Quenching for semidiscretizations of a semilinear heat equation with Dirichlet and Neumann boundary conditions

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Abstract. This paper concerns the study of the numerical approximation for the following boundary value problem:

\[
\begin{cases}
  u_t(x, t) - u_{xx}(x, t) = -u^{-p}(x, t), & 0 < x < 1, \ t > 0, \\
  u_x(0, t) = 0, & u(1, t) = 1, \ t > 0, \\
  u(x, 0) = u_0(x) > 0, & 0 \leq x \leq 1,
\end{cases}
\]

where \( p > 0 \). We obtain some conditions under which the solution of a semidiscrete form of the above problem quenches in a finite time and estimate its semidiscrete quenching time. We also establish the convergence of the semidiscrete quenching time. Finally, we give some numerical experiments to illustrate our analysis.

Keywords: semidiscretizations, discretizations, heat equations, quenching, semidiscrete quenching time, convergence

Classification: 35K55, 35B40, 65M06

1. Introduction

Consider the following boundary value problem:

(1) \( u_t(x, t) - u_{xx}(x, t) = -u^{-p}(x, t), \ 0 < x < 1, \ t > 0, \)

(2) \( u_x(0, t) = 0, \ u(1, t) = 1, \ t > 0, \)

(3) \( u(x, 0) = u_0(x) > 0, \ 0 \leq x \leq 1, \)

where \( p > 0, \ u_0'(0) = 0, \ u_0(1) = 1, \ u_0(x) < 1 \) for \( x \in [0, 1) \).

Definition 1.1. We say that a solution \( u \) of (1)–(3) quenches in a finite time if there exists a finite time \( T_q \) such that \( \|u(x, t)\|_{\text{inf}} > 0 \) for \( t \in [0, T_q) \), but

\[
\lim_{t \to T_q} \|u(x, t)\|_{\text{inf}} = 0,
\]

where \( \|u(x, t)\|_{\text{inf}} = \min_{0 \leq x \leq 1} u(x, t) \). The time \( T_q \) is called the quenching time of the solution \( u \).
The theoretical study of solutions for semilinear heat equations which quench in a finite time has been the subject of investigations of many authors (see [2], [4]–[8] and the references cited therein). Under some conditions, the authors have proved that the solution $u$ of (1)–(3) quenches in a finite time and have given some estimates of the quenching time.

In this paper, we are interested in the numerical study of the phenomenon of quenching using a semidiscrete form of (1)–(3). We give some conditions under which the solution of the semidiscrete form quenches in a finite time and estimate its semidiscrete quenching time. We also prove that the semidiscrete quenching time converges to the real one when the mesh size goes to zero. A similar study has been undertaken by some authors concerning the phenomenon of blow-up (we say that a solution blows up in a finite time if it takes an infinite value in a finite time)(see [1]). In [3], some schemes have been used to study the phenomenon of extinction.

This paper is organised as follows. In the next section, we construct a semidiscrete scheme and give some lemmas which will be used later. In Section 3, under some conditions, we prove that the solution of a semidiscrete form of (1)–(3) quenches in a finite time and estimate its semidiscrete quenching time. In Section 4, we study the convergence of the semidiscrete quenching time. Finally, in the last section, we give some numerical results to illustrate our analysis.

2. A semidiscrete problem

In this section, we give some lemmas which will be used later. We start by the construction of a semidiscrete scheme as follows. Let $I$ be a positive integer, and define the grid $x_i = ih$, $0 \leq i \leq I$, where $h = 1/I$. Approximate the solution $u$ of the problem (1)–(3) by the solution $U_h(t) = (U_0(t), U_1(t), \ldots, U_I(t))^T$ of the following semidiscrete equations

\begin{align}
\frac{dU_i(t)}{dt} &= \delta^2 U_i(t) - (U_i(t))^{-p}, \quad 0 \leq i \leq I - 1, \quad t \in (0, T^h_q), \\
U_I(t) &= 1, \quad t \in (0, T^h_q), \quad U_i(0) = \varphi_i > 0, \quad 0 \leq i \leq I,
\end{align}

where $\varphi_i < 1$ for $0 \leq i \leq I - 1$,

\begin{align}
\delta^2 U_i(t) &= \frac{U_{i+1}(t) - 2U_i(t) + U_{i-1}(t)}{h^2}, \quad 1 \leq i \leq I - 1, \\
\delta^2 U_0(t) &= \frac{2U_1(t) - 2U_0(t)}{h^2}.
\end{align}

Here $(0, T^h_q)$ is the maximal time interval on which $\|U_h(t)\|_{\text{inf}} > 0$, where $\|U_h(t)\|_{\text{inf}} = \min_{0 \leq i \leq I} U_i(t)$. When $T^h_q$ is finite, then we say that the solution $U_h(t)$ of (4)–(5) quenches in a finite time, and the time $T^h_q$ is called the semidiscrete quenching time of the solution $U_h(t)$.

