3 Obviously, if $\int_{\Omega} |u_0|^{p^*} dx < S^{N/p}$, then $J(u_0) > 0$. Indeed, if

$$\int_{\Omega} |u_0|^{p^*} dx < S^{N/p},$$

then

$$\int_{\Omega} |\nabla u_0|^p dx > \int_{\Omega} |u_0|^{p^*} dx.$$

Thus, from $u_0(x) \neq 0$, we have

$$J(u_0) = \frac{1}{p} \int_{\Omega} |\nabla u_0|^p dx - \frac{1}{p^*} \int_{\Omega} |u_0|^{p^*} dx > \frac{1}{N} \int_{\Omega} |\nabla u_0|^p dx > 0.$$

Now we state the main results of the second object:

THEOREM 1.3 *If $u(x, t; u_0)$ is a global solution of (1.1), and uniformly bounded in $W_0^{1,p}(\Omega)$ with respect to t , then, for any subsequence $t_n \rightarrow \infty$, there exists a stationary solution w such that $u(x, t_n; u_0) \rightharpoonup w$ in $W_0^{1,p}(\Omega)$.*

THEOREM 1.4 *If $u(x, t; u_0)$ is a global solution of (1), then the ω -limit set of u contains a stationary solution w .*

The rest of this paper is organized as follows: In Section 2, we prove Theorems 1.1 and 1.2. In Section 3 we prove Theorems 1.3 and 1.4.

2 THE PROOF OF THEOREMS 1.1 AND 1.2

In this section we consider the existence and the time-asymptotic estimates of global solutions and finite time blowup of (1.1). We first prove Theorem 1.1.

Proof of Theorem 1.1 In fact, we can prove a more general result:

If there exists some t_0 such that $J(u(t_0)) \leq 0$, then $u(x, t; u_0)$ blows up in finite time.

We shall employ the classical concavity method (see [4,5,8,12,17]). Suppose that $t_{\max} = \infty$ and denote $f(t) = \frac{1}{2} \int_{t_0}^t \|u\|_2^2 ds$. Performing standard manipulations

$$\int_{t_0}^t \int_{\Omega} u_t^2 dx ds + \frac{1}{p} \int_{\Omega} |\nabla u|^p dx - \frac{1}{p^*} \int_{\Omega} |u|^{p^*} dx = J(u(t_0)), \quad (2.1)$$

$$f'(t) = \frac{1}{2} \|u_0\|_2^2 + \int_{t_0}^t \int_{\Omega} (-|\nabla u|^p + |u|^{p^*}) dx ds, \quad (2.2)$$

$$f''(t) = - \int_{\Omega} |\nabla u|^p dx + \int_{\Omega} |u|^{p^*} dx.$$

By (2.1), we have

$$\begin{aligned} f''(t) &\geq - \int_{\Omega} |\nabla u|^p dx + \frac{p^*}{p} \int_{\Omega} |\nabla u|^p dx - p^* J(u(t_0)) + p^* \int_{t_0}^t \int_{\Omega} u_t^2 dx ds \\ &= \left(\frac{p^*}{p} - 1 \right) \int_{\Omega} |\nabla u|^p dx - p^* J(u(t_0)) + p^* \int_{t_0}^t \int_{\Omega} u_t^2 dx ds. \end{aligned} \quad (2.3)$$

From the assumption, $J(u(t_0)) \leq 0$ such that

$$\left(\frac{p^*}{p} - 1\right) \int_{\Omega} |\nabla u|^p \, dx - p^* J(u(t_0)) > 0,$$

for all $t \geq t_0$. If we had $t_{\max} = \infty$, then this inequality would yield

$$\lim_{t \rightarrow \infty} f'(t) = \lim_{t \rightarrow \infty} f(t) = \infty,$$

and

$$f''(t) \geq p^* \int_{t_0}^t \int_{\Omega} u_t^2 \, dx \, ds,$$

and

$$\begin{aligned} f(t)f''(t) &\geq \frac{p^*}{p} \left(\int_{t_0}^t \|u(s)\|_2^2 \, ds \right) \left(\int_{t_0}^t \|u_s(s)\|_2^2 \, ds \right) \\ &\geq \frac{p^*}{p} \left(\int_{t_0}^t \int_{\Omega} uu_t \, dx \, ds \right)^2 = \frac{p^*}{p} (f'(t) - f'(t_0))^2, \end{aligned}$$

and as $t \rightarrow \infty$ we have for some $\alpha > 0$ and $\forall t \geq t_0$ such that

$$f(t)f''(t) \geq (1 + \alpha)(f'(t))^2.$$

Hence $f(t)^{-\alpha}$ is concave on $[t_0, \infty]$. But $f(t)^{-\alpha} > 0$ and $\lim_{t \rightarrow \infty} f(t)^{-\alpha} = 0$. This contradiction proves that $t_{\max} < \infty$; which completes the proof of Theorem 1.1.