The following lemma is a semidiscrete form of the maximum principle.
**Lemma 2.1.** Let \( \alpha_h \in C^0([0,T), \mathbb{R}^{I+1}) \) and let \( V_h \in C^1([0,T), \mathbb{R}^{I+1}) \) be such that

\[
\frac{dV_i(t)}{dt} - \delta^2 V_i(t) + \alpha_i(t)V_i(t) \geq 0, \quad 0 \leq i \leq I - 1, \quad t \in (0,T),
\]

(6)

\[
V_i(t) \geq 0, \quad t \in (0,T),
\]

(7)

\[
V_i(0) \geq 0, \quad 0 \leq i \leq I.
\]

(8)

Then \( V_i(t) \geq 0 \) for \( 0 \leq i \leq I, \ t \in (0,T) \).

**Proof:** Let \( T_0 < T \) and define the vector \( Z_h(t) = e^{\lambda t}V_h(t) \), where \( \lambda \) is such that \( \alpha_i(t) - \lambda > 0, \ 0 \leq i \leq I, \ t \in [0,T_0] \). Let

\[
m = \min_{0 \leq i \leq I, 0 \leq t \leq T_0} Z_i(t).
\]

For \( i = 0, \ldots, I, \ Z_i(t) \) is a continuous function on the compact \([0,T_0]\). Then, there exist \( i_0 \in \{0,1,\ldots,I\} \) and \( t_0 \in [0,T_0] \) such that \( m = Z_{i_0}(t_0) \). If \( i_0 \in \{0,1,\ldots,I-1\} \), then we observe that

\[
\frac{dZ_{i_0}(t_0)}{dt} = \lim_{k \to 0} \frac{Z_{i_0}(t_0) - Z_{i_0}(t_0 - k)}{k} \leq 0,
\]

(9)

\[
\delta^2 Z_{i_0}(t_0) = \delta^2 Z_0(t_0) = \frac{2Z_1(t_0) - 2Z_0(t_0)}{h^2} \geq 0 \quad \text{if} \quad i_0 = 0,
\]

(10)

\[
\delta^2 Z_{i_0}(t_0) = \frac{Z_{i_0+1}(t_0) - 2Z_{i_0}(t_0) + Z_{i_0-1}(t_0)}{h^2} \geq 0 \quad \text{if} \quad 1 \leq i_0 \leq I - 1.
\]

(11)

Using (6), a straightforward computation yields

\[
\frac{dZ_{i_0}(t_0)}{dt} - \delta^2 Z_{i_0}(t_0) + (\alpha_{i_0}(t_0) - \lambda)Z_{i_0}(t_0) \geq 0 \quad \text{if} \quad 0 \leq i_0 \leq I - 1.
\]

(12)

From the inequalities (9)–(12), it is not hard to see that \( (\alpha_{i_0}(t_0) - \lambda)Z_{i_0}(t_0) \geq 0, \ 0 \leq i_0 \leq I - 1 \). Due to (7) and the fact that \( \alpha_{i_0}(t_0) - \lambda > 0 \), we see that \( Z_h(t_0) \geq 0 \). We deduce that \( V_h(t) \geq 0 \) for \( t \in [0,T_0] \) which leads us to the desired result.

The lemma below shows a property of the semidiscrete solution.