Proof of Theorem 1.2 We divide the proof into several steps

Step 1 Proof of Existence

(i) *A priori estimates and local existence* From [6] and [10], for each $n > 0$, there is a unique classical solution $u_n \in C(Q_T) \cap C^{2,1}(Q_T)$ of the following equation:

$$\begin{aligned} u_t &= \operatorname{div} \left(\left(|\nabla u|^2 + \frac{1}{n} \right)^{(p-2)/2} \nabla u \right) + \min\{n, u^{p^*-1}\}, \quad (x, t) \in \Omega \times (0, T), \\ u(x, t) &= 0, \quad (x, t) \in \partial\Omega \times (0, T), \\ u(x, 0) &= u_{n0}, \end{aligned} \tag{2.4}$$

where $u_{n0} \in C_0^\infty(\Omega)$, such that

$$u_{n0} \rightarrow u_0, \quad \text{strongly in } W_0^{1,p}(\Omega),$$

and $J(u_{n0}) < (1/N)S^{N/p}$, $\int_\Omega |u_{n0}|^{p^*} dx < S^{N/p}$. On the other hand, multiplying (2.4) by u_{nt} and integrating, we have

$$\iint_{Q_{\tau_1}} u_{nt}^2 dx dt + \frac{1}{p} \int_\Omega |\nabla u_n|^p dx - \frac{1}{p^*} \int_\Omega |u_n|^{p^*} dx \leq J(u_{n0}). \quad (2.5)$$

For the sake of convenience, define:

$$\Sigma = \left\{ u \mid u \in W_0^{1,p}(\Omega), u \geq 0, u \neq 0, J(u) < \frac{1}{N}S^{N/p}, \int_\Omega |u|^{p^*} dx < S^{N/p} \right\}.$$

Now we show that

$$u_n(t) \in \Sigma, \quad \text{for any } t \geq 0.$$

Suppose that it does not hold and let t^* be the smallest time for which $u_n(t^*) \notin \Sigma$. Then in virtue of the continuity of $u_n(t)$ we see that $u_n(t^*) \in \partial\Sigma$. Hence

$$J(u_n(t^*)) = \frac{1}{N}S^{N/p},$$

or

$$\int_\Omega |\nabla u_n|^p dx = \int_\Omega |u_n|^{p^*} dx,$$

which contradicts to (2.5). Then from (2.5) and note that if $\int_\Omega |u|^{p^*} dx < S^{N/p}$, then $\int_\Omega |\nabla u|^p dx > \int_\Omega |u|^{p^*} dx$, we have

$$\int_0^t \|u'_n(s)\|_{L^2(\Omega)}^2 ds + \frac{1}{N} \int_\Omega |\nabla u_n|^p dx \leq J(u_{n0}) < \frac{1}{N}S^{N/p}. \quad (2.6)$$

Thus, we obtain

$$\iint_{Q_{\tau_1}} u_{nt}^2 dx dt < \frac{1}{N}S^{N/p}, \quad (2.7)$$

$$\int_\Omega |\nabla u_n|^p dx < S^{N/p}. \quad (2.8)$$

From (2.8) we have

$$|\nabla u_n|_{L^p(Q_{T_1})} \leq C(T_1), \quad (2.9)$$

where $C(T_1)$ is the constant independent of n . From the prior estimates (2.7), (2.8) and (2.9), we see that there exists a subsequence (not relabeled) and a function u such that

$$\begin{aligned} u_n &\rightharpoonup u, \quad u_n^{p^*-1} \rightharpoonup u^{p^*-1}, \quad \text{a.e. on } Q_{T_1}, \\ \nabla u_n &\rightharpoonup \nabla u, \quad \text{weakly in } L^p(Q_{T_1}), \\ u_{n_t} &\rightharpoonup u_t, \quad \text{weakly in } L^2(Q_{T_1}), \\ u_n &\rightharpoonup u, \quad \text{in } L^\infty(0, T_1; W_0^{1,p}(\Omega)) \text{ weak star}, \\ |\nabla u_n|^{p-2} \nabla u_n &\rightharpoonup w, \quad \text{weakly in } L^{p/(p-1)}(Q_{T_1}). \end{aligned}$$

Then well known arguments of the theory of monotone operators yields $w = |\nabla u|^{p-2} \nabla u$; which implies the function u is a desired local solution of the problem (1.1).