**Lemma 2.2.** Let \( U_h \) be the solution of (4)–(5). Then

\[
U_i(t) < 1, \quad 0 \leq i \leq I - 1, \quad t \in (0,T_q^h).
\]

(13)
Proof: Let \( t_0 \) be the first \( t \in (0, T^h_q) \) such that \( U_i(t) < 1 \) for \( t \in [0, t_0) \), \( 0 \leq i \leq I - 1 \), but \( U_{i_0}(t_0) = 1 \) for a certain \( i_0 \in \{0, \ldots, I - 1\} \). We observe that

\[
\frac{dU_{i_0}(t_0)}{dt} = \lim_{k \to 0} \frac{U_{i_0}(t_0) - U_{i_0}(t_0 - k)}{k} \geq 0, \tag{14}
\]

\[
\delta^2 U_{i_0}(t_0) = \frac{U_{i_0+1}(t_0) - 2U_{i_0}(t_0) + U_{i_0-1}(t_0)}{h^2} \leq 0 \quad \text{if} \quad 1 \leq i_0 \leq I - 1, \tag{15}
\]

\[
\delta^2 U_{i_0}(t_0) = \delta^2 U_0(t_0) = \frac{2U_1(t_0) - 2U_0(t_0)}{h^2} \leq 0 \quad \text{if} \quad i_0 = 0, \tag{16}
\]

which implies that

\[
\frac{dU_{i_0}(t_0)}{dt} - \delta^2 U_{i_0}(t_0) + (U_{i_0}(t_0))^{-p} > 0.
\]

But, this contradicts (4) and the proof is complete. \( \square \)

Another version of the maximum principle for semidiscrete equations is the following comparison lemma.

**Lemma 2.3.** Let \( f \in C^0(\mathbb{R} \times \mathbb{R}, \mathbb{R}) \). If \( V_h, W_h \in C^1([0,T), \mathbb{R}^{I+1}) \) are such that

\[
\frac{dV_i(t)}{dt} - \delta^2 V_i(t) + f(V_i(t), t) < \frac{dW_i(t)}{dt} - \delta^2 W_i(t) + f(W_i(t), t), \tag{17}
\]

\( 0 \leq i \leq I - 1, \ t \in (0, T), \)

\[
V_i(t) < W_i(t), \ t \in (0, T), \tag{18}
\]

\[
V_i(0) < W_i(0), \ 0 \leq i \leq I, \tag{19}
\]

then \( V_i(t) < W_i(t), \ 0 \leq i \leq I, \ t \in (0, T). \)

Proof: Let \( Z_h(t) = W_h(t) - V_h(t) \) and let \( t_0 \) be the first \( t > 0 \) such that \( Z_i(t) > 0 \) for \( t \in [0, t_0) \), \( 0 \leq i \leq I \), but \( Z_{i_0}(t_0) = 0 \) for a certain \( i_0 \in \{0, \ldots, I\} \). We observe that

\[
\frac{dZ_{i_0}(t_0)}{dt} = \lim_{k \to 0} \frac{Z_{i_0}(t_0) - Z_{i_0}(t_0 - k)}{k} \leq 0, \tag{17a}
\]

\[
\delta^2 Z_{i_0}(t_0) = \frac{Z_{i_0+1}(t_0) - 2Z_{i_0}(t_0) + Z_{i_0-1}(t_0)}{h^2} \geq 0 \quad \text{if} \quad 1 \leq i_0 \leq I - 1, \tag{17b}
\]

\[
\delta^2 Z_{i_0}(t_0) = \frac{2Z_1(t_0) - 2Z_0(t_0)}{h^2} \geq 0 \quad \text{if} \quad i_0 = 0. \tag{17c}
\]

Therefore if \( i_0 \in \{0, \ldots, I - 1\} \), then we have

\[
\frac{dZ_{i_0}(t_0)}{dt} - \delta^2 Z_{i_0}(t_0) + f(W_{i_0}(t_0), t_0) - f(V_{i_0}(t_0), t_0) < 0,
\]

which contradicts (17). If \( i_0 = I \), then we have a contradiction because of (18). This ends the proof. \( \square \)

The lemma below reveals a property of the operator \( \delta^2 \).
Lemma 2.4. Let $V_h$ and $U_h \in \mathbb{R}^{I+1}$. If $\delta^+(U_0) \geq 0$ and
$\delta^+(U_i) \geq 0, \quad \delta^-(U_i) \geq 0, \quad 1 \leq i \leq I-1,$
then
$$\delta^2(U_i) \geq U_i \delta^2 V_i + V_i \delta^2 U_i, \quad 0 \leq i \leq I-1,$$
where $\delta^+(U_i) = \frac{U_{i+1} - U_i}{h}$, $\delta^-(U_i) = \frac{U_{i-1} - U_i}{h}$.

Proof: A straightforward computation yields
$$\delta^2(U_0 V_0) = 2 \delta^+(U_0) \delta^+(V_0) + U_0 \delta^2 V_0 + V_0 \delta^2 U_0,$$
$$\delta^2(U_i V_i) = \delta^+(U_i) \delta^+(V_i) + \delta^-(U_i) \delta^-(V_i) + U_i \delta^2 V_i + V_i \delta^2 U_i, \quad 1 \leq i \leq I-1.$$
Using the assumptions of the lemma, we obtain the desired result.

The following result shows another property of the semidisc rate solution.

Lemma 2.5. Let $U_h$ be the solution of (4)–(5) such that the initial data at (5) satisfy
$$(20) \quad \varphi_{i+1} > \varphi_i, \quad 0 \leq i \leq I-1.$$ Then, we have
$$(21) \quad U_{i+1}(t) > U_i(t), \quad 0 \leq i \leq I-1, \quad t \in (0,T^h_q).$$

Proof: Let $t_0 \in (0,T^h_q)$ be the first $t > 0$ such that $U_{i+1}(t) > U_i(t)$ for $t \in (0,t_0)$,
$0 \leq i \leq I-1$, but
$$U_{k+1}(t_0) = U_k(t_0) \quad \text{for a certain} \quad k \in \{0, \ldots, I-1\}.$$ Without loss of generality, we may suppose that $k$ is the smallest integer which
satisfies the above equality.
If $k = I-1$ then $U_I(t_0) = U_{I-1}(t_0) = 1$. But, this contradicts Lemma 2.2. If
$k \in \{0, \ldots, I-2\}$, then letting $Z_k(t) = U_{k+1}(t) - U_k(t)$, we observe that
$$\frac{dZ_k(t_0)}{dt} = \lim_{k \to 0} \frac{Z_k(t_0) - Z_k(t_0 - k)}{k} \leq 0,$$
$$\delta^2 Z_k(t_0) = \frac{Z_k(t_0) - Z_k(t_0 - k)}{h^2} \leq 0 \quad \text{if} \quad k = 0,$$
$$\delta^2 Z_k(t_0) = \frac{Z_{k+1}(t_0) - 2Z_k(t_0) + Z_{k-1}(t_0)}{h^2} > 0 \quad \text{if} \quad 1 \leq k \leq I-2.$$ Therefore, if $0 \leq k \leq I-2$, we get
$$\frac{dZ_k(t_0)}{dt} - \delta^2 Z_k(t_0) + (U_{k+1}(t_0))^{-p} - (U_k(t_0))^{-p} < 0,$$
which contradicts (4). This ends the proof.

Remark 2.1. The above result reveals that if the initial data of the semidiscrete solution are increasing in space, then the semidiscrete solution is also increasing in space. This property will be used later to show that the semidiscrete solution
attains its minimum at the first node.
3. Quenching in the semidiscrete problem

In this section, under some assumptions, we show that the solution $U_h$ of (4)–(5) quenches in a finite time and estimate its semidiscrete quenching time.

Let us give another property of the operator $\delta^2$ useful in this section.

**Lemma 3.1.** Let $U_h \in \mathbb{R}^{I+1}$ such that $U_h > 0$. Then, we have

$$\delta^2 u_i^p - p u_i^{p-1} \delta^2 u_i \quad \text{for} \quad 0 \leq i \leq I - 1.$$  

**Proof:** Apply Taylor’s expansion to obtain

$$\delta^2 u_0^p = -p u_0^{p-1} \delta^2 u_0 + (U_1 - U_0) \frac{2(p+1)}{h^2} \theta_0^{p-2},$$

$$\delta^2 u_i^p = -p u_i^{p-1} \delta^2 u_i + (U_{i+1} - U_i) \frac{2(p+1)}{2h^2} \theta_i^{p-2}$$

$$+ (U_{i-1} - U_i) \frac{2(p+1)}{2h^2} \eta_i^{p-2} \quad \text{if} \quad 1 \leq i \leq I - 1,$$

where $\theta_i$ is an intermediate value between $U_{i+1}$ and $U_i$ and $\eta_i$ the one between $U_{i-1}$ and $U_i$. Use the fact that $U_h > 0$ to complete the rest of the proof. \hfill \Box

Our result about the quenching time is the following.