(ii) *Global existence* To prove that it is a global solution. Multiplying (1.1) by u_t and integrating, we obtain

$$\int_0^t \|u'(s)\|_2^2 ds + J(u(x, t)) = J(u_0) < \frac{1}{N} S^{N/p}.$$

Note if $\int_\Omega |u|^{p^*} dx < S^{N/p}$, then $\int_\Omega |\nabla u|^p dx > \int_\Omega |u|^{p^*} dx$. Thus $J(u(x, t)) < (1/N)S^{N/p}$ for any $t > 0$. Now we prove $u(x, t) \in \Sigma$, for any $t > 0$. Indeed, if there exists a t^* such that $u(x, t) \in \partial\Sigma$, then we have $J(u(x, t)) \geq (1/N)S^{N/p}$, which is a contradiction. Hence $\int_\Omega |\nabla u(t)|^p dx > \int_\Omega |u(t)|^{p^*} dx$ for any $t > 0$. Therefore

$$\int_0^t \|u'(s)\|_2^2 dx + \frac{1}{N} \int_\Omega |\nabla u|^p dx \leq J(u_0) < \frac{1}{N} S^{N/p},$$

which implies

$$\int_\Omega |\nabla u|^p dx < S^{N/p}, \quad (2.10)$$

$$\|u'(s)\|_{L^2(0, T; L^2(\Omega))} < \frac{1}{N} S^{N/p}, \quad (2.11)$$

for any $T > 0$. Thus $u(x, t)$ is a global solution of (1.1); which completes Step 1.

Step 2 Proof of (1.3) Note if $\int_{\Omega} |u|^{p^*} dx < S^{N/p}$, then $\int_{\Omega} |\nabla u|^p dx > \int_{\Omega} |u|^{p^*} dx$. Thus $J(u(x, t)) < (1/N)S^{N/p}$ for any $t > 0$. It is easy to prove $u(x, t) \in \Sigma$, for any $t > 0$. Indeed, if there exists a t^* such that $u(x, t) \in \partial\Sigma$, then we have $J(u(x, t)) \geq (1/N)S^{N/p}$, which is a contradiction. Hence $\int_{\Omega} |\nabla u(t)|^p dx > \int_{\Omega} |u(t)|^{p^*} dx$ for any $t > 0$. Let

$$h(u(t)) = \int_{\Omega} |\nabla u|^p dx - \int_{\Omega} u^{p^*} dx,$$

then

$$h(u(t)) > 0, \quad \text{for all } t \geq 0.$$

By Sobolev inequality

$$\int_{\Omega} |u|^{p^*} dx < (1/S^{p^*/p}) \left(\int_{\Omega} |\nabla u|^p dx \right)^{p^*/p},$$

and

$$J(u_0) > \frac{1}{N} \int_{\Omega} |\nabla u|^p dx,$$

implies

$$\int_{\Omega} |u|^{p^*} dx < (1/S^{p^*/p})(NJ(u_0))^{p^*/p-1} \int_{\Omega} |\nabla u|^p dx. \quad (2.12)$$

For simplicity, denote $(1/S^{p^*/p})(NJ(u_0))^{p^*/p-1}$ by $0 < \delta < 1$. Let $\gamma = 1 - \delta$, we have

$$\int_{\Omega} |u(t)|^{p^*} dx \leq (1 - \gamma) \int_{\Omega} |\nabla u(t)|^p dx. \quad (2.13)$$