**Theorem 3.1.** Let $U_h$ be the solution of (4)–(5). Assume that there exists a constant $A > 0$ such that the initial data at (5) satisfy

(22) $$\delta^2 \psi_i - \psi_i^p \leq -A \cos(ih\frac{\pi}{2}) \psi_i^p, \quad 0 \leq i \leq I - 1,$$

(23) $$1 - \frac{\pi^2}{2A(p+1)} ||\psi_h||_{\inf}^{p+1} > 0.$$

If (20) holds, then $U_h$ quenches in a finite time $T_q^h$ which satisfies the following estimate

$$T_q^h < -\frac{8}{\pi^2} \ln \left( 1 - \frac{\pi^2}{2A(p+1)} ||\psi_h||_{\inf}^{p+1} \right).$$

**Proof:** Since $(0, T_q^h)$ is the maximal time interval on which $||U_h(t)||_{\inf} > 0$, our aim is to show that $T_q^h$ is finite and satisfies the above inequality. Introduce the vector $J_h(t)$ such that

$$J_i(t) = \frac{dU_i(t)}{dt} + C_i(t) U_i^{-p}(t), \quad 0 \leq i \leq I, \quad t \in [0, T_q^h),$$
where $C_i(t) = Ae^{-\lambda_h t} \cos(ih \frac{\pi}{2})$ with $\lambda_h = \frac{2-2\cos(h\frac{\pi}{2})}{h^2}$. It is not hard to see that

$$
\frac{dC_i(t)}{dt} - \delta^2 C_i(t) = 0, \quad C_{i+1}(t) < C_i(t), \quad 0 \leq i \leq I - 1.
$$

Using Lemma 2.5, we observe that

$$
\delta^+(U_0^{-p})\delta^+(C_0) \geq 0 \quad \text{and} \quad \delta^+(U_i^{-p})\delta^+(C_i) \geq 0, \quad \delta^-(U_i^{-p})\delta^-(C_i) \geq 0
$$

for $1 \leq i \leq I - 1$. A straightforward computation gives

$$
\frac{dJ_i(t)}{dt} - \delta^2 J_i(t) = \frac{d}{dt} \left( \frac{dU_i(t)}{dt} - \delta^2 U_i(t) \right) + U_i^{-p} \frac{dC_i(t)}{dt} - pC_i(t)U_i^{-p-1} \frac{dU_i(t)}{dt} - \delta^2 (C_i(t)U_i^{-p}(t)), \quad 0 \leq i \leq I - 1.
$$

It follows from (25), Lemmas 2.4 and 3.1 that

$$
\delta^2 (C_i(t)U_i^{-p}(t)) \geq U_i^{-p}(t)\delta^2 C_i(t) - pC_i(t)U_i^{-p-1}(t)\delta^2 U_i(t), \quad 0 \leq i \leq I - 1.
$$

We deduce that

$$
\frac{dJ_i(t)}{dt} - \delta^2 J_i(t) \leq \frac{d}{dt} \left( \frac{dU_i(t)}{dt} - \delta^2 U_i(t) \right) - pC_i(t)U_i^{-p-1} \left( \frac{dU_i(t)}{dt} - \delta^2 U_i(t) \right) + U_i^{-p}(t) \left( \frac{dC_i(t)}{dt} - \delta^2 C_i(t) \right), \quad 0 \leq i \leq I - 1.
$$

In virtue of (4) and (24), we arrive at

$$
\frac{dJ_i(t)}{dt} - \delta^2 J_i(t) \leq pU_i^{-p-1}(t)J_i(t), \quad 0 \leq i \leq I - 1, \quad t \in (0, T_q^h).
$$

Obviously, $J_I(t) = 0$. From the assumption (22), we get $J_h(0) \leq 0$. It follows from Lemma 2.1 that $J_h(t) \leq 0$ for $t \in (0, T_q^h)$. This estimate may be rewritten as follows

$$
\frac{dU_i(t)}{dt} \leq -Ae^{-\lambda_h t} \cos(ih \frac{\pi}{2})U_i^{-p}(t), \quad 0 \leq i \leq I, \quad t \in (0, T_q^h).
$$

We observe that $\lambda_h \leq \frac{\pi^2}{2}$ for $h$ small enough. Hence, we get

$$
U_i^p(0) dU_0(t) \leq -Ae^{-\frac{\pi^2}{2} t} dt \quad \text{for} \quad t \in (0, T_q^h).
$$

From Lemma 2.5, $U_0(t) = \|U_h(t)\|_{\text{inf}}$. Therefore, integrating (26) over $(0, T_q^h)$, we obtain

$$
T_q^h \leq -\frac{8}{\pi^2} \ln(1 - \frac{\pi^2}{2A(p+1)\|U_h(0)\|_{\text{inf}}^{p+1}}).
$$

Use the fact that $U_h(0) = \varphi_h$ and (23) to complete the rest of the proof. \qed
Remark 3.1. Assume that there exists a time $t_0 \in (0, T^h_q)$ such that

$$1 - \frac{\pi^2}{2A(p+1)} e^{-\frac{\pi^2}{2} t_0} \|U_h(t_0)\|_{\inf}^{p+1} > 0.$$ 

Integrating the inequality (26) over $(t_0, T^h_q)$, and using the fact that $U_0(t_0) = \|U_h(t_0)\|_{\inf}$, we arrive at

$$T^h_q - t_0 \leq -\frac{8\pi^2}{\pi^2} \ln \left(1 - \frac{\pi^2}{2A(p+1)} e^{-\frac{\pi^2}{2} t_0} \|U_h(t_0)\|_{\inf}^{p+1}\right).$$

Remark 3.2. It is easy to find a vector $\varphi_h$ and a positive constant $A$ such that (22), (23) hold. In fact, one may find a vector $\psi_h$ and a constant $A \in (0, 1)$ such that

$$\delta^2 \psi_i - \psi_i^{-p} \leq -A \psi_i^{-p}, \quad 0 \leq i \leq I - 1,$$

which implies that

$$\delta^2 \psi_i - \psi_i^{-p} \leq -A \cos(ih\frac{\pi}{2}) \psi_i^{-p}, \quad 0 \leq i \leq I - 1.$$

Let $\varphi_h = \varepsilon \psi_h$ where $0 < \varepsilon < 1$. It is not hard to see that

$$\delta^2 \varphi_i - \varphi_i^{-p} \leq -A \cos(ih\frac{\pi}{2}) \varphi_i^{-p}, \quad 0 \leq i \leq I - 1,$$

and the inequality (22) follows. To obtain (23), it suffices to take $\varepsilon$ small enough.

4. Convergence of the semidiscrete quenching time

In this section, under some assumptions, we prove that the semidiscrete quenching time converges to the real one when the mesh size goes to zero.

We denote

$$u_h(t) = (u(x_0, t), \ldots, u(x_I, t))^T.$$

In order to obtain the convergence of the semidiscrete quenching time, we firstly prove the following theorem about the convergence of the semidiscrete scheme.

Theorem 4.1. Assume that the problem (1)–(3) has a solution $u \in C^{4,1}([0, 1] \times [0, T])$ such that $\min_{0 \leq t \leq T} \|u(x, t)\|_{\inf} = \rho > 0$ and the initial data at (5) satisfy

$$\|\varphi_h - u_h(0)\|_{\infty} = o(1) \quad \text{as} \quad h \to 0.$$ 

Then, for $h$ sufficiently small, the problem (4)–(5) has a unique solution $U_h \in C^1([0, T], \mathbb{R}^{I+1})$ such that

$$\max_{0 \leq t \leq T} \|U_h(t) - u_h(t)\|_{\infty} = O(\|\varphi_h - u_h(0)\|_{\infty} + h^2) \quad \text{as} \quad h \to 0.$$
PROOF: Problem (4)–(5) has for each \( h \) a unique solution \( U_h \in C^1([0,T^h q), \mathbb{R}^{I+1}) \).

Let \( t(h) \) be the greatest value of \( t > 0 \) such that

\[
\|U_h(t) - u_h(t)\|_{\infty} < \frac{\rho}{2} \quad \text{for} \quad t \in (0,t(h)).
\]

Relation (27) implies that \( t(h) > 0 \) for \( h \) sufficiently small. Let \( t^*(h) = \min\{t(h), T\} \). From the triangle inequality, we get

\[
\|U_h(t)\|_{\infty} \geq \|u_h(t)\|_{\infty} - \|U_h(t) - u_h(t)\|_{\infty} \quad \text{for} \quad t \in (0,t^*(h)),
\]

which implies that

\[
\|U_h(t)\|_{\infty} \geq \rho - \frac{\rho}{2} = \frac{\rho}{2} \quad \text{for} \quad t \in (0,t^*(h)).
\]

Consider the error

\[
e_h(t) = U_h(t) - u_h(t).
\]