Let $T > t_0$ be a fixed number, then from

$$\frac{1}{2} \frac{d}{dt} \int_{\Omega} |u(t)|^2 dx = -h(u(t))$$

and Poincaré's inequality, we have

$$\begin{aligned} \int_t^T h(u(s)) \, ds &= \frac{1}{2} \int_{\Omega} |u(t)|^2 \, dx - \frac{1}{2} \int_{\Omega} |u(T)|^2 \, dx \\ &\leq \frac{1}{2} \int_{\Omega} |u(t)|^2 \, dx \leq \frac{1}{2\lambda_1} \int_{\Omega} |\nabla u(t)|^2 \, dx \\ &\leq C(\Omega) \left(\int_{\Omega} |\nabla u(t)|^p \, dx \right)^{2/p}, \end{aligned} \quad (2.14)$$

where λ_1 is the first eigenvalue of $-\Delta_p u = \lambda |u|^{p-2} u$, $x \in \Omega$, $u = 0$, $x \in \partial\Omega$. Furthermore, (2.13) implies

$$\begin{aligned} J(u(t)) &= \frac{1}{p} \int_{\Omega} |\nabla u(t)|^p \, dx - \frac{1}{p^*} \int_{\Omega} |u(t)|^{p^*} \, dx \\ &= \frac{1}{p} \int_{\Omega} |\nabla u(t)|^p \, dx + \frac{1}{p^*} \left[h(u(t)) - \int_{\Omega} |\nabla u(t)|^p \, dx \right] \\ &= \frac{1}{N} \int_{\Omega} |\nabla u(t)|^p \, dx + \frac{1}{p^*} h(u(t)) \geq \frac{1}{N} \int_{\Omega} |\nabla u(t)|^p \, dz, \end{aligned} \quad (2.15)$$

on $[t_0, \infty)$. Therefore, by (2.13) and (2.14) we obtain

$$\int_t^T h(u(s)) \, ds \leq C(\Omega) (J(u(t)))^{2/p}, \quad (2.16)$$

on $[t_0, T]$. On the other hand, (2.13) implies

$$\gamma \int_{\Omega} |\nabla u(t)|^p \, dx \leq h(u(t)), \quad (2.17)$$

on $[t_0, \infty)$. By (2.15) and (2.17), we have

$$J(u(t)) \leq \left(\frac{1}{N\gamma} + \frac{1}{p^*} \right) h(u(t)). \quad (2.18)$$

Further (2.16) and (2.18) give

$$C_1 \int_t^T J(u(s)) \, ds \leq (J(u(t)))^{2/p}$$

on $[t_0, T]$, where

$$C_1 = \left(C(\Omega) \left(\frac{1}{N\gamma} + \frac{1}{p^*} \right) \right)^{-1}.$$

Then, from the arbitrariness of $T > t_0$, we have

$$C_1 \int_t^\infty J(u(s)) \, ds \leq (J(u(t)))^{2/p},$$

i.e.

$$C_1^{p/2} \left(\int_t^\infty J(u(s)) \, ds \right)^{p/2} \leq -\frac{d}{dt} \int_t^\infty J(u(s)) \, ds. \quad (2.19)$$

Setting $y(t) = \int_t^\infty J(u(s)) \, ds$, it follows from (2.19), we have

$$\frac{dy(t)}{dt} \leq -C_1^{p/2} y(t)^{p/2}.$$

Performing standard manipulations, we have

$$y(t) \leq C_2 t^{-2/(p-2)}.$$

Thus, we obtain

$$TJ(u(T+t)) \leq \int_t^{T+t} J(u(s)) \, ds \leq \int_t^\infty J(u(s)) \, ds \leq C_2 t^{-2/(p-2)},$$

By (2.15) we have

$$\frac{1}{N} \int_\Omega |\nabla u(T)|^p \, dx \leq J(u(T)) \leq C_3 t^{-2/(p-2)},$$

with some constant $C_3 > 0$ for enough large $t > T$. Hence

$$\int_\Omega |\nabla u(t)|^p \, dx = O(t^{-2/(p-2)}), \quad \text{as } t \rightarrow \infty.$$

Step 3 Proof of (1.4) Obviously

$$\|\nabla u(x, t; u_0)\|_p^p \leq r^{N/p} < S^{N/p}, \quad (2.20)$$

and

$$\frac{d}{dt} \int_{\Omega} |u(t)|^2 dx + \int_{\Omega} |\nabla u|^p dx = \int_{\Omega} |u|^{p^*} dx, \quad \text{for all } t > 0. \quad (2.21)$$