By a direct calculation, we find that for \( t \in (0,t^*(h)) \),

\[
\frac{d e_i(t)}{dt} - \delta^2 e_i(t) = p(\Theta_i(t))^{-p-1} e_i(t) + \frac{h^2}{12} u_{xxxx}(\tilde{x}_i,t), \quad 0 \leq i \leq I - 1,
\]

where \( \Theta_i \) is an intermediate value between \( U_i(t) \) and \( u(x_i,t) \). Let \( M > 0 \) be such that

\[
\frac{\|u_{xxxx}(x,t)\|_{\infty}}{12} \leq M \quad \text{for} \quad t \in [0,T], \quad p(\frac{\rho}{2})^{p-1} \leq M.
\]

Using (30)–(31), it is not hard to see that

\[
\frac{d e_i(t)}{dt} - \delta^2 e_i(t) \leq M|e_i(t)| + Mh^2, \quad 0 \leq i \leq I - 1, \quad t \in (0,t^*(h)).
\]

Introduce the vector \( z_h(t) \) such that

\[
z_i(t) = e^{(M+1)t}(\|\varphi_h - u_h(0)\|_{\infty} + Mh^2), \quad 0 \leq i \leq I, \quad t \in [0,T].
\]

A straightforward computation yields

\[
\frac{d z_i(t)}{dt} - \delta^2 z_i(t) > M|z_i(t)| + Mh^2, \quad 0 \leq i \leq I - 1, \quad t \in (0,t^*(h)),
\]

\[
z_I(t) > e_I(t), \quad t \in (0,t^*(h)),
\]

\[
z_i(0) > e_i(0), \quad 0 \leq i \leq I.
\]
It follows from Comparison Lemma 2.3 that
\[ z_i(t) > e_i(t) \quad \text{for} \quad t \in (0, t^*(h)), \quad 0 \leq i \leq I. \]

In the same way, we also show that
\[ z_i(t) > -e_i(t) \quad \text{for} \quad t \in (0, t^*(h)), \quad 0 \leq i \leq I, \]
which implies that
\[ \|U_h(t) - u_h(t)\|_{\infty} \leq e^{(M+1)t}(\|\varphi_h - u_h(0)\|_{\infty} + Mh^2), \quad t \in (0, t^*(h)). \]

Let us show that \( t^*(h) = T \). Suppose that \( T > t(h) \). From (29), we obtain
\[ \frac{\rho}{2} = \|U_h(t(h)) - u_h(t(h))\|_{\infty} \leq e^{(M+1)T}(\|\varphi_h - u_h(0)\|_{\infty} + Mh^2). \]

Since the term on the right hand side of the above inequality goes to zero as \( h \) tends to zero, we deduce that \( \frac{\rho}{2} \leq 0 \), which is impossible. Consequently \( t^*(h) = T \) and the proof is complete.

Now, we are in a position to prove the main result of this section.

**Theorem 4.2.** Suppose that the solution \( u \) of (1)–(3) quenches in a finite time \( T_q \) such that \( u \in C^{4,1}([0, 1] \times [0, T_q]) \) and the initial data at (5) satisfy condition (27). Under the assumptions of Theorem 3.1, problem (4)–(5) admits a unique solution \( U_h(t) \) which quenches in a finite time \( T^h_q \) with \( \lim_{h \to 0} T^h_q = T_q \).

**Proof:** Let \( 0 < \varepsilon < T_q/2 \). There exists a constant \( R > 0 \) such that
\[ -\frac{8}{\pi^2} \ln \left(1 - \frac{\pi^2}{2A(p+1)}x^{p+1}\right) < \frac{\varepsilon}{2} \quad \text{for} \quad x \in [0, R]. \]

Since \( u \) quenches in a finite time \( T_q \), then there exists \( T_1 \in (T_q - \varepsilon/2, T_q) \) such that
\[ 0 < \|u(x, t)\|_{\inf} < \frac{R}{2} \quad \text{for} \quad t \in (T_1, T_q). \]

Let \( T_2 = \frac{T_1 + T_q}{2} \). Obviously, we have \( 0 < \|u(x, t)\|_{\inf} < \frac{R}{2} \) for \( t \in [0, T_2] \). It follows from Theorem 4.1 that
\[ \|U_h(t) - u_h(t)\|_{\infty} < \frac{R}{2} \quad \text{for} \quad t \in [0, T_2], \]
which implies that
\[ \|U_h(T_2) - u_h(T_2)\|_{\infty} < \frac{R}{2}. \]
Applying the triangle inequality, we obtain
\[ \| U_h(T_2) \|_\infty \leq \| U_h(T_2) - u_h(T_2) \| + \| u_h(T_2) \|_\infty \leq \frac{R}{2} + \frac{R}{2} = R. \]