By the same argument with Step 2, we have

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} |u|^2 dx &< -(1 - \delta) \int_{\Omega} |\nabla u|^p dx \\ &\leq -\frac{(1 - \delta)}{\lambda_1} \int_{\Omega} |u|^p dx \leq -C \left(\int_{\Omega} |u|^2 dx \right)^{p/2}, \end{aligned}$$

where λ_1 is the first eigenvalue of $-\Delta_p u = \lambda |u|^{p-2} u$, $x \in \Omega$, $u = 0$, $x \in \partial\Omega$,

$$C = \frac{(1 - \delta)}{\lambda_1 |\Omega|^{(p-2)/2}}.$$

Let $y = \|u\|_2^2$, we see that the estimate

$$\frac{dy}{dt} \leq -C y^{p/2}$$

Therefore, we have

$$y^{-(p-2)/2} \geq \left(\int_{\Omega} |u_0|^2 dx \right)^{(p-2)/2} + C \frac{p-2}{2} t,$$

which shows

$$y(t) = O(t^{-2/(p-2)}), \quad \text{as } t \rightarrow \infty,$$

which completes the proof of (1.4).

Step 4 Proof of Theorem 1.2 (2) We divide the proof into two steps.

(i) First of all, we define a set which consists of the functions that satisfy the following conditions:

$$J(u_0) < \frac{1}{N} S^{N/p}, \quad (2.22)$$

$$\int_{\Omega} |u_0|^{p^*} dx = S^{N/p}. \quad (2.23)$$

We claim that the set is an empty set. In fact, let u_0 belong to the set. If u_0 satisfies

$$\int_{\Omega} |\nabla u_0|^p \, dx \leq \int_{\Omega} |u_0|^{p^*} \, dx,$$

then

$$S^{N/p} = \int_{\Omega} |u_0|^{p^*} \, dx \geq \int_{\Omega} |\nabla u_0|^p \, dx \geq S \left(\int_{\Omega} |u_0|^{p^*} \, dx \right)^{p/p^*} = S^{N/p},$$

and hence

$$\begin{aligned} \int_{\Omega} |\nabla u_0|^p \, dx &= \int_{\Omega} |u_0|^{p^*} \, dx = S^{N/p}, \\ J(u_0) &= \frac{1}{p} \int_{\Omega} |\nabla u_0|^p \, dx - \frac{1}{p^*} \int_{\Omega} |u_0|^{p^*} \, dx = \frac{1}{N} S^{N/p}, \end{aligned}$$

which is contradictory to (2.22).

If u_0 satisfies:

$$\int_{\Omega} |\nabla u_0|^p \, dx > \int_{\Omega} |u_0|^{p^*} \, dx,$$

then from (2.22) we see that

$$\frac{1}{N} S^{N/p} > J(u_0) = \frac{1}{p} \int_{\Omega} |\nabla u_0|^p \, dx - \frac{1}{p^*} \int_{\Omega} |u_0|^{p^*} \, dx > \frac{1}{N} \int_{\Omega} |u_0|^{p^*} \, dx.$$

Implies

$$\int_{\Omega} |u_0|^{p^*} \, dx < S^{N/p}$$

which is a contradiction because of (2.23).

(ii) Thus we consider only the case of u_0 satisfies

$$J(u) < \frac{1}{N} S^{N/p}, \quad \int_{\Omega} |u|^{p^*} \, dx > S^{N/p}. \quad (2.24)$$

Obviously, in this case we have

$$S^{N/p} < \int_{\Omega} |\nabla u_0|^p \, dx < \int_{\Omega} |u_0|^{p^*} \cdot dx$$

If $u(x, t)$ is a global solution then we can deduce that $u(x, t)$ does not converge strongly to 0 in $W_0^{1,p}(\Omega)$. Otherwise, $\exists t^*, 0 < t^* < \infty$ such that

$$J(u(t^*)) < \frac{1}{N} S^{N/p}, \quad \int_{\Omega} |u(t^*)|^{p^*} \, dx = S^{N/p},$$

which is a contradiction from the first half (i). Now we prove the following claim:

CLAIM *If u_0 satisfies (2.24) and $u(x, t; u_0)$ is a global solution. Then $\forall t \in [0, T]$ the following inequalities hold:*

$$S^{N/p} < \int_{\Omega} |\nabla u(x, t)|^p \, dx < \int_{\Omega} |u(x, t)|^{p^*} \, dx. \quad (2.25)$$