We deduce from Remark 3.1 and (36) that
\[ |T_q^h - T_q| \leq -\frac{8\pi^2}{\pi^2} \ln \left( 1 - \frac{\pi^2}{2A(p+1)} e^{-\pi^2 T_2 \| U_h(T_2) \|_{\infty}^{p+1}} \right) < \frac{\varepsilon}{2}. \]

Consequently, we find that
\[ |T_q^h - T_q| \leq |T_q^h - T_2| + |T_2 - T_q| \leq \frac{\varepsilon}{2} + \frac{\varepsilon}{2} = \varepsilon, \]
and the proof is complete.

\[ \square \]

5. Numerical experiments

In this section, we consider the problem (1)–(3) in the case where \( p = 1, \)
\( u_0(x) = 0.05 + 0.95\sin(\frac{\pi}{2}x). \) We give some computational results concerning
some approximations of the real quenching time. We start by proposing some
schemes which will be used later for our numerical experiments.

At first, we approximate the solution \( u(x,t) \) of the problem (1)–(3) by the
solution \( U_h(n) = (U_{0}^{(n)}, U_{1}^{(n)}, \ldots, U_{I}^{(n)})^T \) of the following explicit scheme
\[
\frac{U_i^{(n+1)} - U_i^{(n)}}{\Delta t_n} = \delta^2 U_i^{(n)} - (U_i^{(n)})^{-p-1}U_i^{(n+1)}, \quad 0 \leq i \leq I - 1,
\]
\[ U_I^{(n)} = 1, \quad U_0^{(n)} = \varphi_i, \quad 0 \leq i \leq I, \]
where \( n \geq 0. \) In order to permit the discrete solution to reproduce the properties
of the continuous one when the time \( t \) approaches the quenching time \( T_q, \) we need
to adapt the size of the time step so that we take \( \Delta t_n = \min \{ \frac{h^2}{2}, \tau \| U_h^{(n)} \|_{\infty}^{p+1} \} \)
with \( \tau = \text{const} \in (0, 1). \) Let us notice that the restriction on the time step ensures
the positivity of the discrete solution.

At second, we approximate the solution \( u(x,t) \) of the problem (1)–(3) by the
solution \( U_h(n) \) of the implicit scheme below
\[
\frac{U_i^{(n+1)} - U_i^{(n)}}{\Delta t_n} = \delta^2 U_i^{(n+1)} - (U_i^{(n)})^{-p-1}U_i^{(n+1)}, \quad 0 \leq i \leq I - 1,
\]
\[ U_I^{(n)} = 1, \quad U_0^{(n)} = \varphi_i, \quad 0 \leq i \leq I, \]
where \( n \geq 0. \) As in the case of the explicit scheme, here, we choose \( \Delta t_n = \tau \| U_h^{(n)} \|_{\infty}^{p+1} \) with \( \tau = \text{const} \in (0, 1). \) For the implicit scheme, the existence and
positivity of the discrete solution is also guaranteed using standard methods (see,
for instance, [3]).

In both schemes, we take \( \varphi_i = 0.05 + 0.95\sin(\frac{\pi}{2}ih), \quad \tau = h^2. \)
We need the following definition.
Definition 5.1. We say that the solution $U_h^{(n)}$ of the explicit scheme or the implicit scheme quenches in a finite time if $\lim_{n \to +\infty} \|U_h^{(n)}\|_{\text{inf}} = 0$ and the series $\sum_{n=0}^{+\infty} \Delta t_n$ converges. The quantity $\sum_{n=0}^{+\infty} \Delta t_n$ is called the numerical quenching time of the discrete solution $U_h^{(n)}$.

In Tables 1 and 2, in rows, we present the numerical quenching times, the number of iterations, CPU times and the orders of the approximations corresponding to meshes of 16, 32, 64, 128. We take for the numerical quenching time $T^n = \sum_{j=0}^{n-1} \Delta t_j$ which is computed at the first time when

$$\Delta t_n = |T^{n+1} - T^n| \leq 10^{-16}.$$ 

The order $(s)$ of the method is computed from

$$s = \log((T_{4h} - T_{2h})/(T_{2h} - T_h))/\log(2).$$

Table 1:

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References


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