Indeed, if there exists a t^ such that $\int_{\Omega} |\nabla u(x, t^*)|^p \, dx = \int_{\Omega} |u(x, t^*)|^{p^*} \, dx$, then we have $\int_{\Omega} |\nabla u(x, t^*)|^p \, dx = \int_{\Omega} |u(x, t^*)|^{p^*} \, dx \geq S^{N/p}$. But $(1/N)S^{N/p} > J(u(x, t^*)) = (1/N) \int_{\Omega} |\nabla u(x, t^*)|^p \, dx$, with a contradiction. Therefore there exists a constant $\eta > 0$ sufficiently small and independent of t , rely on u_0 such that*

$$\int_{\Omega} |u(x, t)|^{p^*} \, dx \geq (1 + \eta) \int_{\Omega} |\nabla u(x, t)|^p \, dx, \quad (2.26)$$

for any $t \in [0, \infty]$, which completes the proof of the claim.

From the claim and $p > 2$ we have

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} u^2 \, dx &= \int_{\Omega} u^{p^*} \, dx - \int_{\Omega} |\nabla u|^p \, dx \\ &\geq \eta \int_{\Omega} |\nabla u|^p \, dx \geq C \left(\int_{\Omega} |u|^2 \, dx \right)^{p/2}, \end{aligned} \quad (2.27)$$

implies

$$\frac{d}{dt} (\|u\|_2^2) \geq C \|u\|_2^p,$$

i.e.

$$\int_{\|u_0\|_2^2}^{\|u\|_2^2} y^{-(p/2)} dy > Ct,$$

and therefore

$$CT \leq \int_{\|u_0\|_2^2}^{\infty} y^{-(p/2)} dy < +\infty$$

which completes the proof of Theorem 1.2.

3 THE PROOF OF THEOREMS 1.3 AND 1.4

First of all, we prove Theorem 1.3.

Proof of Theorem 1.3 For any $t_n \rightarrow \infty$, let $u_n = u(x, t_n; u_0)$, from the boundedness we know that there exists a subsequence (still denote by $\{u_n\}$) and a function w such that

$$\begin{aligned} u_n &\rightharpoonup w && \text{in } W_0^{1,p}(\Omega), \\ u_n^{p^*-1} &\rightharpoonup w^{p^*-1} && \text{in } (L^{p^*}(\Omega))^*, \\ u_n &\rightarrow w && \text{a.e. in } \Omega. \end{aligned}$$

In order to pass to the limit in (1.1) we first fix some $T < \infty$ and introduce suitable test functions similar to Fila [4]. Take

$$\psi \in W_0^{1,p}(\Omega), \quad \rho \in C_0^2(0, T), \quad \rho \geq 0, \quad \int_0^T \rho(t) dt = 1.$$

Put

$$\varphi(x, t) = \begin{cases} \rho(t - t_n)\psi(x) & \text{for } t > t_n, \quad x \in \bar{\Omega}, \\ 0 & \text{for } 0 \leq t \leq t_n, \quad x \in \bar{\Omega}. \end{cases}$$

Further, we obtain from Definition 1.1 that

$$\int_{t_n}^{t_n+T} \int_{\Omega} [u\rho'(t - t_n)\psi - \rho|\nabla u|^{p-2}\nabla u\nabla\psi + u^{p^*-1}\rho(t - t_n)\psi] dx dt = 0.$$

The transformation $s = t - t_n$, leads to

$$\begin{aligned} & \int_0^T \int_{\Omega} [u(t_n + s)\rho'(s)\psi - \rho|\nabla u(t_n + s)|^{p-2}\nabla u(t_n + s)\nabla\psi \\ & + u(t_n + s)^{p^*-1}\rho(s)\psi] dx ds = 0. \end{aligned} \quad (3.1)$$

Note that the uniform boundedness of $u(t_n + s)$ in $W_0^{1,p}(\Omega)$ for $0 \leq s \leq T$. Therefore, we can choose the same subsequence of $\{t_n\}$ (not relabeled) and functions w_s and w such that

$$u(t_n + s) \rightarrow w_s, \quad \text{strongly in } L^q(\Omega) (p \leq q < p^*)$$

and

$$u(t_n) \rightarrow w, \quad \text{strongly in } L^q (p \leq q < p^*).$$

Now we claim: $w_s = w$. Indeed, by the energy inequality we have

$$\int_{\Omega} |u(t_n + s) - u(t_n)|^2 dx = s \int_{t_n}^{t_n+s} \int_{\Omega} \left| \frac{\partial u}{\partial \tau} \right|^2 dx d\tau \rightarrow 0, \quad \text{as } t_n \rightarrow \infty$$

as $0 \leq s \leq T$ for any fixed $T < \infty$. Thus, we have

$$u(t_n + s) - u(t_n) \rightarrow 0, \quad \text{strongly in } L^2(\Omega), \quad \text{as } t_n \rightarrow \infty$$

for $0 \leq s \leq T$ for any fixed $T < \infty$. Hence

$$w_s = w$$

which prove the claim.

Now we rewrite (3.1) as follows:

$$\begin{aligned} & \int_0^T \int_{\Omega} [u(t_n)\rho'(s)\psi - \rho|\nabla u(t_n)|^{p-2}\nabla u(t_n)\nabla\psi + u(t_n)^{p^*-1}\rho(s)\psi] dx ds \\ & + \int_0^T \int_{\Omega} [u(t_n + s) - u(t_n)]\rho'(s)\psi dx ds \\ & - \int_0^T \int_{\Omega} [|\nabla u(t_n + s)|^{p-2}\nabla u(t_n + s) - |\nabla u(t_n)|^{p-2}\nabla u(t_n)]\nabla\psi dx ds \\ & + \int_0^T \int_{\Omega} [u(t_n + s)^{p^*-1} - u(t_n)^{p^*-1}]\rho(s)\psi dx ds = 0 \end{aligned} \quad (3.2)$$

By the dominated convergence theorem and the choice of ρ and $u(t_n) \rightarrow w$ strongly in $L^q(\Omega)$ ($p \leq q < p^*$), we have

$$\int_0^T \rho \left[\int_{\Omega} |\nabla u(t_n)|^{p-2} \nabla u(t_n) \nabla \psi \, dx - \int_{\Omega} u(t_n)^{p^*-1} \psi \, dx \right] ds = o(1),$$

as $n \rightarrow \infty$

Denote $u(t_n)$ by u_n , From the choice of ρ , we obtain

$$\int_{\Omega} |\nabla u(t_n)|^{p-2} \nabla u(t_n) \nabla \psi \, dx - \int_{\Omega} u(t_n)^{p^*-1} \psi \, dx = o(1), \quad \text{as } n \rightarrow \infty.$$

Thus, we have

$$\int_{\Omega} |\nabla w|^{p-2} \nabla w \nabla \psi \, dx = \int_{\Omega} w^{p^*-1} \psi \, dx, \quad \text{for all } \psi \in W_0^{1,p}(\Omega)$$

which completes the proof of Theorem 1.3.

Proof of Theorem 1.4 From now on, denote $u(x, t; u_0)$ by u , we have

$$\int_0^\infty \int_{\Omega} u_t^2 \, dx \, ds \leq C < \infty.$$

Then there exists a sequence $\{t_n\}$ satisfying $t_n \rightarrow \infty$ as $n \rightarrow \infty$ such that

$$\int_{\Omega} |u_t(x, t_n; u_0)|^2 \, dx \rightarrow 0 \quad \text{as } n \rightarrow \infty. \tag{3.3}$$

For the sake of convenience, denote $u(x, t_n; u_0)$ by u_n . From Theorem 1.1, $J(u(t)) > 0$ for all $t \geq 0$, and

$$0 < J(u(t)) \leq J(u_0). \tag{3.4}$$

If we consider the time sequence $\{t_n\}$ as

$$0 < J(u(t_n)) \leq J(u_0). \tag{3.5}$$

The statement of (3.3) and (3.5) says that $u_n = u(t_n)$, $t_n \rightarrow \infty$ is a Palais–Smale sequence related to the statement problem of (1.1). Such a situation has been well studied in the theory of nonlinear elliptic

equations. It is easy to prove that there exists a constant $C < +\infty$ such that

$$\int_{\Omega} |\nabla u_n|^p dx \leq C.$$

Thus, there exists a subsequence (not relabeled) and a function w such that

$$\begin{aligned} u_n &\rightharpoonup w, \quad \text{weakly in } W_0^{1,p}(\Omega), \\ u_n &\rightarrow w, \quad \text{strongly in } L^q(\Omega) (p \leq q < p^*). \end{aligned}$$

From the theory of elliptic equation we can obtain that w is a stationary solution, which completes the proof of Theorem 1.4.

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